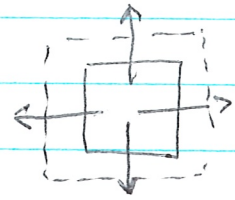


Recap

Deformation in Fluids



- Dilatation (volume changing deformation)

It preserves "shape" but not "volume"

The rate of strain from dilatation process is:

$$\frac{\dot{|\Delta\Gamma|}}{\Delta\Gamma} = \frac{1}{3} (\vec{\nabla} \cdot \vec{v}) \quad \text{Eq. 6.4-14}$$

in matrix form:
$$\underline{\underline{\left(\frac{\dot{|\Delta\Gamma|}}{\Delta\Gamma} \right)}} = \frac{1}{3} (\vec{\nabla} \cdot \vec{v}) \underline{\underline{\delta}}$$

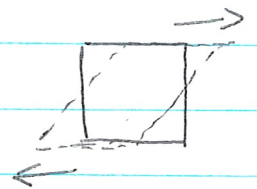
$$= \frac{1}{3} \begin{pmatrix} \vec{\nabla} \cdot \vec{v} & & 0 \\ & \vec{\nabla} \cdot \vec{v} & \\ 0 & & \vec{\nabla} \cdot \vec{v} \end{pmatrix}$$

Shear

- Deformation (shape changing deformation)

It preserves "volume" but not "shape"

Difference between total deformation $\underline{\underline{\Gamma}}$ and dilatation $\underline{\underline{\varepsilon}}_{ii}$



$$\underline{\underline{\Gamma}} - \frac{1}{3} (\vec{\nabla} \cdot \vec{v}) \underline{\underline{\delta}} = \begin{pmatrix} \frac{\partial v_i}{\partial x_1} - \frac{1}{3} (\vec{\nabla} \cdot \vec{v}) & & \\ & \frac{\partial v_2}{\partial x_2} - \frac{1}{3} (\vec{\nabla} \cdot \vec{v}) & \\ & & \frac{\partial v_3}{\partial x_3} - \frac{1}{3} (\vec{\nabla} \cdot \vec{v}) \end{pmatrix}$$

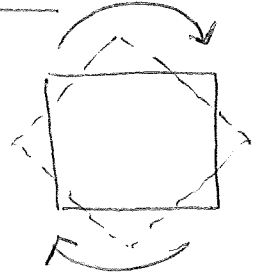
$$\frac{1}{2} \left(\frac{\partial v_i}{\partial x_j} - \frac{\partial v_j}{\partial x_i} \right)$$

Recap cont.

2

Other stress-induced fluid motion

- Rigid Body Rotation



The rigid body rotation is fully expressed by the vorticity tensor

$$\underline{\underline{\Omega}} = \frac{1}{2} \left[\underline{\underline{\nabla V}} - (\underline{\underline{\nabla V}})^T \right] = \frac{1}{2} \begin{pmatrix} 0 & w_z & -w_y \\ -w_z & 0 & w_x \\ w_y & -w_x & 0 \end{pmatrix}$$

i.e., the antisymmetric component of the velocity gradient $\underline{\underline{\nabla V}}$.

w_x, w_y, w_z are the component of the vorticity vector $\vec{w} (w_x, w_y, w_z) = \vec{\nabla} \times \vec{v}$ given

as

$$w_x = \frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z}$$

$$w_y = \frac{\partial v_x}{\partial z} - \frac{\partial v_z}{\partial x}$$

$$w_z = \frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial z}$$

Note:

$$\vec{w} = 2\vec{\omega}$$

$\vec{\omega}$ angular velocity



It turns out that rate of strain of the rigid body rotation is zero.

Recap cont.

3

Constitutive Equation for viscous stress

We search for relationships between stress and strain rate, i.e., $\underline{\underline{\tau}} = f(\underline{\underline{\Gamma}})$

Assume viscous stresses are linearly related to the two strain rate components for (a) dilatation, and (b) shape changing shear deformation

\Rightarrow Newtonian Fluid (Constitutive Eq.)

$$\underline{\underline{\tau}} = 2\mu \left[\underline{\underline{\Gamma}} - \frac{1}{3} (\vec{\nabla} \cdot \vec{v}) \underline{\underline{\delta}} \right] + 3K \left[\frac{1}{3} (\vec{\nabla} \cdot \vec{v}) \underline{\underline{\delta}} \right]$$

Annotations:

- shear viscosity (points to 2μ)
- total deformation (points to $\underline{\underline{\Gamma}}$)
- volume change (points to $\frac{1}{3} (\vec{\nabla} \cdot \vec{v}) \underline{\underline{\delta}}$)
- shape changing shear deformation (bracketed under total deformation and volume change)
- dilatation (bracketed under $\frac{1}{3} (\vec{\nabla} \cdot \vec{v}) \underline{\underline{\delta}}$)
- dilatational viscosity ("bulk viscosity") (points to $3K$)

$$\left. \begin{array}{l} \mu = f(T, P) \\ K = g(T, P) \end{array} \right\} \text{but not dependent on } \underline{\underline{\Gamma}} \text{ or } \vec{\nabla} \cdot \vec{v}$$

As $K \ll \mu$ \Rightarrow $\underline{\underline{\tau}} \approx 2\mu \left[\underline{\underline{\Gamma}} - \frac{1}{3} (\vec{\nabla} \cdot \vec{v}) \underline{\underline{\delta}} \right]$

for $g = \text{const}$ (incompressible fluid) $\Rightarrow \vec{\nabla} \cdot \vec{v} = 0$ (continuity eq.)

\Rightarrow $\underline{\underline{\tau}} \approx 2\mu \underline{\underline{\Gamma}}$ Newtonian fluid for $K \rightarrow 0$ and $g = \text{const}$ (incompressible)

Recap cont.

4

Navier - Stokes Eq. :
Conservation of Momentum for
Newtonian incompressible fluids

Subst. in Cauchy Eq.

$$\rho \frac{D\vec{v}}{Dt} = -\vec{\nabla} P + \rho \vec{g} + \vec{\nabla} \cdot \underline{\underline{\tau}}$$

$$\text{for } \underline{\underline{\tau}} = 2\mu \underline{\underline{\Gamma}} = \mu \left[\vec{\nabla} \vec{\nabla} + \vec{\nabla} \vec{\nabla} \right]^{\leftarrow}$$

$$\Rightarrow \boxed{\rho \frac{D\vec{v}}{Dt} = -\vec{\nabla} P + \rho \vec{g} + \mu \nabla^2 \vec{v}}$$

∇^2 Laplace Operator

Navier - Stokes Eq.