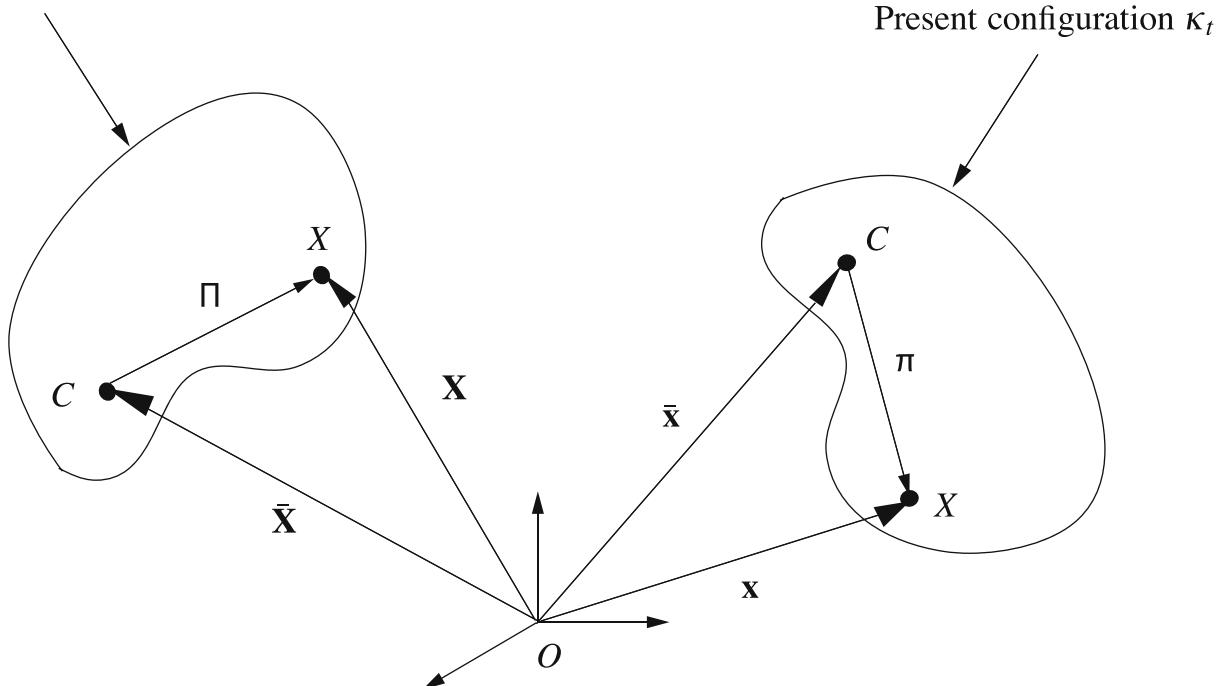


## 8.7 Angular Momenta

Before we can discuss the balance laws of Chapter 9, we must introduce the angular momentum of a rigid body.

Reference configuration  $\kappa_0$



Definition: the angular momentum of B about its center of mass C is

Definition: the angular momentum of B about the fixed point O is

These equations can be related as follows.



See O'Reilly p. 153 for an expression for the angular momentum about an arbitrary point.

## 8.8 Inertia tensor

Expressing the relative position vectors in bases:

Also, we write the angular velocity of the body:

## 8.8.1 The Inertia Tensor

If we consider the angular momentum of  $\text{B}$  about its center of mass  $\bar{H}$ , we can derive the following important relationship:



where

$I$  is the **inertia tensor**, which is written as a matrix in the corotational basis as:

The following components are called **moments of inertia**:

$$I_{xx} = \int_R (\Pi_y^2 + \Pi_z^2) \rho dv = \int_{R_0} (\Pi_y^2 + \Pi_z^2) \rho_0 dV,$$

$$I_{yy} = \int_R (\Pi_x^2 + \Pi_z^2) \rho dv = \int_{R_0} (\Pi_x^2 + \Pi_z^2) \rho_0 dV,$$

$$I_{zz} = \int_R (\Pi_x^2 + \Pi_y^2) \rho dv = \int_{R_0} (\Pi_x^2 + \Pi_y^2) \rho_0 dV,$$

The remaining components are called **products of inertia**:

$$I_{xy} = - \int_R \Pi_x \Pi_y \rho dv = - \int_{R_0} \Pi_x \Pi_y \rho_0 dV,$$

$$I_{xz} = - \int_R \Pi_x \Pi_z \rho dv = - \int_{R_0} \Pi_x \Pi_z \rho_0 dV,$$

$$I_{yz} = - \int_R \Pi_y \Pi_z \rho dv = - \int_{R_0} \Pi_y \Pi_z \rho_0 dV.$$

$I$  is positive-definite, and so its eigenvalues are positive. Its components depend on the basis chosen. If the basis  $(\bar{E}_x, \bar{E}_y, \bar{E}_z)$  is chosen such that  $\bar{E}_x, \bar{E}_y, + \bar{E}_z$  are the eigenvectors of  $I$ , then  $\{\bar{E}_x, \bar{E}_y, \bar{E}_z\}$  and  $\{\bar{e}_x, \bar{e}_y, \bar{e}_z\}$  are called the **principal axes** of the body in its reference and present configurations, respectively.

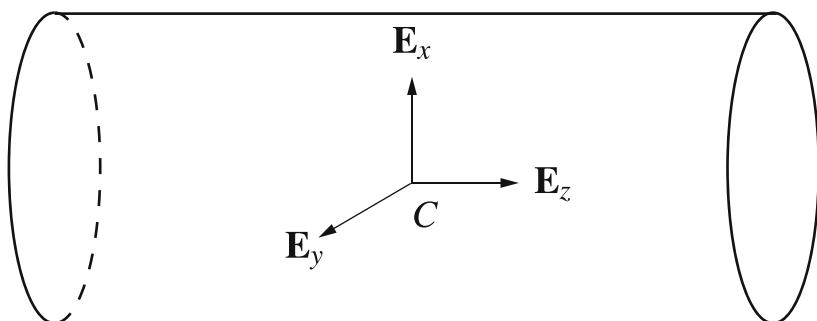
In this case, in the corotational basis,

therefore, we always try to choose  $\bar{E}_x, \bar{E}_y, + \bar{E}_z$  as the principal axes.

### 8.8.3 A Circular Cylinder (example)

What is the inertia tensor for the homogeneous cylinder of mass  $m$ , radius  $R$ , and length  $L$ ?

Choose the principle axes, as shown.



If  $I_{xy} = I_{xz} = I_{yz} = 0$ , (and they are) we have correctly chosen the principal axes.

Computing the remaining integrals,

If we set  $R=0$ , we have the inertia tensor for a "slender rod." With  $L=0$ , we have the inertia tensor for a "thin disk."

#### 8.8.4 The Parallel Axis Theorem + Practical Notes

The Parallel Axis Theorem is commonly used to find the inertia tensor of a point that is not the center of mass. However, we circumvent the need for it by expressing the angular momentum about an arbitrary point A:

We will use this in Chapter 9. Note that this approach is more general than the PAT, which

only applies to points **on** the body.

To find the moments of inertia, we often refer to a table of common shapes of bodies, like that in the cover of Hibbler.