

Lecture 10

Oct 17, 2018

$$\# \rho \frac{D\xi_z}{Dt} = \mu \nabla^2 \xi_z \quad \text{if } \rho \neq 0, \text{ incompressible}$$

$$\# \xi_z = -\nabla^2 \psi$$

$$\# \vec{v} = \nabla \phi \quad \text{velocity potential}$$

$$u = \frac{\partial \phi}{\partial x} \quad v = \frac{\partial \phi}{\partial y} \quad w = \frac{\partial \phi}{\partial z} \quad \text{irrotational}$$

$$\vec{\xi} = \nabla \times \vec{v} = \nabla \times (\nabla \phi) = \vec{0}$$

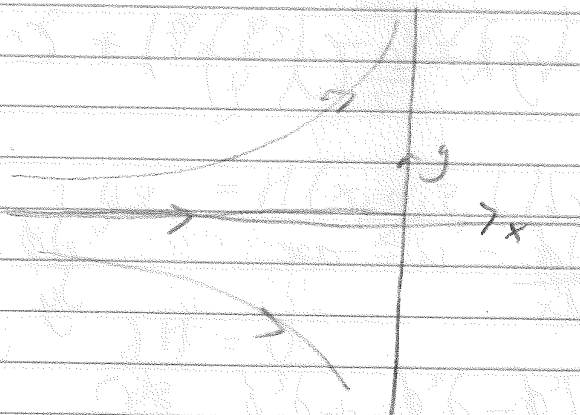
For incompressible flow, continuity becomes

$$\nabla \cdot \vec{v} = \nabla \cdot (\nabla \phi) = \boxed{\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2} = 0}$$

The velocity potential ϕ satisfies Laplace eqn.
e.g. stagnation point flow

$$u(x,y) = -ax$$

$$v(x,y) = ay$$



$$\nabla \cdot \vec{v} = \frac{du}{dx} + \frac{dv}{dy} = -a + a = 0 \quad \checkmark$$

$$\xi_z = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 0$$

Incompressible, irrotational \rightarrow a velocity potential ϕ exists, satisfying: $u = \frac{\partial \phi}{\partial x} = -ax$

$$v = \frac{\partial \phi}{\partial y} = ay$$

$$\int u dx \Rightarrow \phi = \int u dx + f(y) \\ = \frac{-ax^2}{2} + f(y)$$

$$\int v dy \Rightarrow \phi = \int v dy + g(x) \\ = \frac{ay^2}{2} + g(x)$$

Solutions must be the same

$$f(y) = \frac{ay^2}{2} + c_1$$

$$g(x) = \frac{-ax^2}{2} + c_2$$

$$\phi(x,y) = \frac{a}{2}(y^2 - x^2) + C$$

Check

$$\frac{\partial \phi}{\partial x} = u = -ax$$

$$\frac{\partial \phi}{\partial y} = v = ay$$

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0$$

Bernoulli's eqn

Smits 4.2

Consider incompressible, irrotational flow

$$\nabla^2 \phi = 0$$

Consider momentum eqn. for incompressible, inviscid flow ($\mu=0$)

$$\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} = -\frac{1}{\rho} \nabla p + \vec{g} + \frac{\mu}{\rho} \nabla^2 \vec{v}$$

Irrotational + inviscid go hand in hand. $\Rightarrow \nabla^2 \vec{v} = 0$

Irrotational: $\vec{v} = \nabla \phi$

Vector identity: $\nabla \cdot \vec{v} = 0$

$$(\vec{v} \cdot \nabla) \vec{v} = \frac{\nabla |\vec{v}|^2}{2} - \vec{v} \times \zeta$$

Gravity: $\vec{g} = \nabla(-gz) = \langle 0, 0, -g \rangle$

$$\Rightarrow \frac{\partial(\nabla \phi)}{\partial t} + \frac{\nabla |\vec{v}|^2}{2} - \vec{v} \times \zeta = \nabla \left(\frac{p}{\rho} \right) - \nabla(gz)$$

$$H = \frac{\partial \phi}{\partial t} + \frac{|\vec{v}|^2}{2} + \frac{p}{\rho} + gz = H(x, y, z, t)$$

$$\frac{\partial H}{\partial x} = 0$$

$$\frac{\partial H}{\partial y} = 0$$

$$\frac{\partial H}{\partial z} = 0$$

$$\frac{\partial \phi}{\partial t} + \frac{|\vec{v}|^2}{2} + \frac{p}{\rho} + gz = H(t)$$

unsteady
Bernoulli

Steady flow $\frac{d\phi}{dt} = 0$

$H_0 = \text{constant}$

$$\frac{|\vec{v}|^2}{2} + \frac{p}{\rho} + gz = 0$$

Notes

→ assumed irrotational, incompressible, inviscid

→ dimensions of energy although it comes from momentum

→ other versions only require inviscid

→ Bernoulli's $\Rightarrow p$ from \vec{v} or ϕ

eg. stagnation point

$$u = -ax \quad v = ay$$

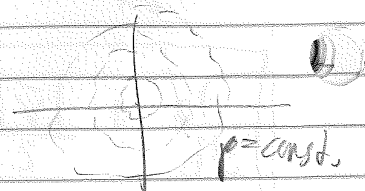
steady, incomp., neglect gravity

$$\frac{1}{2}(u^2 + v^2) + \frac{p}{\rho} = \frac{1}{2}a^2(x^2 + y^2) + \frac{p}{\rho} = H_0$$

Evaluate H_0 at stagnation point $(0,0)$
 $p = p_0$

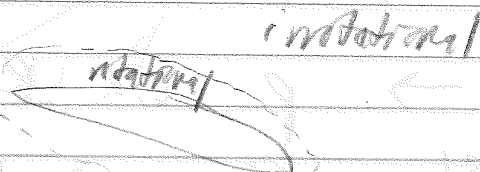
$$\frac{1}{2}a^2(x^2 + y^2) + \frac{p}{\rho} = \frac{p_0}{\rho}$$

$$p(x,y) = p_0 - \frac{1}{2}\rho a^2(x^2 + y^2)$$



Pressure is highest @ stagnation pt, decreases
as $v^2 = v_x^2 + v_y^2$

A general approach to flows w/ substantial
inviscid region

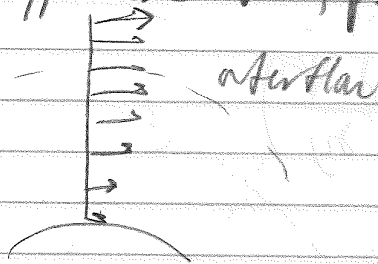


1) Solve for $\nabla^2 \phi = 0$ + BC \rightarrow gives $\vec{v}_I = \nabla \phi$

2) Use Bernoulli's $\rightarrow p_I(\vec{x})$

3) Compute pressure force on object

4) If necessary, use \vec{v}_I, p_I to compute BL flow



5) Compute viscous force

6) If necessary, go back to #1 with the object
to change arc to BL