

322 MidTerm 2 Solutions

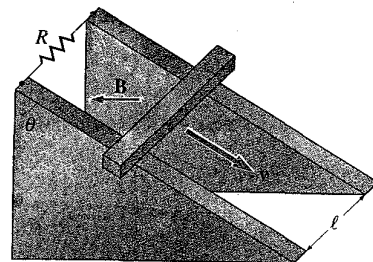
1. (a) In terms of free charge and current densities, ρ_f, \mathbf{J}_f state the partial differential equation version of Maxwell's equations.

$$\begin{aligned} \vec{\nabla} \cdot \vec{D} &= \rho_f, & \vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t}, & \vec{\nabla} \times \vec{H} &= \vec{J}_f + \frac{\partial \vec{D}}{\partial t} \end{aligned}$$

(b) Use Maxwell's equations to show that $\vec{\nabla} \cdot \mathbf{J}_f + \frac{\partial \rho_f}{\partial t} = 0$

$$\begin{aligned} \vec{\nabla} \cdot (\vec{\nabla} \times \vec{H}) &= 0 = \vec{\nabla} \cdot \vec{J}_f + \vec{\nabla} \cdot \frac{\partial \vec{D}}{\partial t} = \vec{\nabla} \cdot \vec{J}_f + \frac{\partial}{\partial t} \vec{\nabla} \cdot \vec{D} \\ &= \vec{\nabla} \cdot \vec{J}_f + \frac{\partial \rho_f}{\partial t} \end{aligned}$$

(c) A conducting bar of mass m , slides down the conducting wedges shown in the figure. The wedges are separated by a distance l , connected at the top by a resistance R and make an angle θ with the vertical. A uniform magnetic field, \mathbf{B} points horizontally as shown. The bar is released from rest and slides down the rails, moving with speed v such that $|\frac{\partial \Phi}{\partial t}| = Blv \cos \theta$. Determine the terminal speed (final constant speed) of the bar.



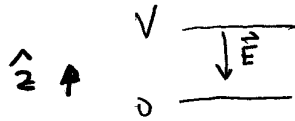
downward force of gravity along wedge
 $= mg \cos \theta$ is balanced by magnetic force $I l B \cos \theta$

\vec{F}_{mg}

$$I R = \left| \frac{\partial \Phi}{\partial t} \right| = Blv \cos \theta$$

$$\text{so } mg \cos \theta = \frac{(Bl)^2 v \cos \theta}{R}$$

$$v = \frac{mg R}{(Bl)^2 \cos \theta}$$



2. A capacitor with parallel plates, with radius a and separation $d \ll a$ has a potential difference $V(t) = V_0 \frac{t}{\tau_0}$, which is the potential at the top ($z = d$) minus that of the bottom ($z = 0$), and τ_0 is a constant with dimensions of time.

(a) Show that the magnetic field \mathbf{B} obeys the equation: $\nabla \times \mathbf{B}(s, z, t) = -\frac{\mu_0 \epsilon_0}{\tau_0 d} V_0 \hat{\mathbf{k}}$, where s is the distance from the axis of symmetry.

$$\vec{E} = -\frac{V(t)}{d} \hat{\mathbf{k}}$$

$$\vec{\nabla} \times \vec{B} = + \frac{\partial \vec{E}}{\partial t} \mu_0 \epsilon_0$$

$$= \frac{\mu_0 \epsilon_0}{d} \left(+\frac{V_0}{\tau_0} \right) \hat{\mathbf{k}}$$

(b) Determine \mathbf{B} for $0 < z < d$ and $s < a$. \vec{B} is along $\hat{\phi}$ direction

$$\vec{B} = B \hat{\phi}$$

$$\oint_C \vec{B} \cdot d\vec{\ell} = \epsilon_0 \frac{\mu_0 V_0}{d \tau_0} \text{ area}$$

$$2\pi B s = \frac{\epsilon_0 \mu_0 V_0}{d \tau_0} \pi s^2 \Rightarrow$$

$$\vec{B} = \frac{\epsilon_0 \mu_0 V_0}{d \tau_0} \frac{s}{2} \hat{\phi}$$

(c) Determine \mathbf{B} for $0 < z < d$ and $s > a$.

\vec{B} is along $\hat{\phi}$ direction

$$\vec{B} = B \hat{\phi}$$

$$\oint_C \vec{B} \cdot d\vec{\ell} = \frac{\epsilon_0 \mu_0 V_0}{d \tau_0} \text{ area}$$

$$B = \frac{1}{2\pi s} \frac{\epsilon_0 \mu_0 V_0}{d \tau_0} \pi a^2$$

$$\vec{B} = \frac{\mu_0 \epsilon_0 V_0}{\tau_0 d} \frac{a^2}{2s} \hat{\phi}$$

3. An infinitely long solenoid of radius R , n turns per unit length, and current I_0 lies along the z axis, which also is the axis of the solenoid. $\mathbf{B} = \mu_0 n I \hat{\mathbf{k}}$. Let s be the distance away from the z axis.

(a) Determine the magnetic energy density for $s < R$.

$$u = \frac{1}{2\mu_0} B^2 = \frac{\mu_0 n^2 I^2}{2}$$

(b) Now suppose that the current is time-dependent, so that, if $t > 0$, it decays slowly with time: $I(t) = I_0 e^{-t/\tau_0}$. Determine $\mathbf{E}(s, t)$ for positions inside the solenoid, where s is the distance away from the center. $\vec{E} = E \hat{\phi}$

$$\oint_C \vec{E} \cdot d\vec{l} = -\frac{\partial \Phi}{\partial t} \quad \text{C is circle of radius } s$$

$$2\pi s E = -\pi s^2 \mu_0 n \frac{\partial I}{\partial t} = \pi s^2 \frac{\mu_0 n I_0}{\tau_0} e^{-t/\tau_0}$$

$$\vec{E} = \frac{s}{2} \frac{\mu_0 n I_0}{\tau_0} e^{-t/\tau_0} \hat{\phi}$$

(c) Determine the density of electromagnetic momentum and the vector $\nabla \cdot \vec{T}$.

$$\vec{P}_{em} = \epsilon_0 (\vec{E} \times \vec{B}) = \epsilon_0 \frac{(\mu_0 n)^2 I^2}{\tau_0} \frac{s}{2} (\hat{\phi} \times \hat{z})$$

$$= \epsilon_0 \frac{(\mu_0 n)^2 I^2}{\tau_0} \frac{s}{2} \hat{r}$$

$$\vec{\nabla} \cdot \vec{T} = \frac{\partial \vec{P}_{em}}{\partial t} = -\epsilon_0 \frac{(\mu_0 n)^2}{\tau_0^2} s I_0^2 e^{-2t/\tau_0} \hat{r}$$