You have 60 minutes. End on the buzzer at 9:20. Answer all questions.
You are strongly recommended to read through all the questions before you begin.
Write your name on every page and your ID on the first page.
Watch the blackboard for corrections or clarifications during the exam.
In this exam you are allowed no books, no note sheets, and no calculators!

Write all your working on these question sheets. You will get credit for it. It is important to show your calculation or derivation, and where appropriate to write a few words to indicate the reasoning. You usually won’t get full marks just for stating the correct answer.

Do not turn this page until the buzzer goes.
1. A straight, neutral wire along the z-axis carries a current \( I(t) = \begin{cases} 0 & t < 0 \\ I_0 & t \geq 0 \end{cases} \) in the +\( \hat{z} \) direction.

1. [8] State the relations for getting the fields \( \mathbf{B} \) and \( \mathbf{E} \) from the potentials \( \psi \) and \( \mathbf{A} \) in electrodynamics.

\[
\mathbf{E} = -\nabla \psi - \frac{\partial \mathbf{A}}{\partial t} \quad \mathbf{B} = \mathbf{\nabla} \times \mathbf{A}
\]

2. [10] Show that there exists a gauge transformation parameterized by a scalar field \( \xi(\mathbf{r}) \) such that neither \( \mathbf{E} \) nor \( \mathbf{B} \) is affected by choice of gauge.

Consider \( \mathbf{A} \rightarrow \mathbf{A}' = \mathbf{A} + \nabla \xi \quad \therefore \quad \mathbf{E}' = \mathbf{\nabla} \times \mathbf{A}' = \mathbf{\nabla} \times (\mathbf{A} + \nabla \xi) = \mathbf{\nabla} \times \mathbf{A} = \mathbf{E} \)

\[
\mathbf{V}' = \mathbf{V} - \frac{\partial \xi}{\partial t}
\]

\[
\mathbf{E}' = -\mathbf{\nabla} \mathbf{V}' - \frac{\partial \mathbf{A}'}{\partial t} = -\nabla(V - \frac{\partial \xi}{\partial t}) - \frac{\partial}{\partial t} \left( \mathbf{A} + \nabla \xi \right) = -\nabla \mathbf{V} - \frac{\partial \mathbf{A}}{\partial t} + \nabla \frac{\partial \xi}{\partial t} \frac{\partial \xi}{\partial t}
\]

3. [5] State the Lorentz gauge condition. What is the key physical advantage of the Lorentz gauge over the Coulomb gauge?

\[
\mathbf{\nabla} \cdot \mathbf{A} = -\frac{1}{c^2} \frac{\partial \mathbf{V}}{\partial t}
\]

4. [16] Find the scalar potential \( \psi(\mathbf{r}, t) \) and vector potential \( \mathbf{A}(\mathbf{r}) \) for this wire in the Lorentz gauge. Your answer for \( \mathbf{A} \) should include a fully explicit integral over the distance \( z' \) along the \( z \) axis. Do not attempt to evaluate the integral.

\[
\psi(\mathbf{r}, t) = \int \frac{\rho(\mathbf{r}', t)}{4\pi \varepsilon_0 |\mathbf{r} - \mathbf{r}'|} \, d^3 r' \\
\psi = 0 \quad \text{because} \quad \rho = 0 \quad \text{at all} \ t.
\]

\[
\mathbf{A}(\mathbf{r}, t) = \int \frac{\mathbf{\mu}_0 \mathbf{I}(\mathbf{r}', t)}{4\pi |\mathbf{r} - \mathbf{r}'|} \, d^3 r' \\
= \int_0^{t_c} \frac{\mathbf{I}_0 t_c}{4\pi |\mathbf{r} - \mathbf{r}'|} \, d^3 r' \\
= \frac{\mu_0 I_0}{2\pi} \hat{z} \int \frac{I_0 t_c \, dz'}{\sqrt{r^2 + z'^2}} \\
= \frac{\mu_0 I_0}{2\pi} \hat{z} \left[ \int_0^{\sqrt{r^2 + z'^2}} \frac{dz'}{(\sqrt{r^2 + z'^2})^2} \right] \\
= \frac{\mu_0 I_0}{2\pi} \hat{z} \left[ \int_0^{\sqrt{r^2 + z'^2}} dz' \left( \frac{\sqrt{r^2 + z'^2}}{z'} - \frac{z'}{\sqrt{r^2 + z'^2}} \right) \right] \\
= \frac{\mu_0 I_0}{2\pi} \hat{z} \left[ \int_0^{\sqrt{r^2 + z'^2}} \frac{dz'}{z'} - \frac{z'}{\sqrt{r^2 + z'^2}} \right] \\
= \frac{\mu_0 I_0}{2\pi} \hat{z} \left[ \ln \left( \frac{\sqrt{r^2 + z'^2}}{z'} \right) - \frac{z'}{\sqrt{r^2 + z'^2}} \right] \\
= \frac{\mu_0 I_0}{2\pi} \hat{z} \left[ \ln \left( \frac{\sqrt{r^2 + z'^2}}{z'} \right) \right] \\
= \frac{\mu_0 I_0}{2\pi} \hat{z} \left[ \ln \left( \frac{\sqrt{r^2 + z'^2}}{z'} \right) \right]
\]
5. Sketch the variation of the magnitude of \( A \) with time \( t \) at a distance \( r \) from the wire.

\[ |\vec{A}| \text{ keeps increasing as does } I(t) \]

6. Would the result of the previous question be different in the Coulomb gauge?

Yes. (Change in \( \vec{A} \) would be instantaneous.)

7. Is power radiated during this process?

Yes. (\( \vec{B} = \vec{V} \times \vec{A} \) is \( \perp \hat{z} \), \( \vec{E} \sim \int \vec{V} \times \vec{B} \, dt \) is \( \parallel \hat{z} \), so \( \vec{S} \) is outward.)

8. A point charge \( q \) moves along the \( z \)-axis at constant velocity \( u \), reaching the origin at time \( t = 0 \).

\[ \vec{w}(t) = ut \hat{z}. \]

9. Sketch the \( \vec{E} \) and \( \vec{B} \) fields of the point charge if \( u \) is close to \( c \) (say, \( u = 0.8c \)).

\[ t = \frac{t - \frac{t_c}{c}}{c} \text{ when } u = 0 \]
10. [10] When the charge reaches the origin it suddenly stops. Sketch the lines of electric field at a later time \( t \), showing the effect of this sharp deceleration. Indicate the features in the field lines which tell you that radiation is emitted, and in what direction the power is concentrated, and point out as many other features as you can.

At a later time, the charge is made to oscillate about the origin, with position given by \( z = d \cos \omega t \).

11. [9] What are the conditions under which the electric field at a point \( \mathbf{r} \) due to this oscillating charge is given by the following equation? 

\[
\mathbf{E} = \frac{\mu_0 \omega^2 d \cos \left[ \omega (t - r/c) \right]}{4\pi r} \mathbf{\hat{r}} \times (\mathbf{\hat{r}} \times \mathbf{\hat{z}})
\]

\( \delta \ll r \) dipole contribution dominates multipole expansion

\( \delta \ll \lambda \) "long wavelength" \( \lambda = \frac{\omega}{2\pi c} \) phase varies little over source

\( r \gg \lambda = \frac{\omega}{2\pi c} \) "radiation zone", radiative terms \( \gg \) static terms.

12. [10] Under these conditions, draw a sketch to indicate the lines of \( \mathbf{E} \) and \( \mathbf{B} \) at a fixed time \( t \) over a large sphere centered at the origin. What is the direction and angular dependence of the radiative energy flow through the sphere?

\( \mathbf{\hat{E}} = -\mathbf{\hat{z}} \cos \left[ \omega (t - \frac{r}{c}) \right] \)

\( \mathbf{\hat{B}} \sim \mathbf{\hat{E}} \sim \mathbf{\hat{z}} \)

This is electric dipole radiation.

\( S \propto \mathbf{E} \mathbf{B} \propto \sin^2 \theta \)

\( \max \text{ at } \theta = \frac{\pi}{2}, \text{ equator} \)

\( \text{zero at poles} \).