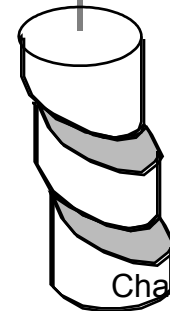
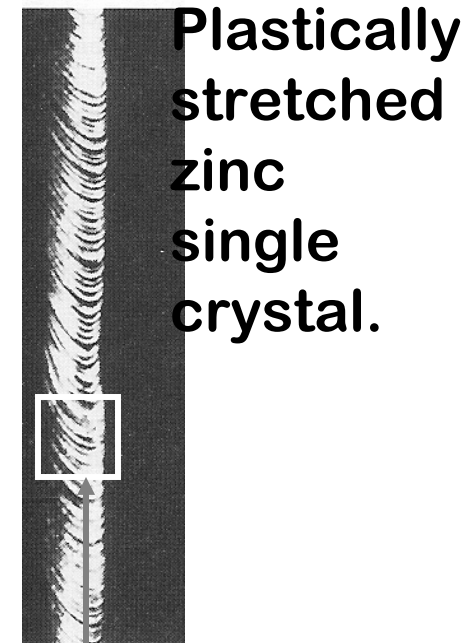


Chapter 7: Dislocations and strengthening mechanisms

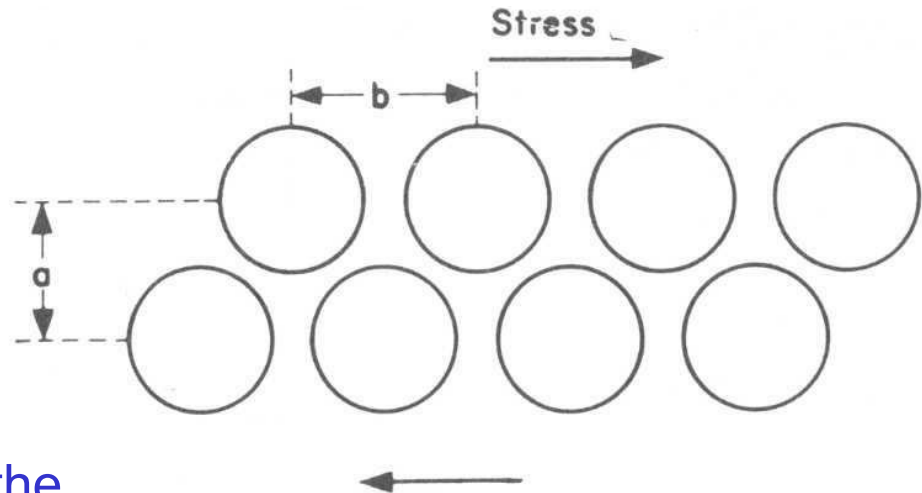
- Introduction
- Basic concepts
- Characteristics of dislocations
- Slip systems
- Slip in single crystals
- Plastic deformation of polycrystalline materials



Theoretical stress

Theoretical stress (Frenkel in 1926)

$$\tau = \frac{Gb}{2\pi a} \sin \frac{2\pi x}{b}$$



G : shear modulus

b : spacing between atoms in the direction of shear stress

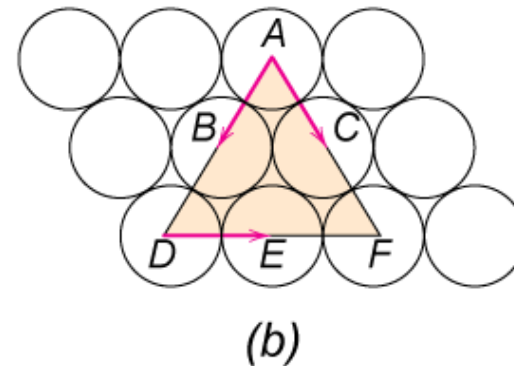
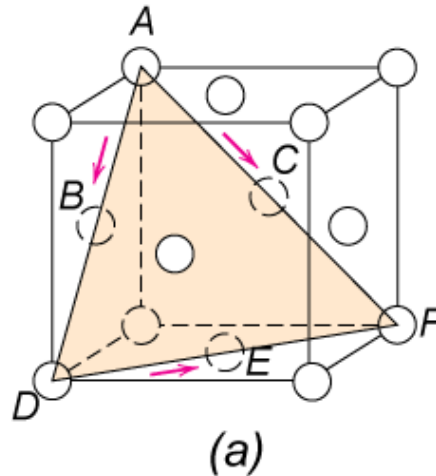
a : spacing of the rows of atoms

x : shear translation

Deformation Mechanisms

Slip System

- Slip plane - plane allowing easiest slippage
 - Wide interplanar spacings - highest planar densities
- Slip direction - direction of movement - Highest linear densities



Adapted from Fig. 7.6, Callister 7e.

- FCC Slip occurs on $\{111\}$ planes (close-packed) in $\langle 110 \rangle$ directions (close-packed)
 - => total of 12 slip systems in FCC
- in BCC & HCP other slip systems occur

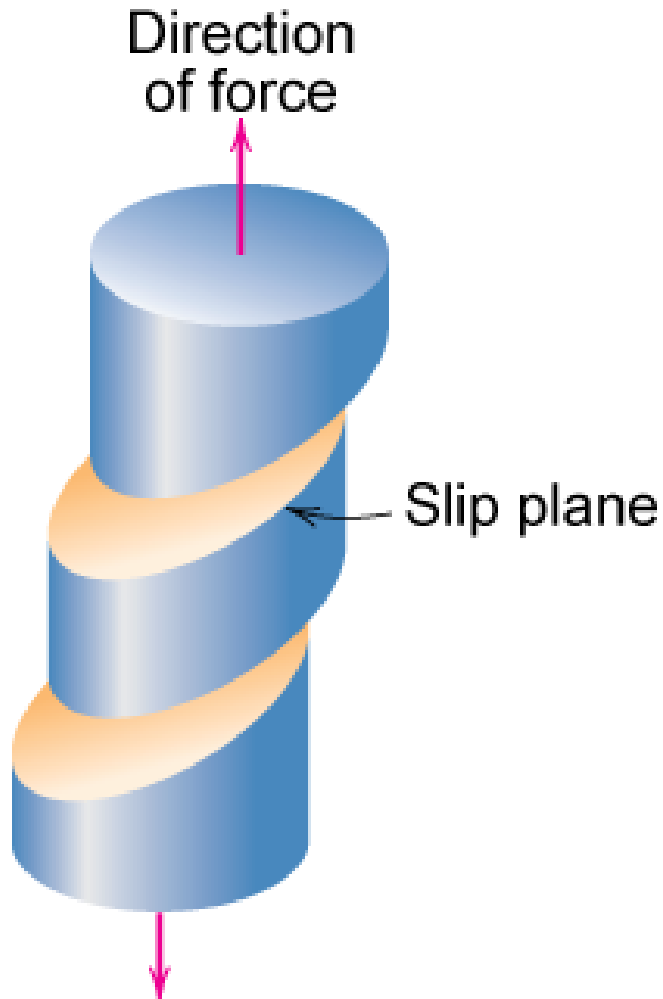
Slip planes and directions for common crystal structure

Table 7.1 Slip Systems for Face-Centered Cubic, Body-Centered Cubic, and Hexagonal Close-Packed Metals

<i>Metals</i>	<i>Slip Plane</i>	<i>Slip Direction</i>	<i>Number of Slip Systems</i>
Face-Centered Cubic			
Cu, Al, Ni, Ag, Au	{111}	$\langle 1\bar{1}0 \rangle$	12
Body-Centered Cubic			
α -Fe, W, Mo	{110}	$\langle \bar{1}11 \rangle$	12
α -Fe, W	{211}	$\langle \bar{1}11 \rangle$	12
α -Fe, K	{321}	$\langle \bar{1}11 \rangle$	24
Hexagonal Close-Packed			
Cd, Zn, Mg, Ti, Be	{0001}	$\langle 11\bar{2}0 \rangle$	3
Ti, Mg, Zr	{10 $\bar{1}0$ }	$\langle 11\bar{2}0 \rangle$	3
Ti, Mg	{10 $\bar{1}1$ }	$\langle 11\bar{2}0 \rangle$	6



Single Crystal Slip



Adapted from Fig. 7.9, *Callister 7e*.

Adapted from Fig. 7.8, *Callister 7e*.



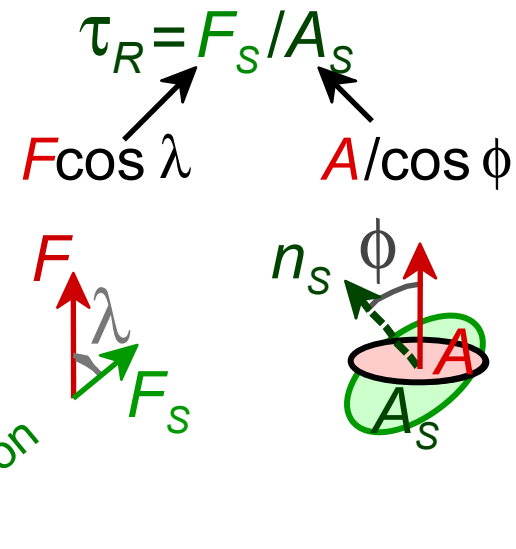
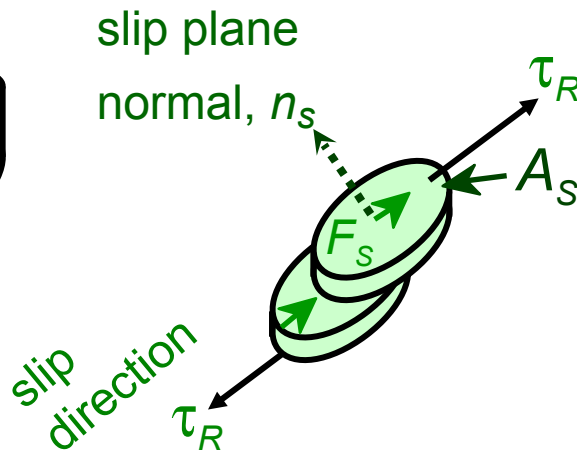
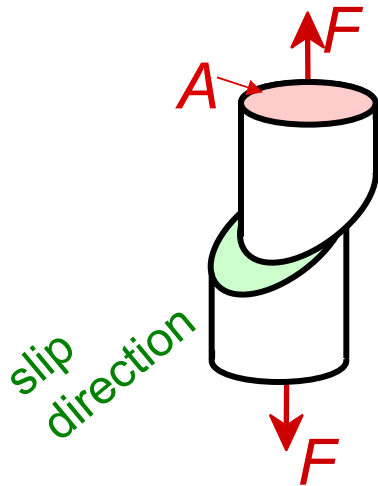
Stress and Dislocation Motion

- Crystals slip due to a **resolved shear stress**, τ_R .
- Applied tension can produce such a stress.

Applied **tensile**
stress: $\sigma = F/A$

Resolved shear
stress: $\tau_R = F_S/A_S$

Relation between
 σ and τ_R

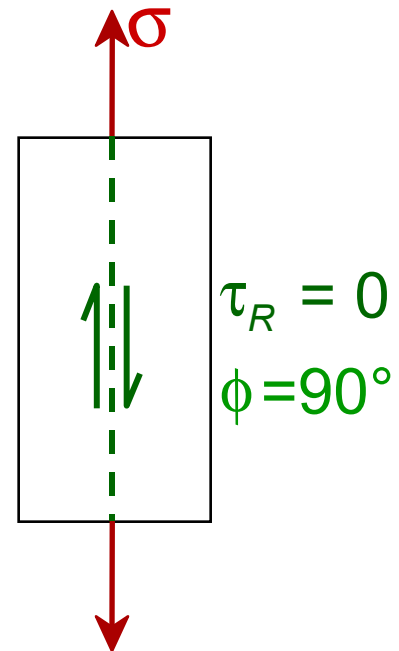
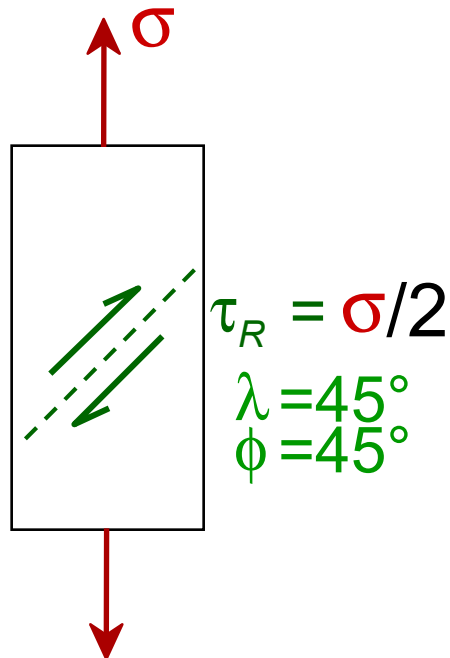
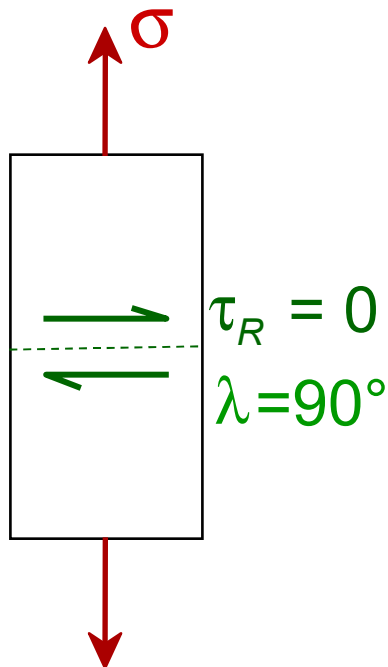


$$\tau_R = \sigma \cos \lambda \cos \phi$$

Critical Resolved Shear Stress

- Condition for dislocation motion: $\tau_R > \tau_{\text{CRSS}}$
 - Crystal orientation can make it easy or hard to move dislocation
- \uparrow
 typically
 $10^{-4} \text{ GPa to } 10^{-2} \text{ GPa}$

$$\tau_R = \sigma \cos \lambda \cos \phi$$



τ maximum at $\lambda = \phi = 45^\circ$

Theoretical & experimental strength

There is much difference between theoretical and experimental strength

Reasons are:

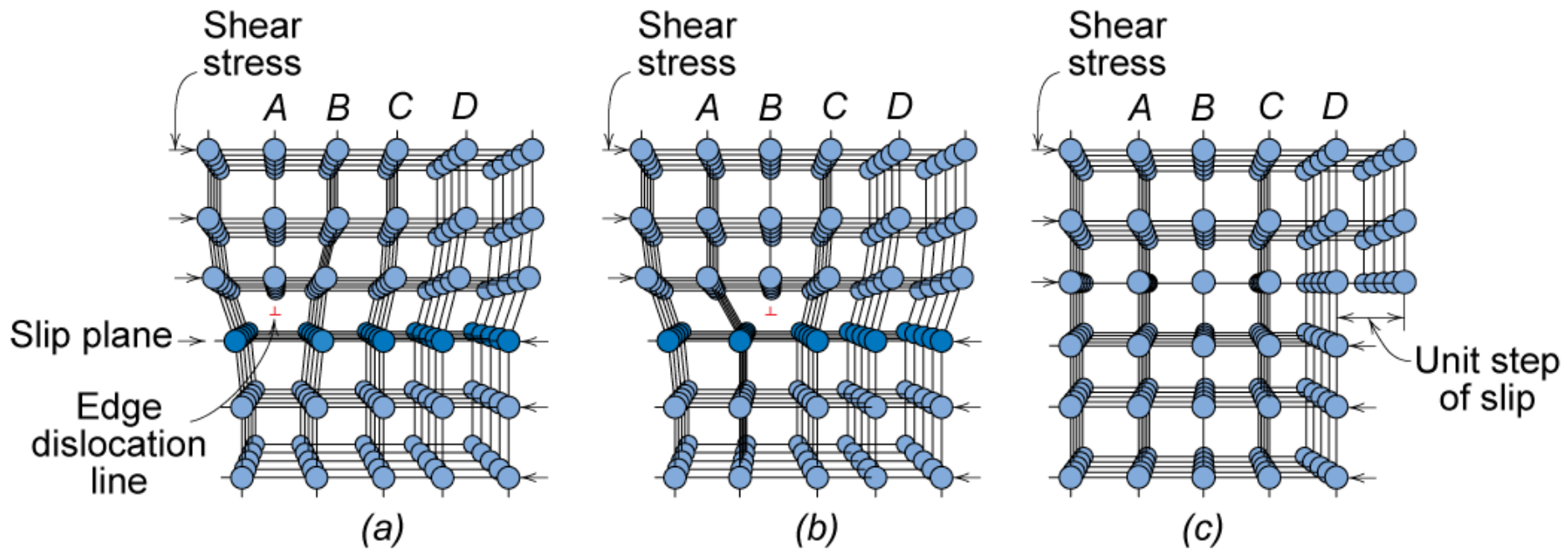
- **Defects are present in all perfect crystal**
- **Dislocation movement makes plastic deformation easier than that predicted by the Frenkel calculation**



Dislocation Motion

Dislocations & plastic deformation

- Cubic & hexagonal metals - plastic deformation by **plastic shear or slip** where one plane of atoms slides over adjacent plane by defect motion (dislocations).



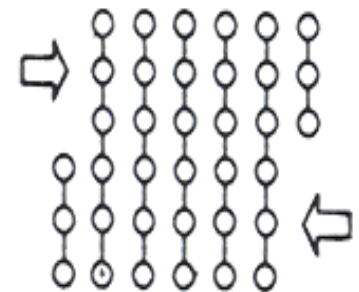
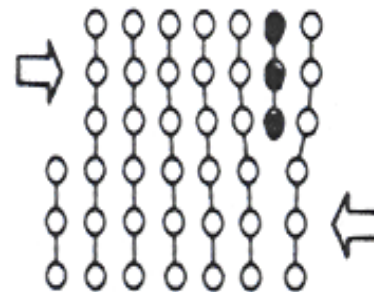
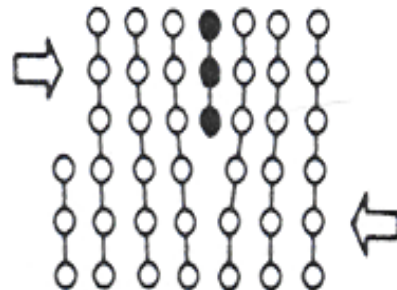
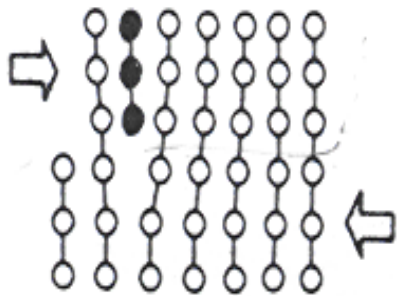
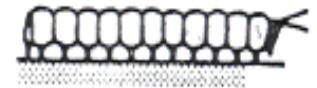
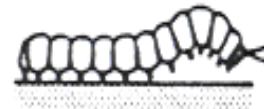
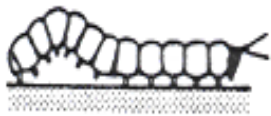
- If dislocations don't move, deformation doesn't occur!

Adapted from Fig. 7.1,
Callister 7e.



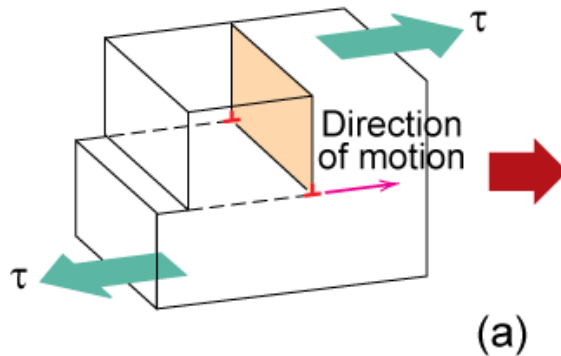
Analogy between caterpillar and dislocation motion

- Dislocation density: total dislocation length per unit volume
- 10^3 mm^{-2} for pure metal crystals; $10^9\text{-}10^{10}\text{mm}^{-2}$ for heavily deformed metals; $10^5\text{-}10^6\text{mm}^{-2}$ for heat-treated deformed metals



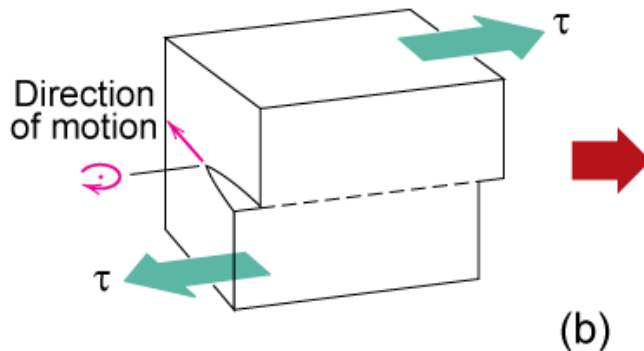
Dislocation Motion

- Dislocation moves along **slip plane** in **slip direction** perpendicular to dislocation line
- Slip direction same direction as **Burgers vector**



Edge dislocation

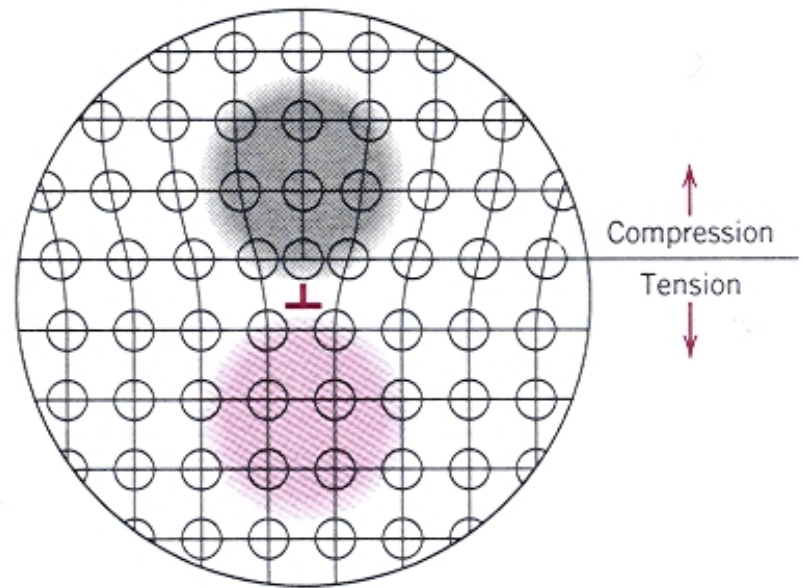
Adapted from Fig. 7.2,
Callister 7e.



Screw dislocation

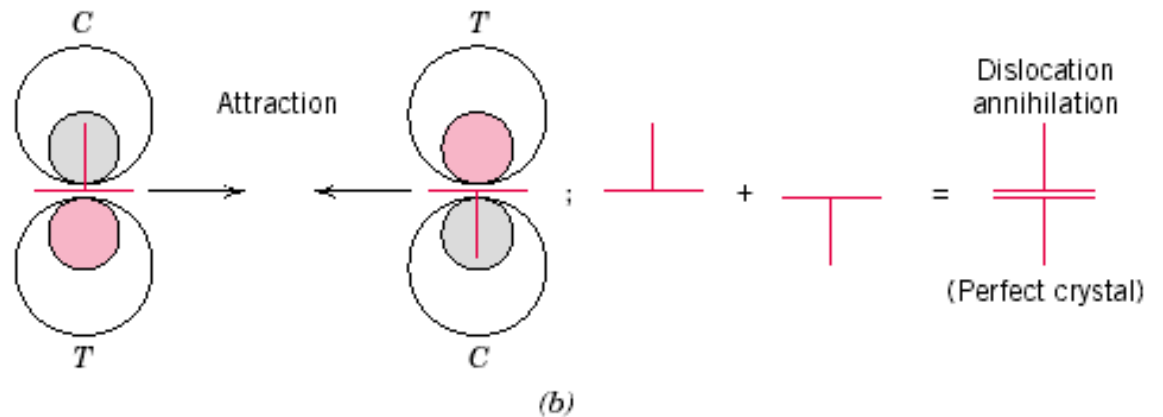
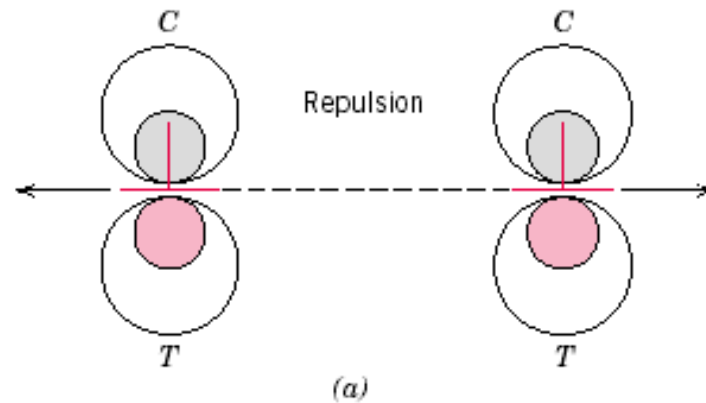
Characteristics of dislocations

- **Strain fields: determining the mobility of the dislocations and their ability to multiply**
- **Compressive, tensile, and shear lattice stains**



Dislocation interaction

- Edge dislocation
- Positive sign
- Negative sign



Slip Motion in Polycrystals

- Stronger - grain boundaries pin deformations
- Slip planes & directions (λ , ϕ) change from one crystal to another.
- τ_R will vary from one crystal to another.
- The crystal with the largest τ_R yields first.
- Other (less favorably oriented) crystals yield later.



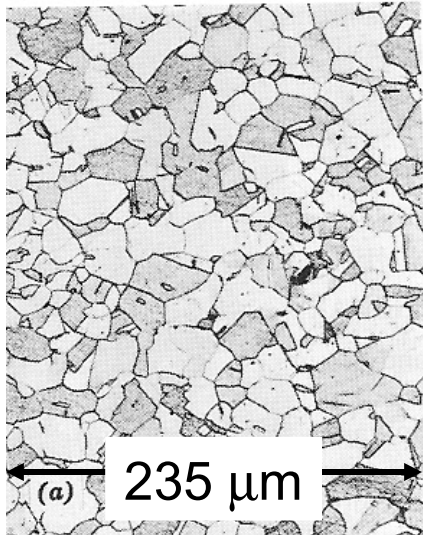
Adapted from Fig. 7.10, *Callister 7e*. (Fig. 7.10 is courtesy of C. Brady, National Bureau of Standards [now the National Institute of Standards and Technology, Gaithersburg, MD].)



Anisotropy in σ_y

- Can be induced by rolling a polycrystalline metal

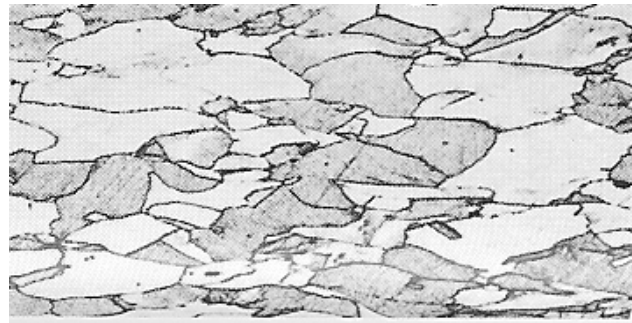
- before rolling



- isotropic

since grains are approx. spherical & randomly oriented.

- after rolling



rolling direction

- anisotropic

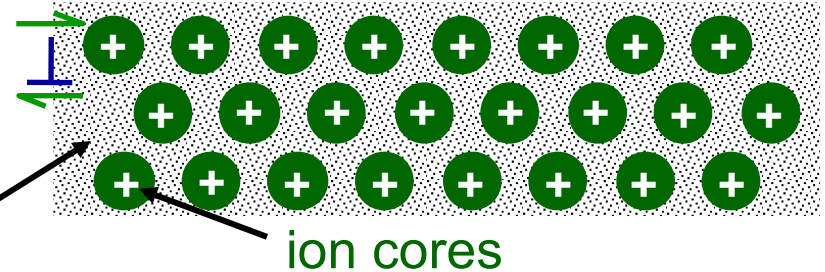
since rolling affects grain orientation and shape.

Adapted from Fig. 7.11, *Callister 7e*. (Fig. 7.11 is from W.G. Moffatt, G.W. Pearsall, and J. Wulff, *The Structure and Properties of Materials*, Vol. I, *Structure*, p. 140, John Wiley and Sons, New York, 1964.)

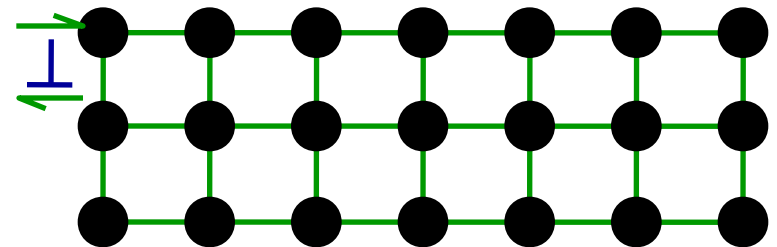
Dislocations & Materials Classes

- Metals: Disl. motion easier.
 - non-directional bonding
 - close-packed directions for slip.

electron cloud



- Covalent Ceramics
(Si, diamond): Motion hard.
 - directional (angular) bonding



- Ionic Ceramics (NaCl):
Motion hard.
 - need to avoid ++ and -- neighbors.

