

(found in W.M. Deen, Analysis of Transport Phenomena)

Fluid Properties	Continuity Eq.	Equation of Momentum
general stress vs. strain rate relationship: $\underline{\underline{\tau}}(\underline{\underline{\Gamma}}) = \frac{1}{2}[\vec{\nabla}\vec{v} + (\vec{\nabla}\vec{v})^t]$	$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0 \quad (2.3-1)$	Cauchy's Momentum Eq. $\rho \frac{D\vec{v}}{Dt} = \rho \vec{g} - \vec{\nabla}P + \vec{\nabla} \cdot \underline{\underline{\tau}} \quad (6.3-10)$
assume the following stress vs. strain-rate relationship: $\underline{\underline{\tau}} = 2\mu(\underline{\underline{\Gamma}})\underline{\underline{\Gamma}}$ $\underline{\underline{\Gamma}} = \frac{1}{2} \left(\frac{\partial v_i}{\partial x_j} - \frac{\partial v_j}{\partial x_i} \right)_{ij}$ 9 element symm. tensor		Generalized Momentum Eq. for Newtonian and non-Newtonian Fluids $\rho \frac{D\vec{v}}{Dt} = \rho \vec{g} - \vec{\nabla}P + \vec{\nabla} \cdot 2\mu(\underline{\underline{\Gamma}})\underline{\underline{\Gamma}} \quad (6.5-16)$ magnitude of the strain-rate tensor $\underline{\underline{\Gamma}} \equiv \left[\frac{1}{2} [\underline{\underline{\Gamma}} : \underline{\underline{\Gamma}}] \right]^{1/2} \quad (6.5-12)$ dissipation function $\Phi = (2\underline{\underline{\Gamma}})^2$
Static Fluid: $v = 0; \underline{\underline{\tau}} = 0$		Table 6-10
Fully Developed Creeping Flow ($Re \ll 1$)		Static Pressure Eq. $\vec{\nabla}P = \rho \vec{g} \quad (6.3-11)$
Newtonian: stress vs. strain-rate relationship $\underline{\underline{\tau}} = 2\mu\underline{\underline{\Gamma}}$		Stokes Equation $\vec{\nabla}P + \rho \vec{g} = \mu \nabla^2 \vec{v} \quad (8.2-1)$ dynamic pressure: $\vec{P} = \vec{\nabla}P - \rho \vec{g}$
neglected molecular resistance to isotropic expansion or compression, $\kappa \ll \mu$		Navier-Stokes Representations: $\rho \left(\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} \right) = \rho \vec{g} - \vec{\nabla}P + \mu \nabla^2 \vec{v} \quad (6.5-10)$
assumed $\mu = \text{constant}$, and incompressibility ($\rho = \text{const.}$)	$\vec{\nabla} \cdot \vec{v} = \text{div} \vec{v} = 0$	Vorticity Transport Equation (equivalent to Navier-Stokes) $\frac{D\vec{w}}{Dt} = \vec{w} \cdot \vec{\nabla} \vec{v} + \frac{\mu}{\rho} \nabla^2 \vec{w} \quad (6.8-7)$ for inviscid flow ($\mu = 0$ or $Re \rightarrow \infty$) $\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} + \frac{1}{\rho} \vec{\nabla}P - \vec{g} = 0 \quad (\text{Euler Eq.}) \quad (9.2-2)$ Gromeka-Lamb Equation (equivalent to Euler) $\frac{\partial \vec{v}}{\partial t} + \vec{w} \times \vec{v} = \vec{g} - \vec{\nabla} \left(\frac{P}{\rho} + \frac{v^2}{2} \right) - v(\vec{\nabla} \times \vec{w})$
Planar Flow (equivalent to Navier-Stokes in 2D)		Stream Function Equation $\frac{\partial}{\partial t} (\nabla^2 \psi) - \frac{\partial(\psi, \nabla^2 \psi)}{\partial(x, y)} = v \nabla^4 \psi \quad \text{Table (6-12)}$
		Jacobian determinant $\frac{\partial(f, g)}{\partial(x, y)} = \begin{vmatrix} \frac{\partial f}{\partial x} & \frac{\partial f}{\partial y} \\ \frac{\partial g}{\partial x} & \frac{\partial g}{\partial y} \end{vmatrix}$

cont. Fluid Properties	Continuity Equation	Equation of Momentum
2D Boundary Layer Equation $v_x(x,y) \gg v_y$, $dP/dx = f(x)$ neglect derivatives involving v_y	$\frac{\partial}{\partial x} v_x + \frac{\partial}{\partial y} v_y = 0$	$\rho \left(v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} \right) = -\frac{dP}{dx} + \mu \frac{\partial^2}{\partial y^2} v_x$ ----- Prandtl Equation: $v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} - u \frac{du}{dx} = v \frac{\partial^2}{\partial y^2} v_x$ (9.4-1) ----- Blasius solution for laminar boundary layer flow over a flat plate: Prandtl-Blasius eq. $v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} = v \frac{\partial^2}{\partial y^2} v_x$ (9.2-25) or with dimensionless stream functions (to reduce variables): $\frac{\partial \hat{\psi}}{\partial \hat{y}} \frac{\partial^2 \hat{\psi}}{\partial \hat{x} \partial \hat{y}} - \frac{\partial \hat{\psi}}{\partial \hat{x}} \frac{\partial^2 \hat{\psi}}{\partial \hat{y}^2} = v \frac{\partial^3 \hat{\psi}}{\partial \hat{y}^3}$ (9.4-4)
for known pressure gradient from the outer inviscid flow region (any geometry)		
restriction to planar geometry (pressure gradient vanishes)		
Lubrication Approx.(2D flow) $v_x(x,y) \gg v_y$, $dP/dx = f(x)$; neglect v_y		$\frac{\partial^2}{\partial y^2} v_x = \frac{1}{\mu} \frac{dP}{dx}$
steady state, inviscid, incompressible constant density	$\vec{\nabla} \cdot \vec{v} = \text{div } \vec{v} = 0$	Bernoulli Equation (true along a given streamline) $\frac{v_2^2 - v_1^2}{2} + \int_1^2 \frac{dp}{\rho} + g(h_2 - h_1) = 0$ (9.2-6) Differential Form of Bernoulli equation: Lamb's Equation (steady state and inviscid Gromeka-Lamb Eq.) $\vec{w} \times \vec{v} = -\vec{\nabla} \left(\frac{P}{\rho} + \frac{v^2}{2} + \phi_g \right)$ gravitational potential: $\vec{\nabla} \phi_g \equiv \vec{g}$
irrotational, steady state, inviscid, incompressible, constant density		Bernoulli Equation (true everywhere in the liquid) $\frac{v_2^2 - v_1^2}{2} + \frac{P_2 - P_1}{\rho} + g(h_2 - h_1) = 0$ (9.2-13)