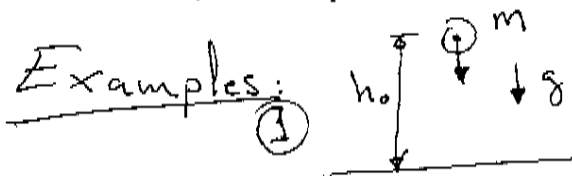


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DIMENSIONAL ANALYSIS

Take the characteristic variables that the problem depends on and build a non-dimensional expression to account for that dependency.



What is the time it takes the stone to reach the ground?

$$\underline{t_f = f(m, g, h_0)}$$

It does not depend on anything else.

Dimensions:

$$[h_0] = L$$

$$[g] = L T^{-2}$$

$$[m] = M$$

$$[t_f] = T$$

We need to have a time on both sides of the equation:

$$\left[\sqrt{\frac{h}{g}} \right] = T$$

$$t_f = \sqrt{h_0/g} \cdot f(m)$$

$\frac{t_f}{\sqrt{h_0/g}} = f(m)$: to be correct this expression needs to be non-dimensional.

Since we cannot make m non-dimensional (no combination of g, h, t_f contains M)

the problem cannot depend on M .

$$\frac{t_f}{\sqrt{h_0/g}} = \text{constant}$$

We know that the solution is $\frac{dV}{dt} = -g; \frac{dh}{dt} = V(t) \Rightarrow t_f = \sqrt{2h_0/g}$ from the solution to the O.D.E

② Pressure loss in a pipe: $\Delta P = f(D, L, S, \mu, v)$



How about Q ?

Q depends on $\frac{\pi D^2}{4}$ and v , so it is not an independent variable. We can use it instead of v (or D) but we can not use all three of them.

$$[\Delta P] = \frac{MLT^{-2}}{L^2} = ML^{-1}T^{-2}$$

$$[D] = L$$

$$[L] = L$$

$$[S] = ML^{-3}$$

$$[\mu] = ML^{-1}T^{-1}$$

$$[v] = LT^{-1}$$

} form a group that has dimensions of pressure: ρv^2

$$\frac{\Delta P}{\rho v^2} = f(D, L, \mu, S, v^2)$$

To complete this we need to form dimensionless groups with the variables we haven't used yet: D, L, μ

An easy one: $\frac{L}{D}$;

the other one $\frac{\mu}{\rho v D}$

$$\frac{\Delta P}{\rho v^2} = f\left(\frac{L}{D}, \frac{\mu}{\rho v D}\right)$$

If we know that the pressure loss is proportional to the length (conditions are uniform along the pipe in terms of geometry, roughness, flow perturbations, etc.) then

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$$\frac{\Delta P}{\rho v^2} = \frac{L}{D} f\left(\frac{\mu}{\rho v D}\right)$$

$$\frac{\Delta P/L}{\rho v^2/D} = f\left(\frac{\mu}{\rho v D}\right)$$

$\frac{\rho v D}{\mu}$ = Reynolds number, and is
THE MOST IMPORTANT
NON-DIMENSIONAL PARAMETER
IN FLUID MECHANICS.
For reasons completely different
from this analysis of flow in
pipes.

To carry out this analysis rigorously, we have
Buckingham - Pi Theorem to guide us:

If an equation involving k variables is dimensionally homogeneous, then it can be reduced to a relationship between $k-r$ dimensionless groups, where r is the minimum number of dimensions involved in the problem.

It is a constructive theorem:

1. Identify the physical variables involved in the problem: (k)
2. Express each variable in term of its fundamental dimensions: $MLT\Theta$ system. (r)
3. Determine $(k-r)$ pi terms are necessary to describe the problem.

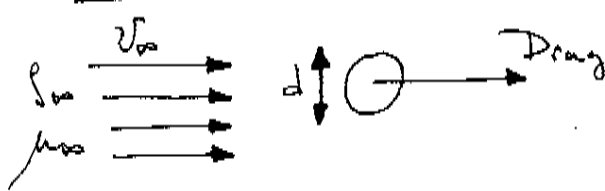
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4. Choose r variables to form the non-dimensional π groups. They need to be chosen with care so that all dimensions are represented independently:

$$\begin{aligned}
 [\phi_1] &= M^{a_1} L^{b_1} T^{c_1} \Theta^{d_1} \dots \\
 [\phi_2] &= M^{a_2} L^{b_2} T^{c_2} \Theta^{d_2} \dots \\
 &\vdots \\
 [\phi_r] &= M^{a_r} L^{b_r} T^{c_r} \Theta^{d_r} \dots
 \end{aligned}
 \quad \text{rank} \begin{pmatrix} a_1 & b_1 & c_1 & d_1 & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ a_r & b_r & c_r & d_r & \dots \end{pmatrix} = r$$

5. Form $k-r$ non dimensional π groups using the k variables chosen in 4 and everyone of the other $k-r$ variables

Example: Drag on a cylinder: $D = f(d, \rho, \mu, v)$



$K = 5$ ①

$$[D] = M L T^{-2}$$

$$[d] = L$$

$$[\rho] = M L^{-3}$$

$$[\mu] = M L^{-1} T^{-1}$$

$$[v_\infty] = L T^{-1}$$

$r = 3$ M, L, T

②

③ $k-r = 2$ non-dimensional parameters.

④ Choose 3 variables: d, ρ, v_∞ (Reynolds number)

$$\frac{D}{\rho v_\infty^2 d^2} = f\left(\frac{\mu}{\rho v_\infty d}\right) \Rightarrow C_D = f(Re)$$

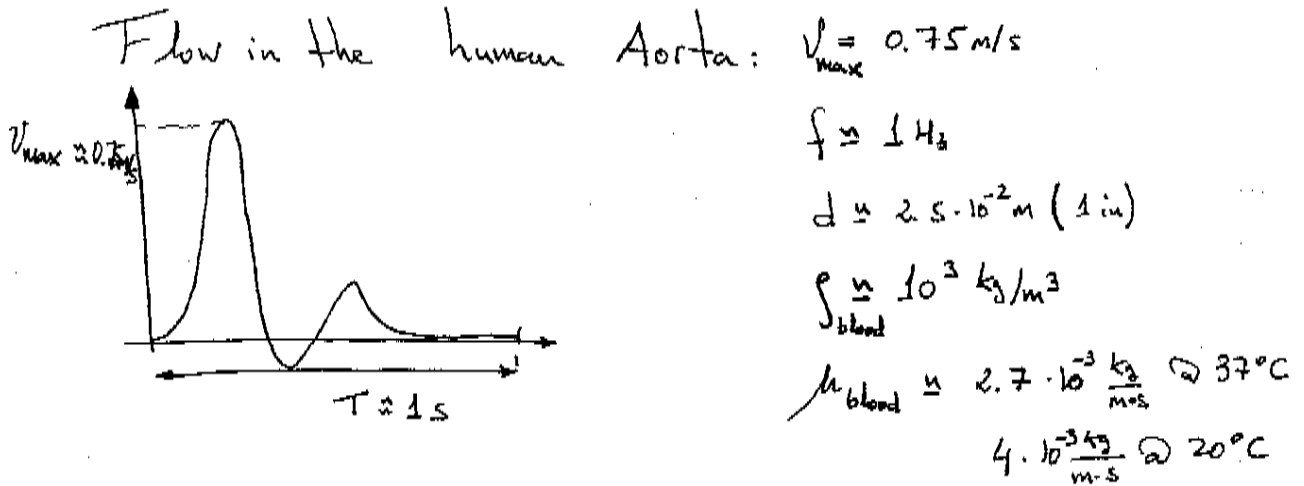
$$C_D = \frac{D}{\frac{1}{2} \rho v_\infty^2 d^2} \quad Re = \frac{\rho v_\infty d}{\mu}$$

↓
Cross Area of the object

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Dynamic Similarity

If we want to test the behaviour of a phenomenon in a laboratory experiment, we need to match the values of the relevant physical parameters between the experiment and the real phenomenon.



Experiment: $d = 10^{-1} \text{ m (4 in)}$

Characteristic non dimensional numbers are Re , St .

$$Z_w = \left\{ \left(d, \mu, \rho, T, v_{max} \right) \quad \left. \begin{array}{l} k=6 \\ r=3 \end{array} \right\} \quad k-r=3$$

$$[Z_w] = \frac{MLT^{-2}}{L^3} = ML^{-1}T^{-2}$$

$$[d] = L \quad \leftarrow$$

$$[f] = T^{-1}$$

$$[v_{max}] = LT^{-1} \quad \leftarrow$$

$$[\rho] = ML^{-3} \quad \leftarrow$$

$$[\mu] = ML^{-1}T^{-1}$$

$$\frac{Z_w}{\rho v_{max}^2} = f \left(\frac{\mu}{\rho v_{max} d}, \frac{f \cdot d}{v_{max}} \right)$$

$$\text{if } \left. \frac{\mu}{\rho v_{max} d} \right|_{\text{real}} = \left. \frac{\mu}{\rho v_{max} d} \right|_{\text{lab}}$$

$$\text{and } \left. \frac{f \cdot d}{v_{max}} \right|_{\text{real}} = \left. \frac{f \cdot d}{v_{max}} \right|_{\text{lab}} \Rightarrow \left. \frac{Z_w}{\rho v_{max}^2} \right|_{\text{real}} = \left. \frac{Z_w}{\rho v_{max}^2} \right|_{\text{lab}}$$

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$$d_{lab} = 4 d_{real} \Rightarrow \left(\frac{f}{v_{max}} \right)_{lab} = \frac{1}{4} \left(\frac{f}{v_{max}} \right)_{real}$$

$$\left(\frac{\mu}{\rho v_{max}} \right)_{lab} = 4 \left(\frac{\mu}{\rho v_{max}} \right)_{real}$$

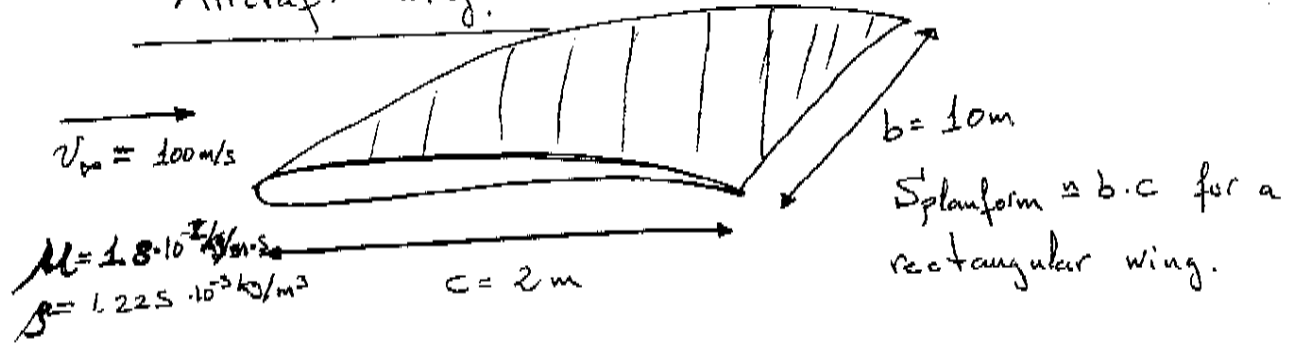
If we choose $f_{lab} = f_{real} \Rightarrow v_{max,lab} = 4 v_{max,real}$ (≈ 3 m/s)

and $\left(\frac{\mu}{\rho} \right)_{lab} = 16 \left(\frac{\mu}{\rho} \right)_{real} \approx \begin{pmatrix} 45 \cdot 10^{-3} \frac{kg}{m \cdot s} \\ (45 \text{ cP}) \end{pmatrix}$
45 times the
viscosity of water

$Z_{w,lab} = 16 Z_{w,real}$ \rightarrow This could be good as Z_w is typically a very small quantity and therefore difficult to measure.

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Aircraft wing.



$$\frac{L}{D} = f(\mu_{air}, v_{\infty}, \rho_{air}, c, b)$$

$$\frac{L}{\frac{1}{2} \rho_{air} v_{\infty}^2 \cdot c \cdot b} = f\left(\frac{\mu_{air}}{\rho_{air} \cdot v_{\infty} \cdot c}, \frac{b}{c}\right)$$

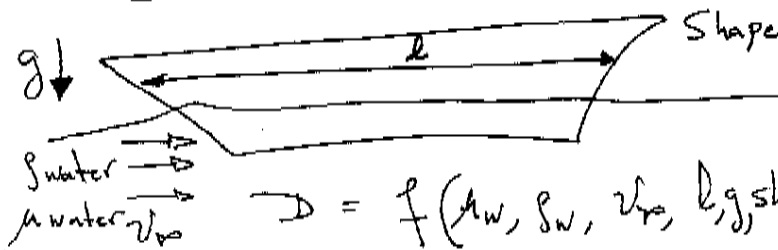
$$c_{lab} = 2 \text{ m} \rightarrow b_{lab} = 10 \text{ m}$$

$$\left(\frac{\mu_{air}}{\rho_{air} \cdot v_{\infty} \cdot c}\right)_{lab} = \left(\frac{\mu_{air}}{\rho_{air} \cdot v_{\infty} \cdot c}\right)_{real} \Rightarrow v_{\infty, lab} = v_{\infty, real} \cdot \frac{c_{real}}{c_{lab}} \cdot \frac{(\frac{\mu_{air}}{\rho_{air}})_{lab}}{(\frac{\mu_{air}}{\rho_{air}})_{real}}$$

$$v_{\infty, lab} \approx 10 \cdot v_{\infty, real} \approx 1000 \text{ m/s.}$$

!!!

Boat



$$D = f(\mu_w, \rho_w, v_{\infty}, l, g, \text{shape})$$

$$\frac{D}{\rho_w v_{\infty}^2 l^2} = f\left(\frac{\mu_w}{\rho_w v_{\infty} l}, \frac{g}{v_{\infty}^2 / l}\right)$$

$$\frac{v_{\infty}}{\sqrt{g \cdot l}} = F_r$$

$$Re_{lab} = Re_{real} ; F_{r, lab} = F_{r, real}$$

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$$\left(\frac{V_{\infty}}{\sqrt{g \cdot l}} \right)_{\text{lab}} = \left(\frac{V_{\infty}}{\sqrt{g \cdot l}} \right)_{\text{real}}$$

if $g_{\text{lab}} = g_{\text{real}}$
and $l_{\text{lab}} = l_{\text{real}}$

$$\left(\frac{V_{\infty} \cdot l}{\nu} \right)_{\text{lab}} = \left(\frac{V_{\infty} \cdot l}{\nu} \right)_{\text{lab}}$$

$$\left(\frac{V_{\infty}}{\sqrt{l}} \right)_{\text{lab}} = \left(\frac{V_{\infty}}{\sqrt{l}} \right)_{\text{real}}$$

$$(V_{\infty} \cdot l)_{\text{lab}} = (V_{\infty} \cdot l)_{\text{real}}$$

You can satisfy one or the other but not both.

$$l_{\text{real}} \approx 100 \text{ m}$$

$$l_{\text{lab}} \approx 1 \text{ m}$$

$$V_{\infty \text{ real}} \approx 10 \text{ m/s}$$

$$V_{\infty \text{ lab}} = V_{\infty \text{ real}} \cdot \sqrt{\frac{l_{\text{lab}}}{l_{\text{real}}}}$$

$$V_{\infty \text{ lab}} = 10 \text{ m/s} \cdot \frac{1}{10} \approx 1 \text{ m/s}$$

but then: $\frac{1 \text{ m/s} \cdot 1 \text{ m}}{\nu_{\text{lab}}} = \frac{10 \text{ m/s} \cdot 100 \text{ m}}{\nu_{\text{real}}} \Rightarrow \frac{\nu_{\text{lab}}}{\nu_{\text{real}}} \approx \frac{1000}{1}$
not feasible.

$$Re_{\text{lab}} \approx 10^6, \quad Re_{\text{real}} \approx 10^9$$

Thanks to the asymptotic behaviour for large Reynolds numbers, the behaviour is reproduced in the lab to some extent.