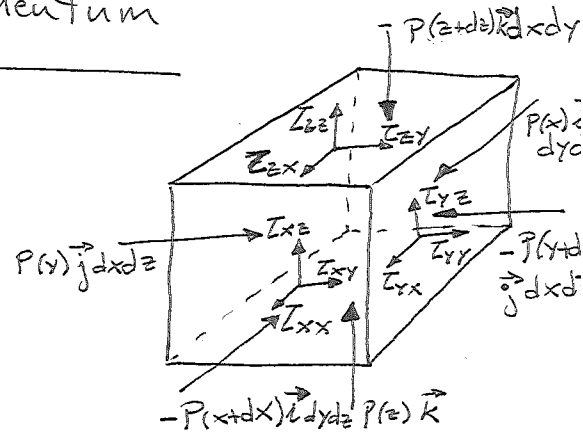
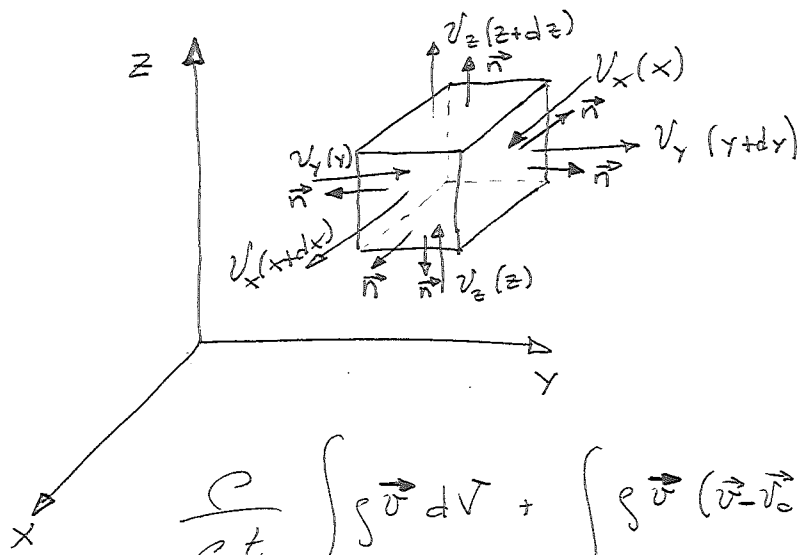


# Conservation of Momentum



$$\frac{\rho}{\rho t} \int_{c.v.} \rho \vec{v} dV + \int_{c.s.} \rho \vec{v} (\vec{v} \cdot \vec{v}_0) \cdot \vec{n} dA = - \int_{c.s.} p \vec{n} dA + \int_{c.s.} \vec{T} \cdot \vec{n} dA + \int_{c.v.} \vec{f} dV$$

$$\lim_{dx \rightarrow 0} \left[ \frac{\rho}{\rho x} \left( \rho \vec{v} v_x \right) dx dy dz + \left[ \rho(x+dx) \vec{v}(x+dx) v_x(x+dx) - \rho(x) \vec{v}(x) v_x(x) \right] dy dz + \left[ \rho(y+dy) \vec{v}(y+dy) v_y(y+dy) - \rho(y) \vec{v}(y) v_y(y) \right] dx dz + \left[ \rho(z+dz) \vec{v}(z+dz) v_z(z+dz) - \rho(z) \vec{v}(z) v_z(z) \right] dx dy \right]$$

$$= - \left[ \frac{\rho p}{\rho x} dx + o(dx)^2 \right] \vec{i} dy dz - \left[ \frac{\rho p}{\rho y} dy + o(dy)^2 \right] \vec{j} dx dz - \left[ \frac{\rho p}{\rho z} dz + o(dz)^2 \right] \vec{k} dx dy$$

$$+ \left[ (T_{xx} \vec{i} + T_{xy} \vec{j} + T_{xz} \vec{k})(x+dx) dy dz - (T_{xx} \vec{i} + T_{xy} \vec{j} + T_{xz} \vec{k})(x) dy dz \right] + \left[ (T_{yx} \vec{i} + T_{yy} \vec{j} + T_{yz} \vec{k})(y+dy) dx dz - (T_{yx} \vec{i} + T_{yy} \vec{j} + T_{yz} \vec{k})(y) dx dz \right] + \left[ (T_{zx} \vec{i} + T_{zy} \vec{j} + T_{zz} \vec{k})(z+dz) dx dy - (T_{zx} \vec{i} + T_{zy} \vec{j} + T_{zz} \vec{k})(z) dx dy \right]$$

$$+ \rho \vec{g} dx dy dz$$

$$\lim_{dx \rightarrow 0} [ ] = \left[ \frac{\rho}{\rho x} (T_{xx} \vec{i} + T_{xy} \vec{j} + T_{xz} \vec{k}) dx + o(dx)^2 \right] dy dz$$

$$\lim_{dy \rightarrow 0} [ ] = \left[ \frac{\rho}{\rho y} (T_{yx} \vec{i} + T_{yy} \vec{j} + T_{yz} \vec{k}) dy + o(dy)^2 \right] dx dz$$

$$\lim_{dz \rightarrow 0} [ ] = \left[ \frac{\rho}{\rho z} (T_{zx} \vec{i} + T_{zy} \vec{j} + T_{zz} \vec{k}) dz + o(dz)^2 \right] dx dy$$

$$\frac{\rho}{\rho t} (\rho \vec{v}) + \frac{\rho}{\rho x} (\rho \vec{v} v_x) + \frac{\rho}{\rho y} (\rho \vec{v} v_y) + \frac{\rho}{\rho z} (\rho \vec{v} v_z) = -\nabla p +$$

$$+ \left( \frac{\rho}{\rho x} \tau_{xx} + \frac{\rho}{\rho y} \tau_{yx} + \frac{\rho}{\rho z} \tau_{zx} \right) \vec{i} + \left( \frac{\rho}{\rho x} \tau_{xy} + \frac{\rho}{\rho y} \tau_{yy} + \frac{\rho}{\rho z} \tau_{zy} \right) \vec{j} +$$

$$+ \left( \frac{\rho}{\rho x} \tau_{xz} + \frac{\rho}{\rho y} \tau_{yz} + \frac{\rho}{\rho z} \tau_{zz} \right) \vec{k} + \rho \vec{g}$$

$$\int \frac{\rho \vec{v}}{\rho t} + \underbrace{\vec{v} \cdot \left( \frac{\rho \rho}{\rho t} + \frac{\rho \rho v_x}{\rho x} + \frac{\rho \rho v_y}{\rho y} + \frac{\rho \rho v_z}{\rho z} \right)}_{\substack{\vec{v} \cdot \left( \frac{\rho \rho}{\rho t} + \rho \nabla \cdot \vec{v} \right) \\ \text{Continuity}}} + \underbrace{\rho v_x \frac{\rho \vec{v}}{\rho x} + \rho v_y \frac{\rho \vec{v}}{\rho y} + \rho v_z \frac{\rho \vec{v}}{\rho z}}_{\rho (\vec{v} \cdot \nabla) \vec{v}} =$$

$$= -\nabla p + \nabla \cdot \vec{\tau}' + \rho \vec{g}$$

$$\boxed{\rho \left[ \frac{\rho \vec{v}}{\rho t} + (\vec{v} \cdot \nabla) \vec{v} \right] = -\nabla p + \nabla \cdot \vec{\tau}' + \rho \vec{g}}$$

$\rho$ ,  $p$ ,  $\vec{v}$  are unknowns in these equations, but we need a constitutive equation for  $\vec{\tau}'$  as a function of the deformation (strain) of the fluid.

$$\vec{\tau}' = f(\vec{\epsilon})$$

In tensor form this translates into

$$\tau'_{ij} = K_{ijkl} \epsilon_{kl}$$

$\uparrow$  4<sup>th</sup> order tensor (3 components in each direction)  
 $27 = 3 \times 3 \times 3$  possible coefficients.

The simplest assumption is that the stress tensor is linearly proportional to the rate of strain.

$K_{ijkl}$  is composed of 81 constants.

If we further assume that the fluid is isotropic, that is to say that the stress developed by the fluid due to the rate of strain in a certain direction does not depend on the direction (the stress/rate of strain relationship does not change under rotation of the reference frame).

$$K_{ijkl} = \lambda \delta_{ij} \delta_{kl} + \mu \delta_{ik} \delta_{jl} + \gamma \delta_{il} \delta_{jk}$$

3 independent variables (vs. 81 possible)

We can show that  $\tau'_{ij} = \tau'_{ji}$ , the stress tensor is symmetric and then  $K_{ijkl} = K_{jikl} = \lambda \delta_{ji} \delta_{kl} + \mu \delta_{jk} \delta_{il} + \gamma \delta_{jl} \delta_{ik}$

$\lambda \delta_{ij} \delta_{kl} + \mu \delta_{ik} \delta_{jl} + \gamma \delta_{jk} \delta_{il}$

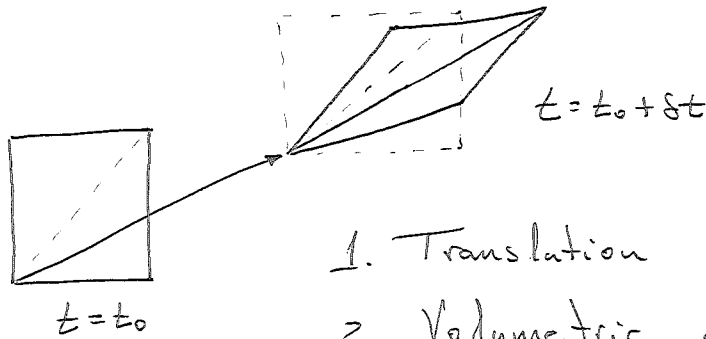
$\mu = \gamma$

$$\tau'_{ij} = \lambda \delta_{ij} \dot{\epsilon}_{kk} + 2\mu \dot{\epsilon}_{ij}$$

$$\underline{\underline{\tau}}' = 2\mu \frac{1}{2} (\nabla \vec{v} + \nabla \vec{v}^T) + \lambda \nabla \cdot \vec{v} \underline{\underline{I}}$$

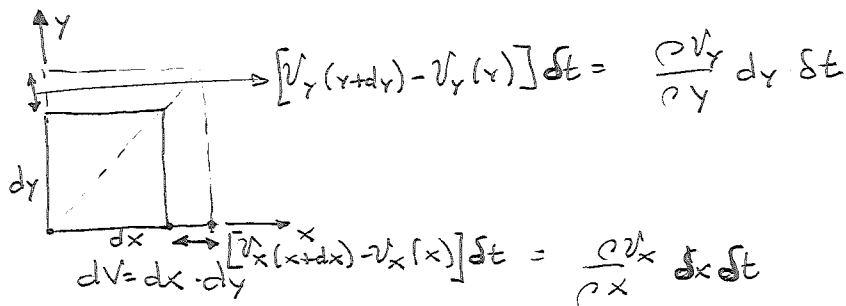
Now we need to relate the rate of strain to some fluid variable used in our analysis.

F.M. White: "Viscous flow" 2<sup>nd</sup> Edition 1991. McGraw Hill  
Section 1.3.3.



1. Translation  $\vec{v} \cdot \delta t$

2. Volumetric deformation:  $\frac{1}{V} \frac{\delta V}{\delta t}$



$$\delta V = \left[ dx + \frac{\partial v_x}{\partial x} dx \delta t \right] \left[ dy + \frac{\partial v_y}{\partial y} dy \delta t \right] - dx dy$$

$$\frac{1}{dV} \frac{\delta dV}{\delta t} = \frac{1}{dx dy} \frac{\frac{\partial v_x}{\partial x} dx dy \delta t + \frac{\partial v_y}{\partial y} dx dy \delta t + \frac{\partial v_x}{\partial x} \frac{\partial v_y}{\partial y} dx dy \delta t}{\delta t}$$

$$\frac{1}{V} \frac{DV}{Dt} = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_x}{\partial x} \frac{\partial v_y}{\partial y} \delta t \xrightarrow{\delta t \rightarrow 0}$$

$$\boxed{\frac{1}{V} \frac{DV}{Dt} = \nabla \cdot \vec{v}} = \dot{\epsilon}_{xx} + \dot{\epsilon}_{yy} + \dot{\epsilon}_{zz}$$

