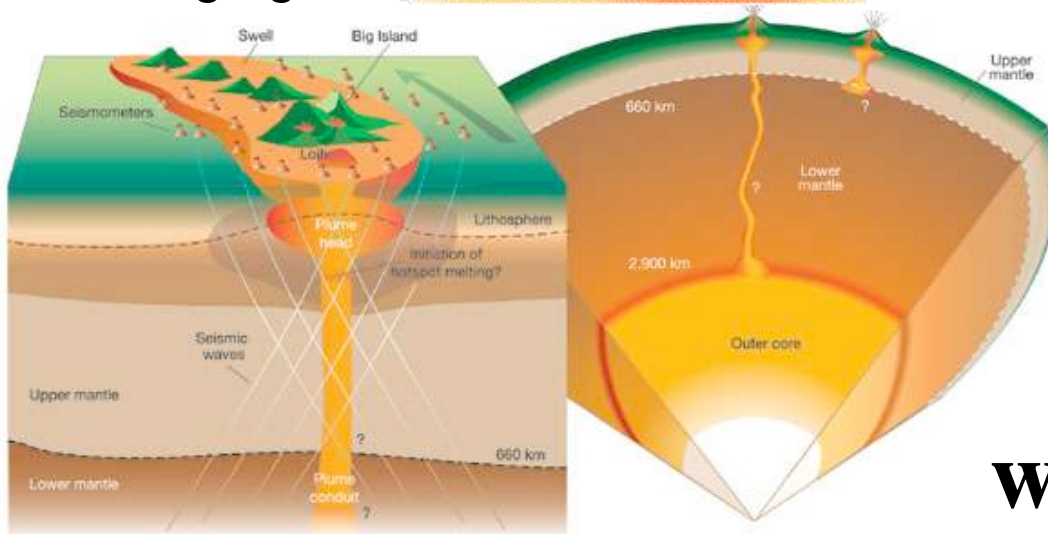
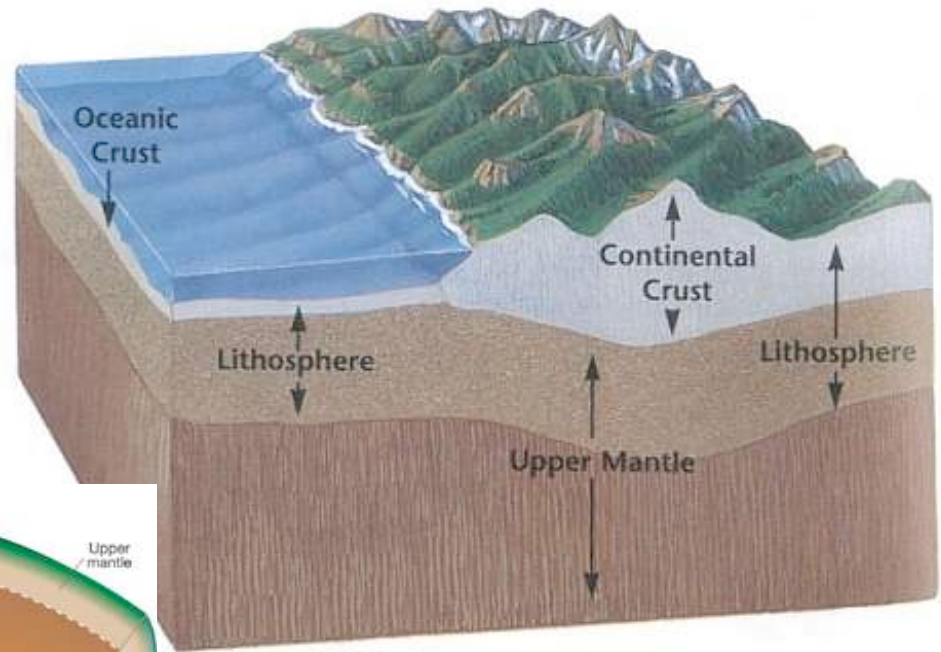
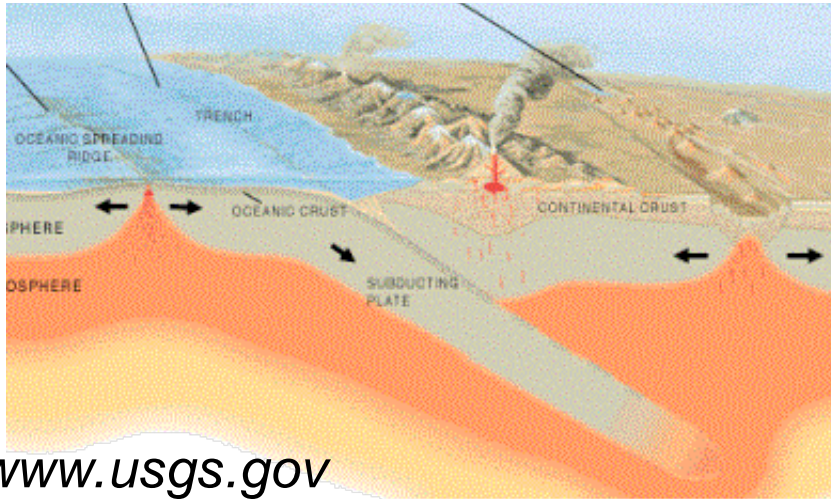


# Seismic tomography: Art or science?

Frederik J Simons  
*Princeton University*

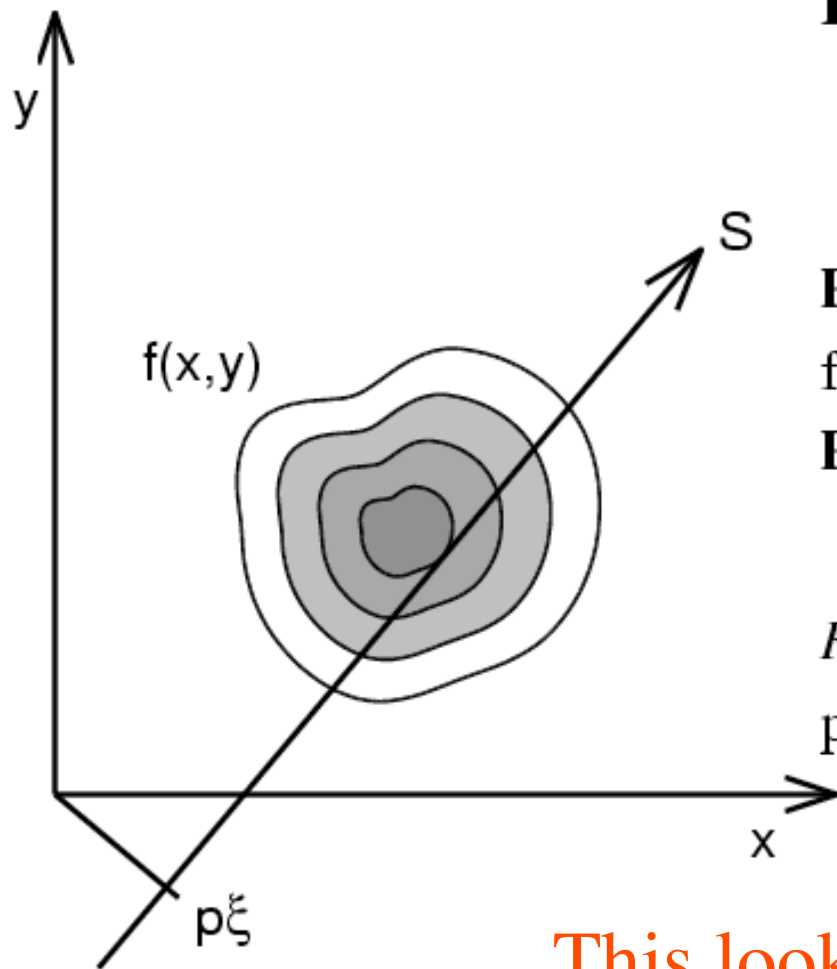


# What's inside the Earth?



**Only seismic  
waves have actually  
*been there, done that***

# The seismic tomography problem



## Inverting the Radon transform

$$\mathcal{R}[f(x, y)](p, \xi) = \int_S f(x, y) ds \quad (1)$$

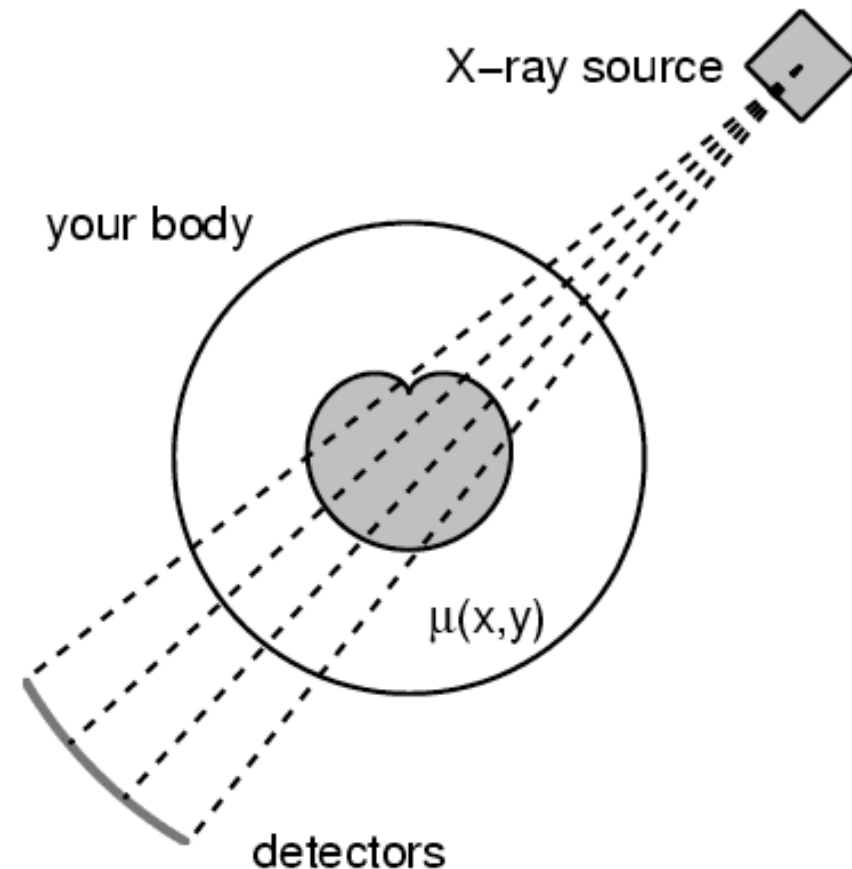
**Purpose:** Reconstruction of functions from their line integrals (projections).

**Problem:** Given  $\mathcal{R}[f(x, y)](p, \xi)$ , find  $f(x, y)$ .

*Radon* [1917] derived a solution to this problem, giving an expression for  $\mathcal{R}^{-1}$ .

This looks more complicated than it is;  
and that's my point.

# What is $f(x, y)$ ? Medical applications.



## X-ray absorption & scattering

Tissues and bones have  $\neq$  absorption and scattering coefficients  $\mu(x, y)$ .

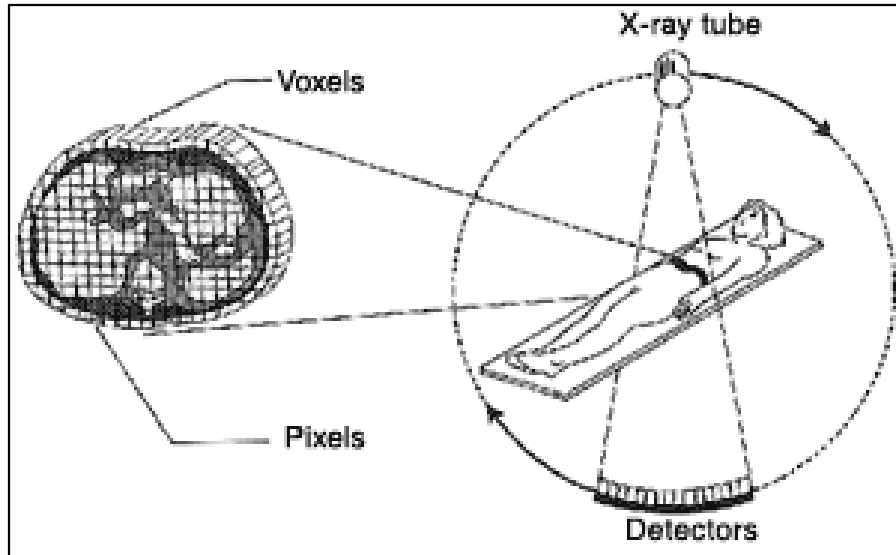
Recorded intensity goes as

$$I = I_0 \exp \left[ \int_{\text{ray}} -\mu(x, y) ds \right]. \quad (2)$$

Sources and detectors rotate to achieve perfect “coverage”.

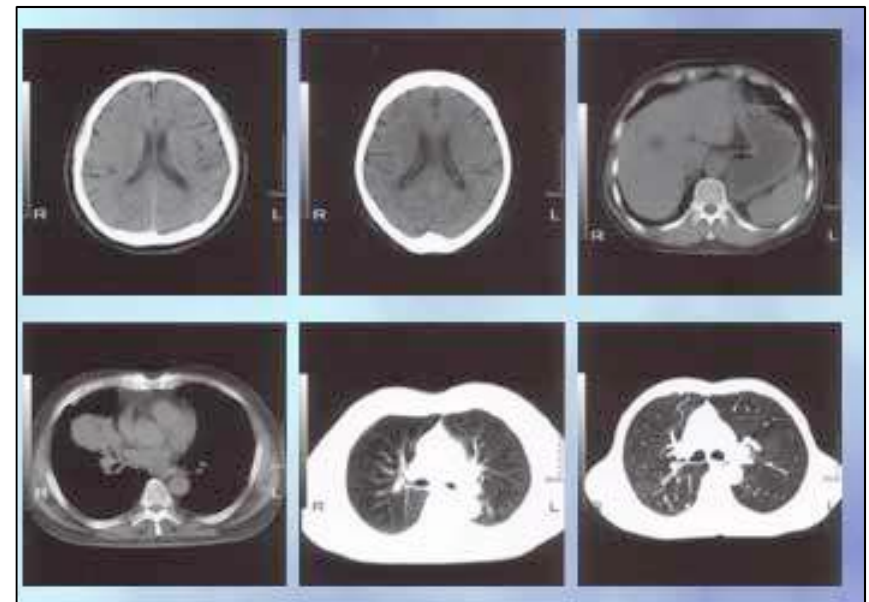
This looks simpler than it is;  
and that's my point.

# X-Ray attenuation tomography

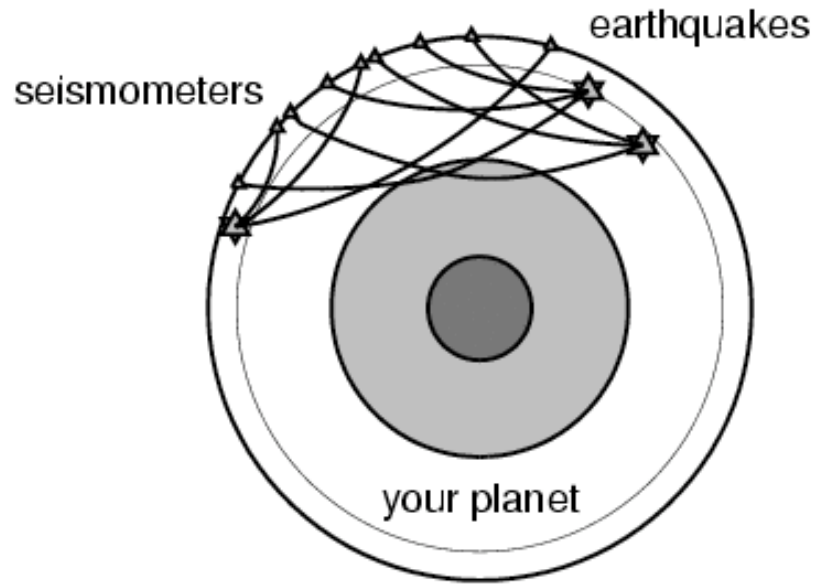


Projections from all angles:  
*X-ray intensity*

Reconstructed image:  
*X-ray attenuation constants*



# What is $f(x, y, z)$ ? Seismic wavespeeds.



## Travel-time tomography

The Earth is made of a heterogeneity of seismic velocities  $v(x, y, z)$ .

Travel-time anomalies go as

$$\delta t = \int_{\text{ray}} \frac{1}{\delta v(x, y, z)} ds. \quad (3)$$

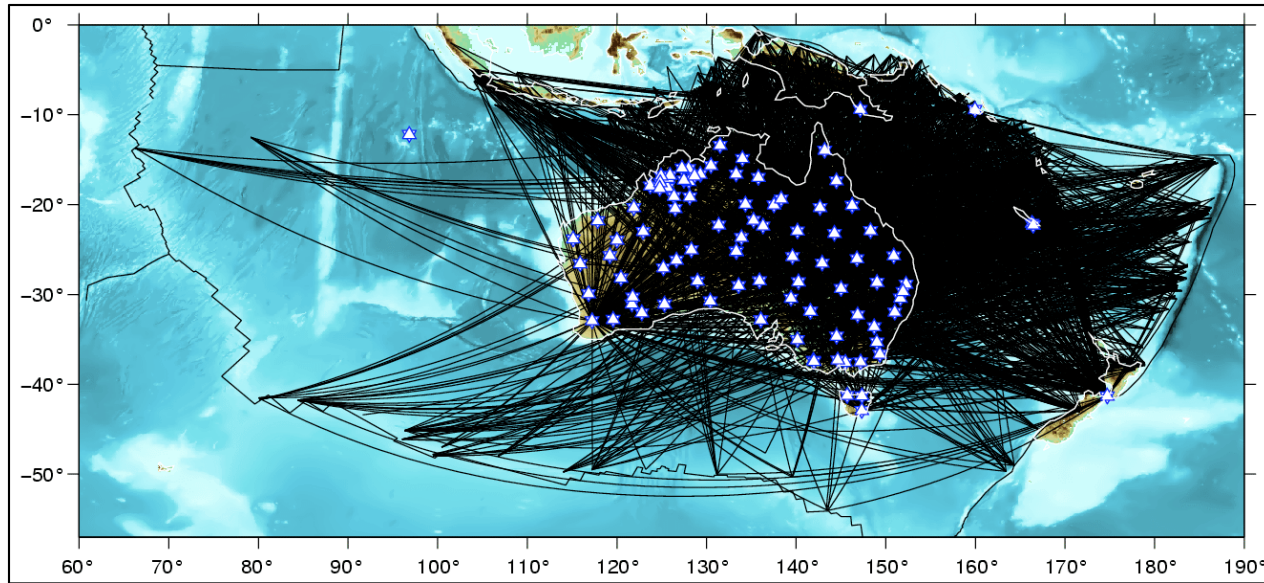
## Waveform tomography

Arrival times depend on the wavelength of the seismic phases.

All raypaths curve and coverage is far from perfect.

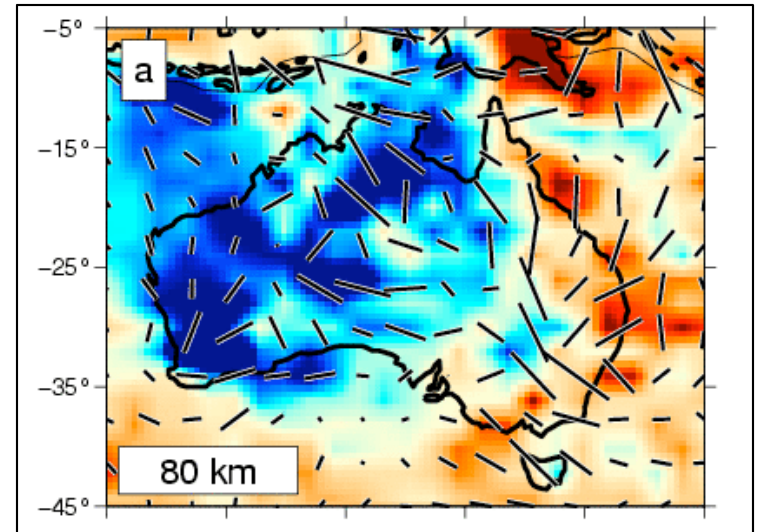


# Seismic wavespeed tomography



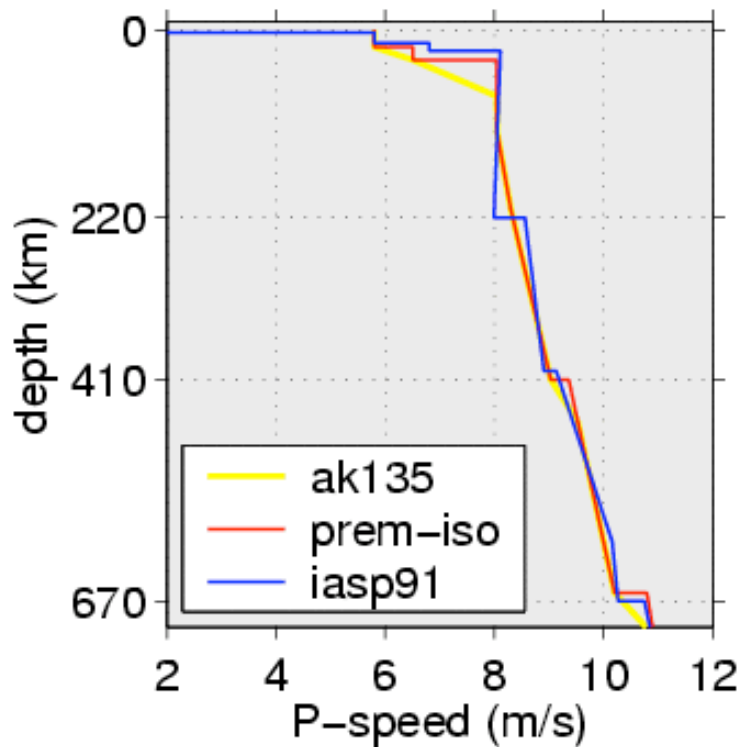
Projections from all angles:  
*Waveforms and arrival times*

Reconstructed image:  
*Wavespeed variations*

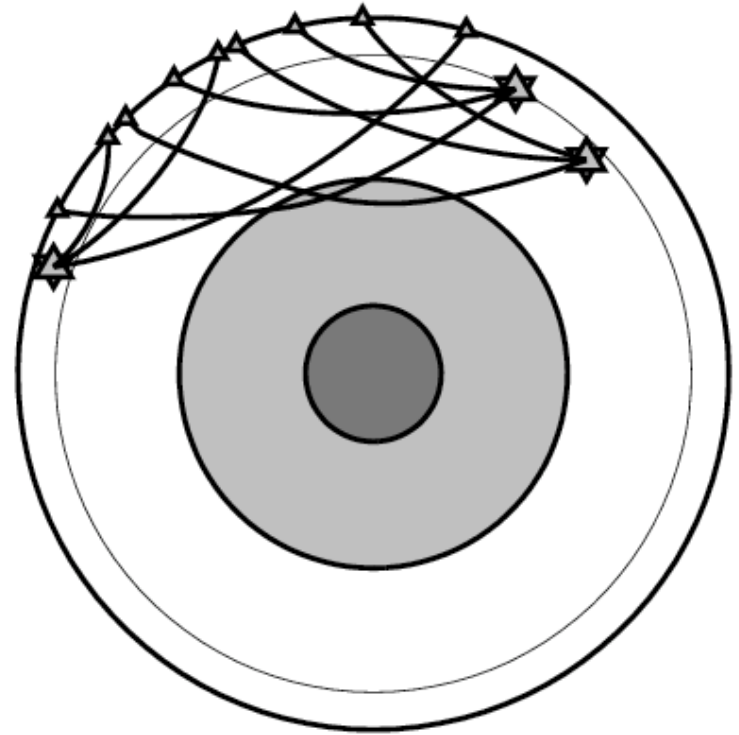


# Forward modeling of the wave field, Part I:

## Ray tracing, **most 1-D**



Before

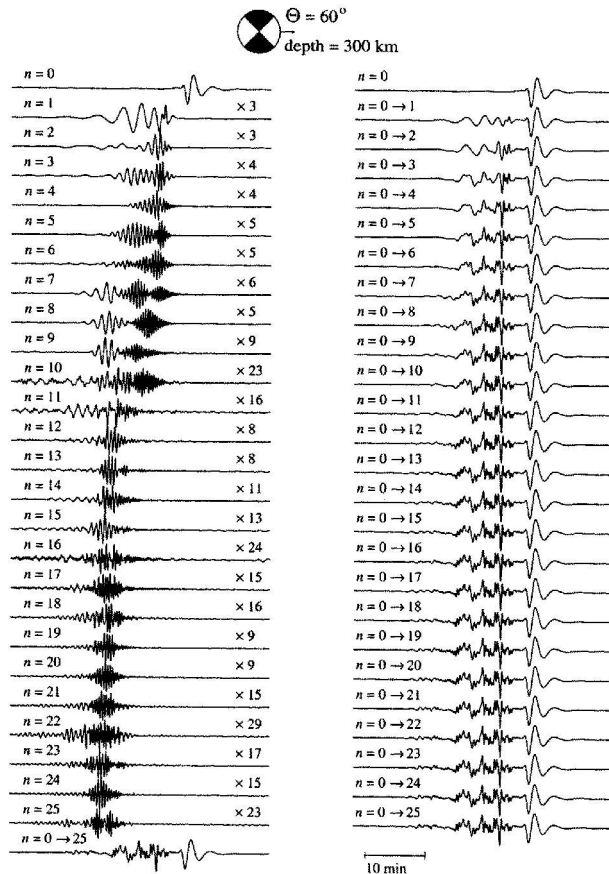


After

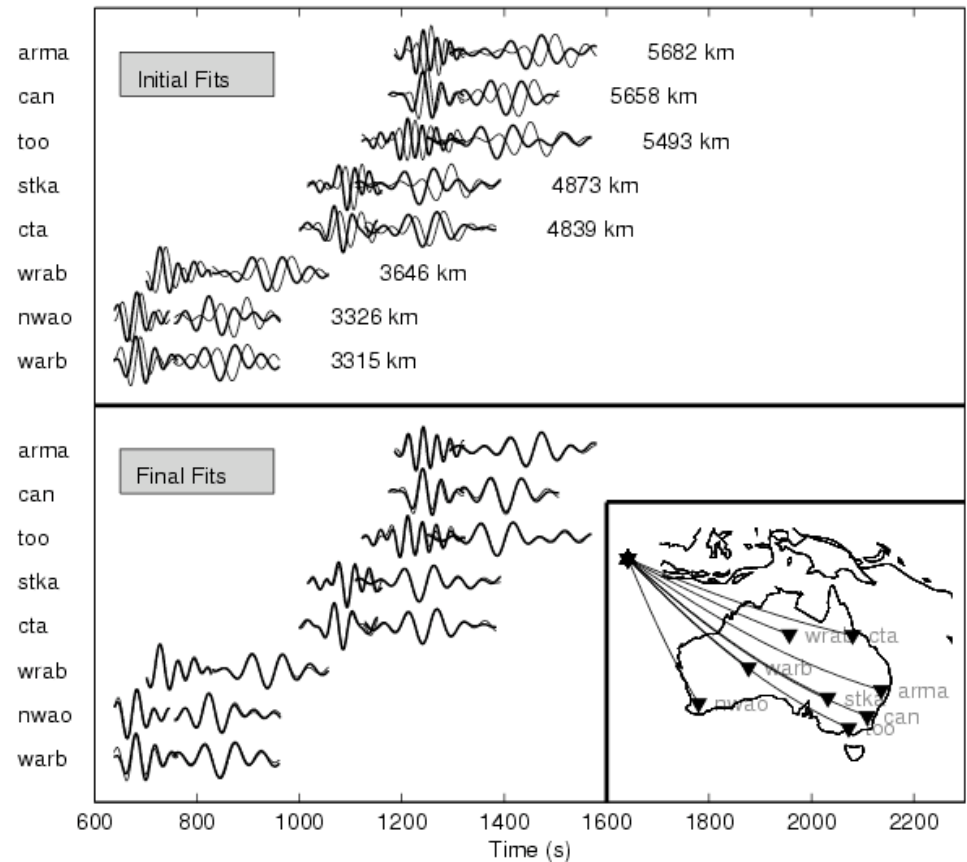


# Forward modeling of the wave field, Part II:

## Normal-mode summation, 1-D



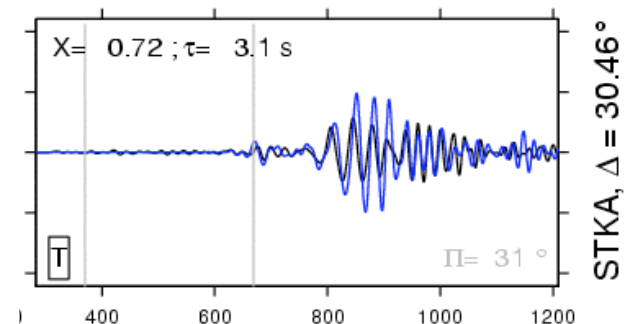
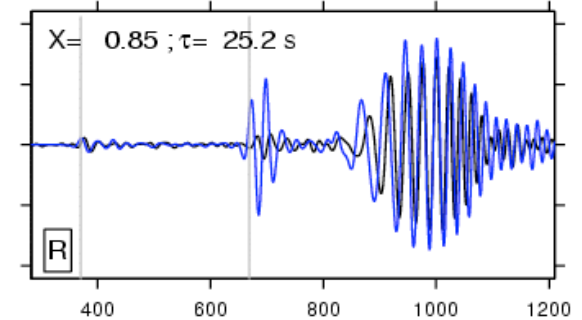
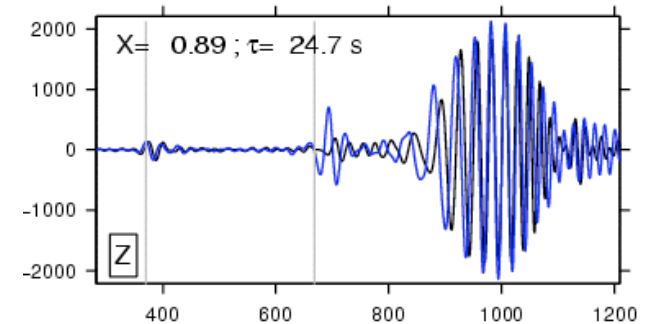
Before



After

# Forward modeling of the wave field, Part III: Spectral-element methods, 3-D

Before



After

# That's all there is to it. Gooby!

## Except:

---

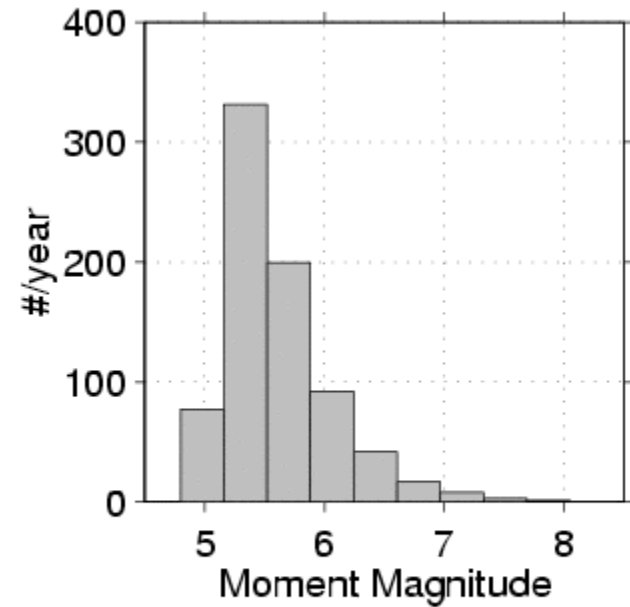
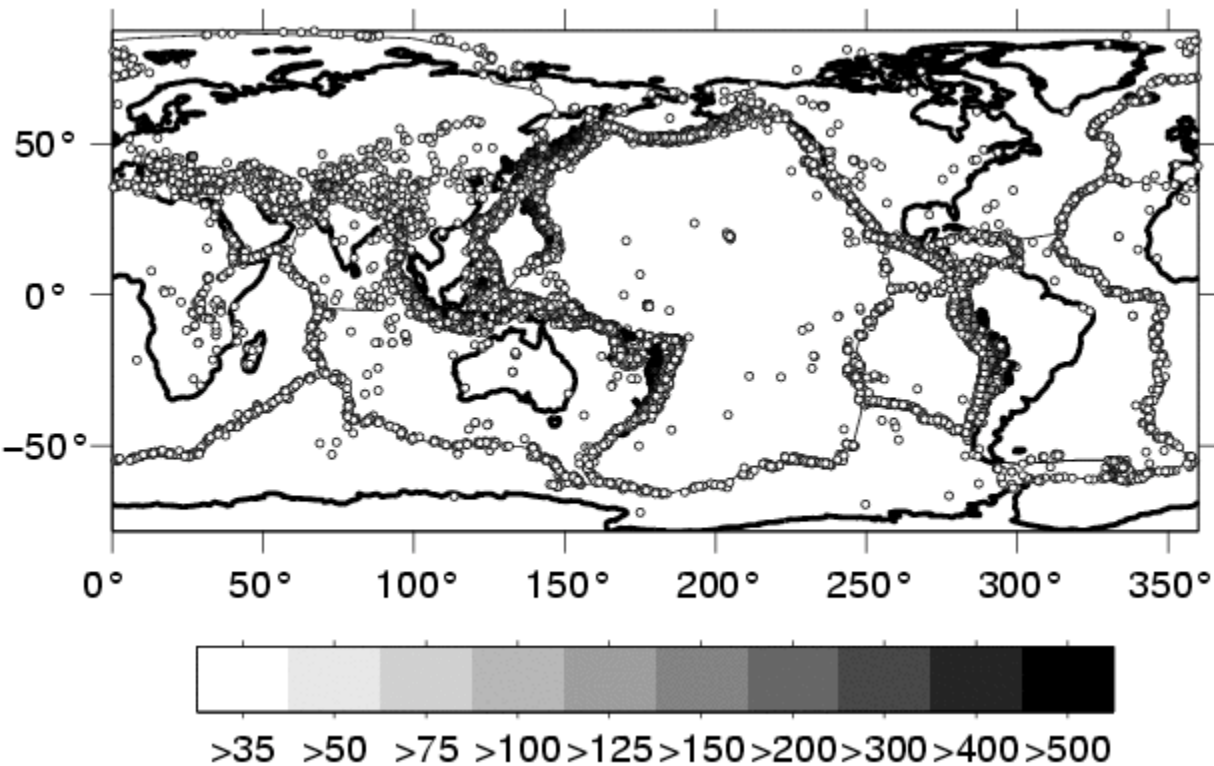
- **X-ray**: exponential of a line integral
  - **S-ray**: raypath itself is a function of velocity
- } **non-linear functions!**
- Earth **coverage** is non-continuous
  - “Experiment” is done by nature and **not repeatable**
  - Earthquake **source parameters** (location, time) is uncertain

## Remedy:

---

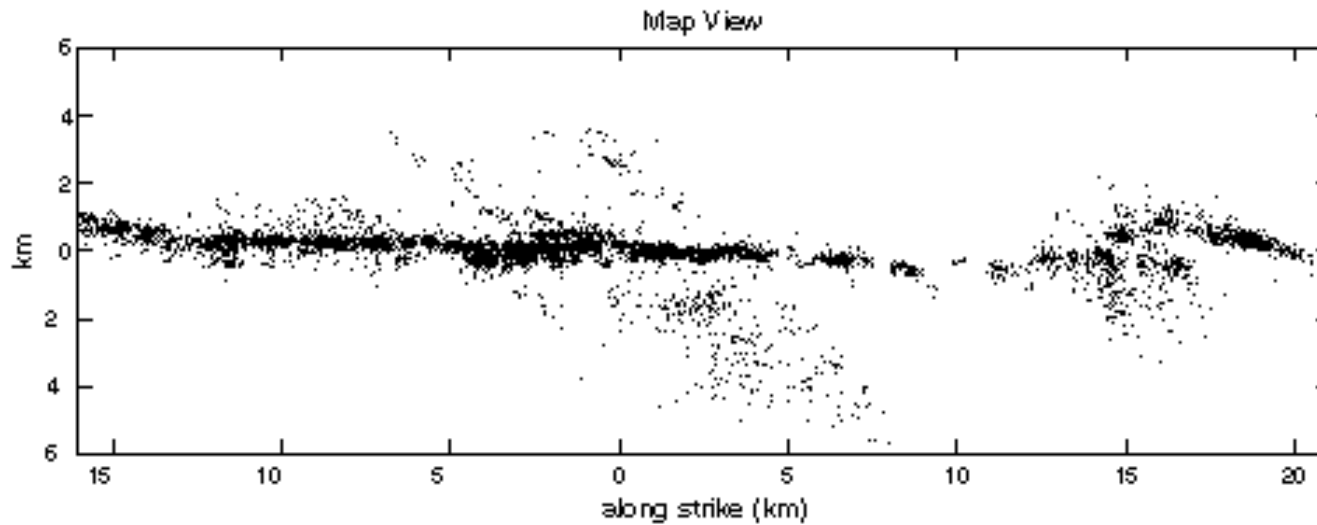
- Linearization
- Discretization
- Regularization (*a priori* information)

# Non-continuous **source coverage**

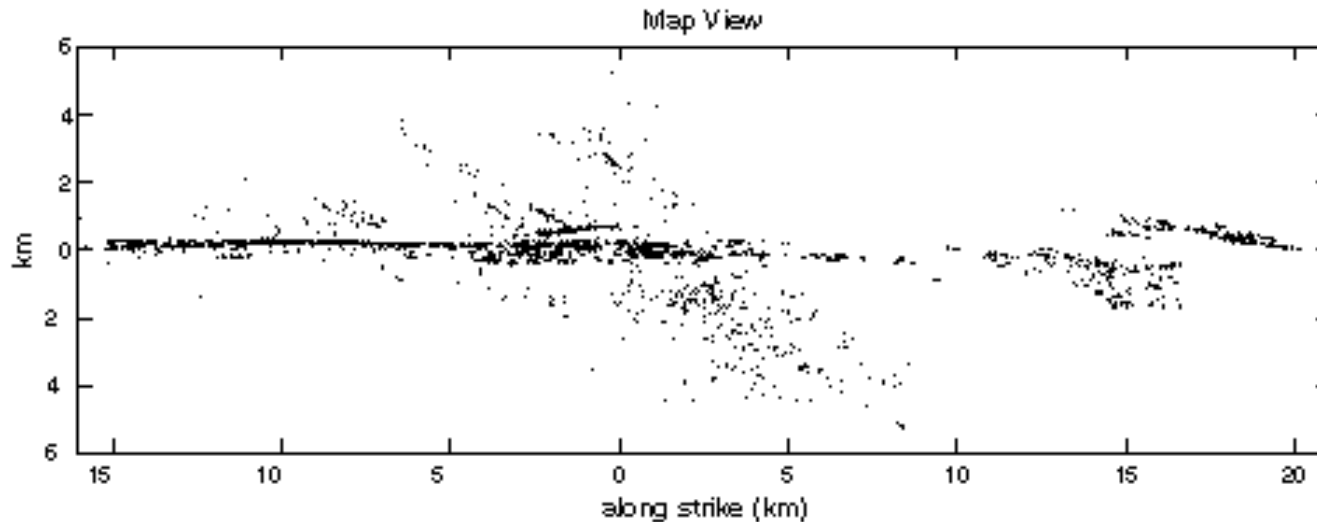


**The CMT catalog of large events**

# Source location – (in)extricably linked



Before



After

Source relocation is big business.



## *Recipe, Step 1: Linearize!*

---

### **X-ray**

Approximate  $\exp(-x) \approx 1 - x$ .

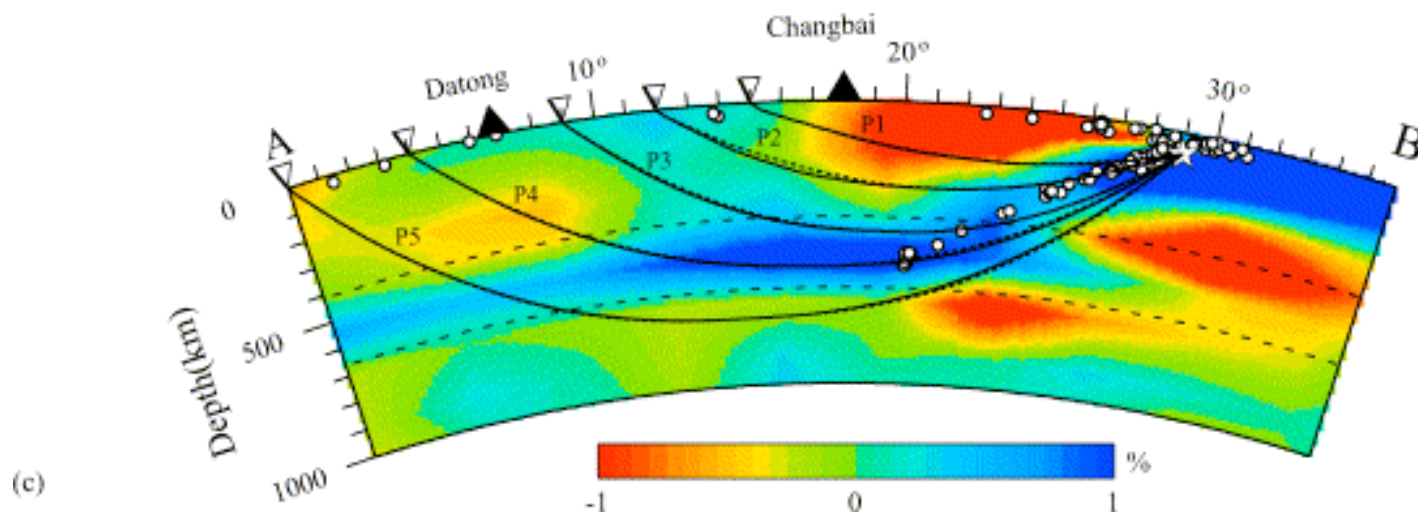
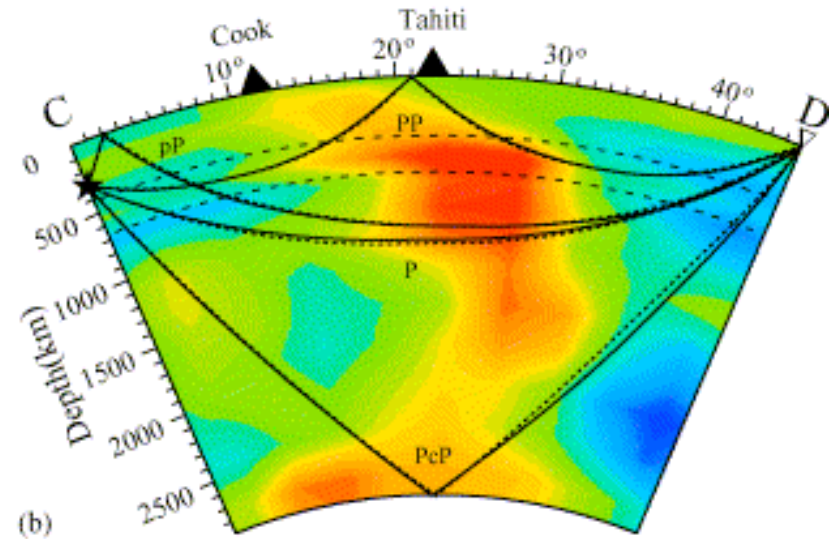
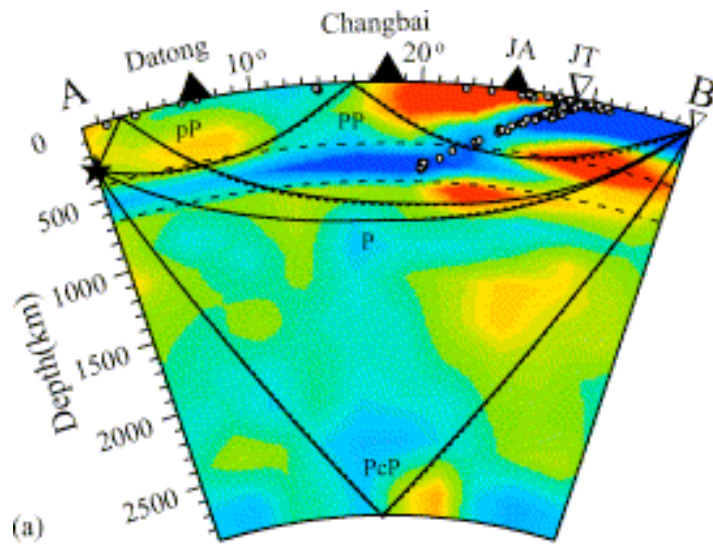
### **S-ray**

**Fermat's principle:** For a small perturbation of the path, the travel-time (anomaly) is stationary. Using the *slowness*:

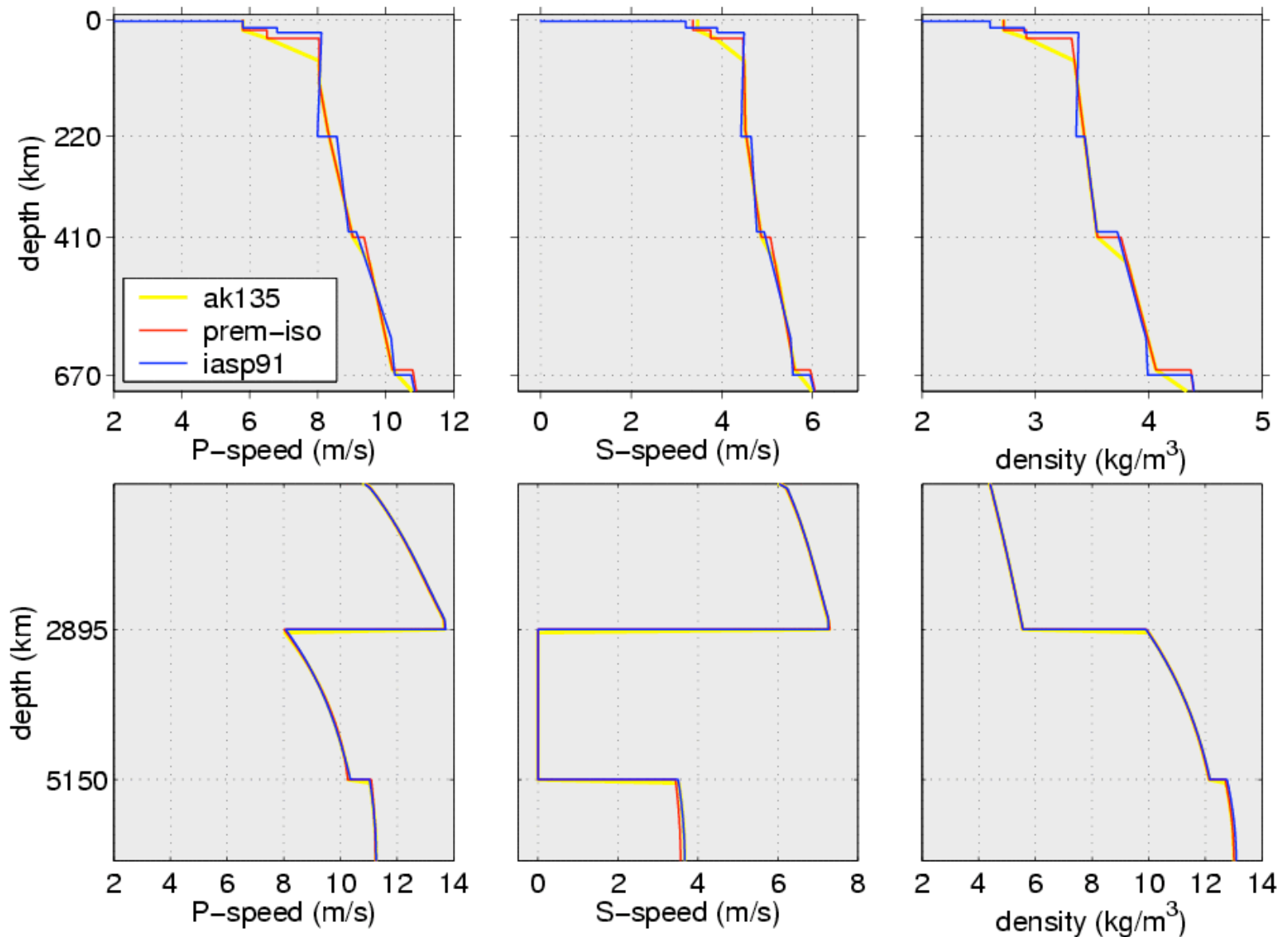
$$\delta s = \frac{1}{\delta v} \rightarrow \delta(\delta t) + \mathcal{O}[(\delta t)^2]. \quad (4)$$

This highlights the importance of the **reference model**, usually a radial model  $v(r)$ , such as PREM, AK135, IASP91.

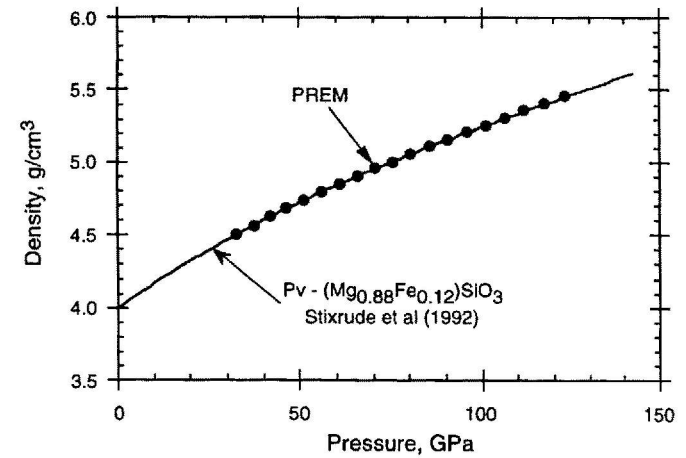
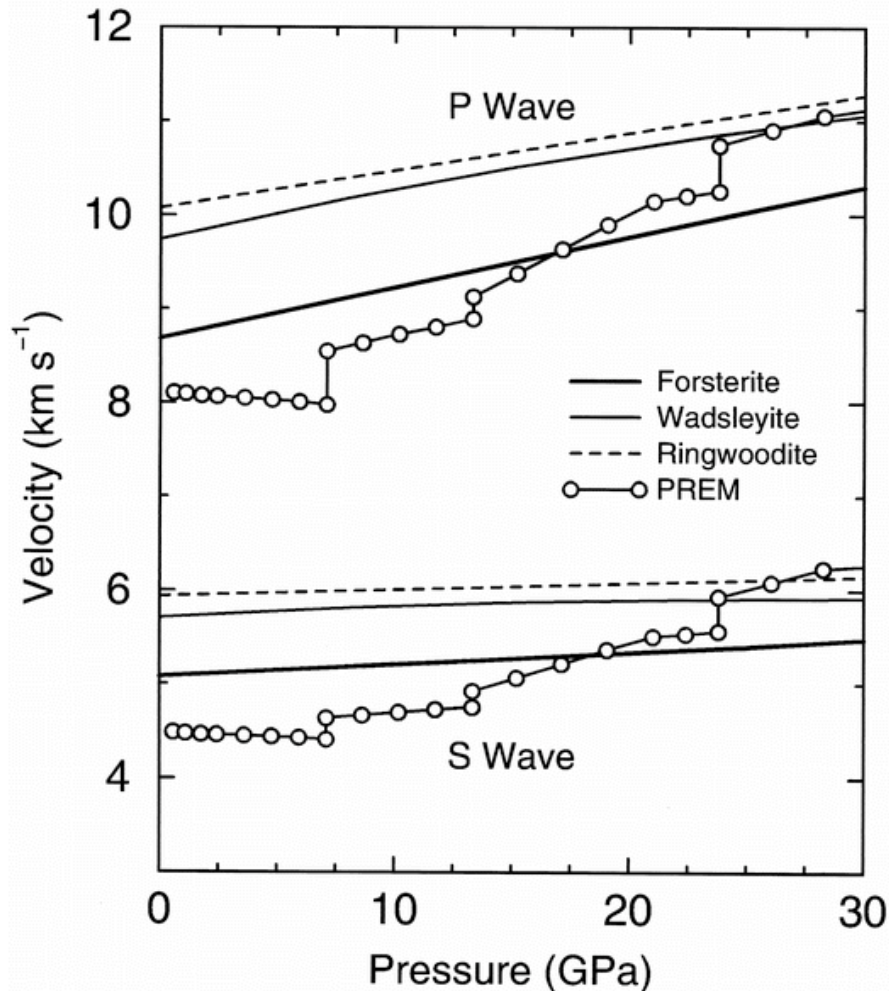
# Fermat's Principle at Work for you



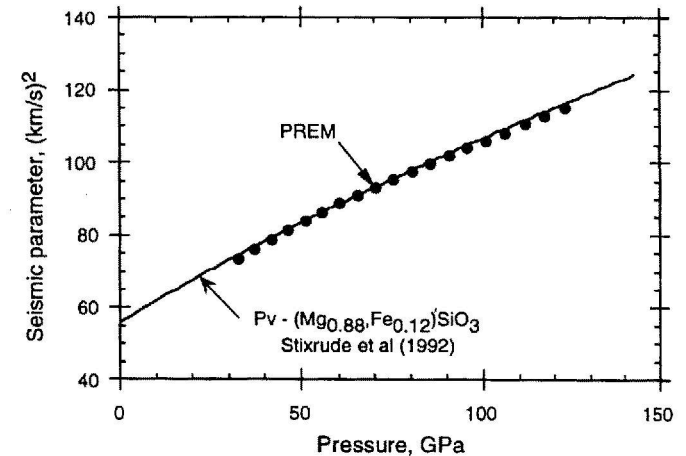
# The reference Earth: Radial models



# ... and at least some of it is true...



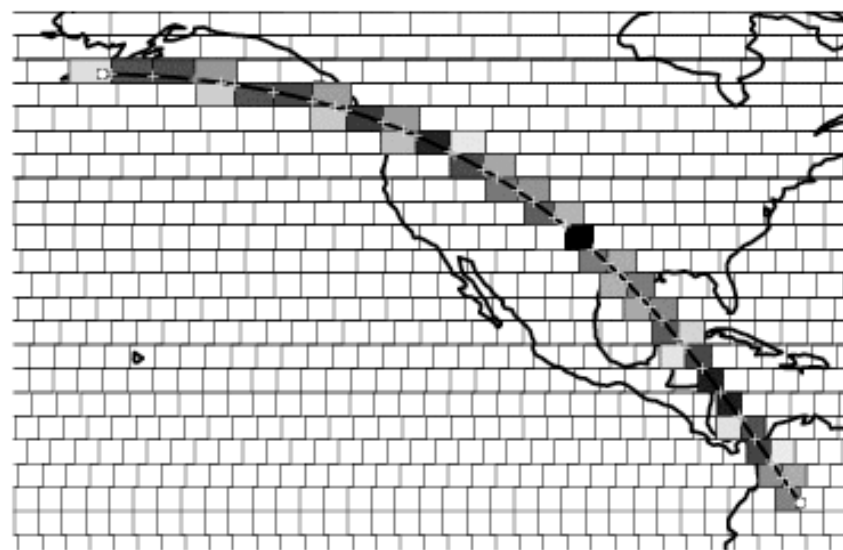
(a)



(b)

## Recipe, Step 2: Discretize!

---



For a set of seismic rays  $i = 1 \rightarrow M$ , calculate the length spent in each of  $j = 1 \rightarrow N$  grid boxes, in each of which it accumulates a proportional fraction of the total travel-time anomaly  $\delta t$ .

$$\delta t_i = L_{ij} \delta s_j \quad \text{or} \quad \delta \mathbf{t} = \mathbf{L} \cdot \delta \mathbf{s} \quad (5)$$

$$\begin{array}{l} \text{M travel-time} \\ \text{anomalies} \end{array} \begin{bmatrix} \vdots \\ \delta t_i \\ \vdots \end{bmatrix} = \begin{array}{c} \begin{bmatrix} \vdots \\ \dots & L_{ij} & \dots \\ \vdots \end{bmatrix} \\ \text{M} \times \text{N sensitivity matrix} \end{array} \times \begin{array}{c} \begin{bmatrix} \vdots \\ \delta s_j \\ \vdots \end{bmatrix} \\ \text{N slowness} \\ \text{perturbations} \end{array} \quad (6)$$



# Letting it simmer: Solving inverse problems

---

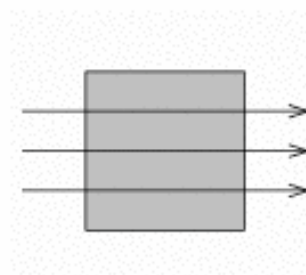
We have:  $\mathbf{G} \cdot \mathbf{m} = \mathbf{d}$ , which is **linear**.

You think:  $\mathbf{m} = \mathbf{G}^{-1} \cdot \mathbf{d}$ , but we **can't invert** a non-square  $M \times N$  matrix.

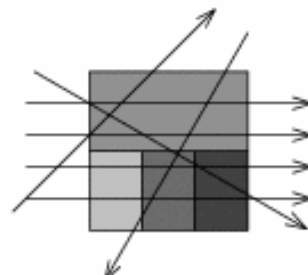
You think:  $\mathbf{G}^T \cdot \mathbf{G}$  is square, let's solve  $\mathbf{G}^T \cdot \mathbf{G} \cdot \mathbf{m} = \mathbf{G}^T \cdot \mathbf{d}$ .

You try:  $\mathbf{m} = (\mathbf{G}^T \cdot \mathbf{G})^{-1} \cdot \mathbf{G}^T \cdot \mathbf{d}$ .

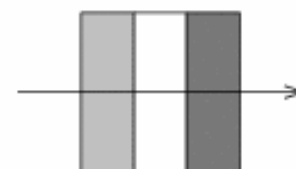
Alas!  $\mathbf{G}^T \cdot \mathbf{G}$  may be singular, ill-conditioned, under/over-determined, have (near-)zero eigenvalues, and thus be **not-invertible**.



**over-determined,  $M > N$**



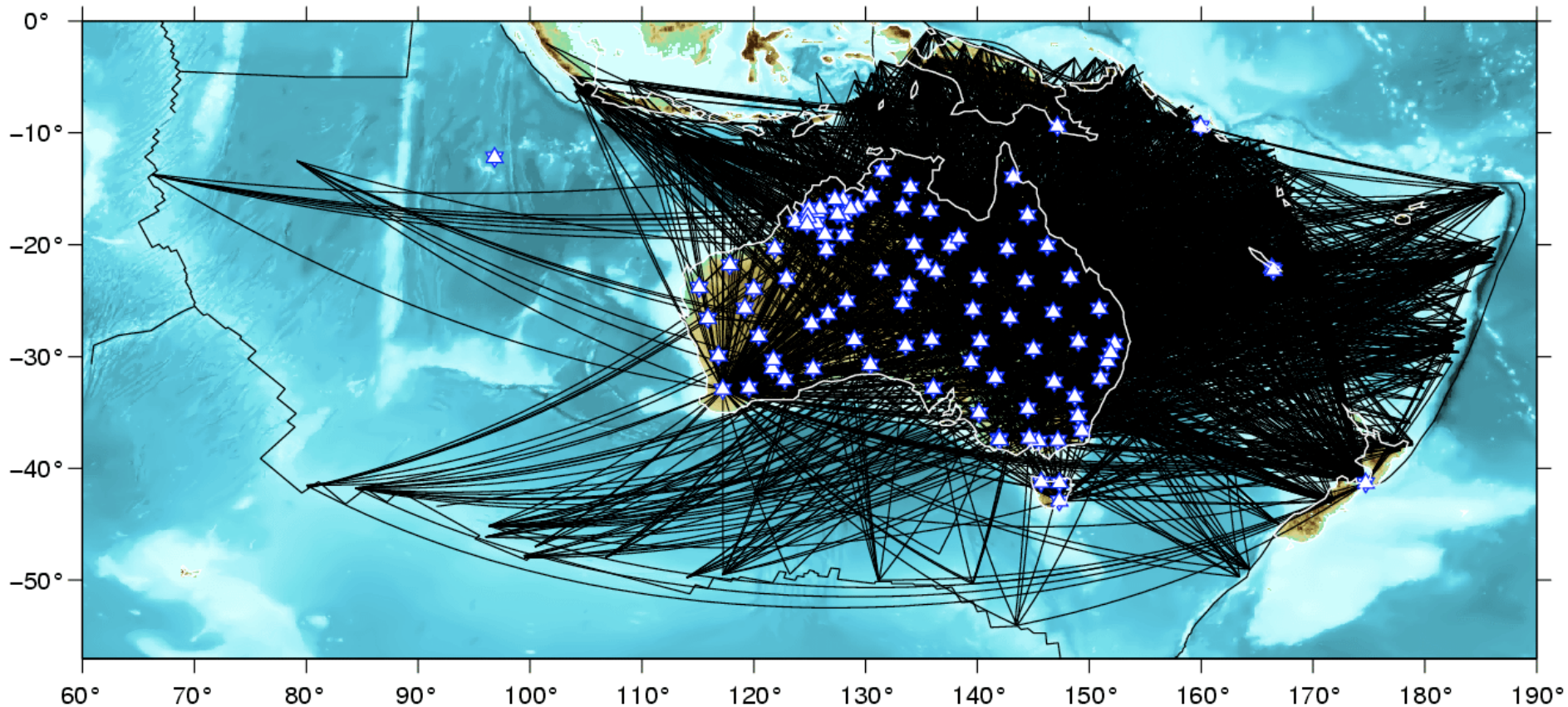
**mixed-determined**



**under-determined,  $M < N$**

# Receiver coverage

## Picking the right continent

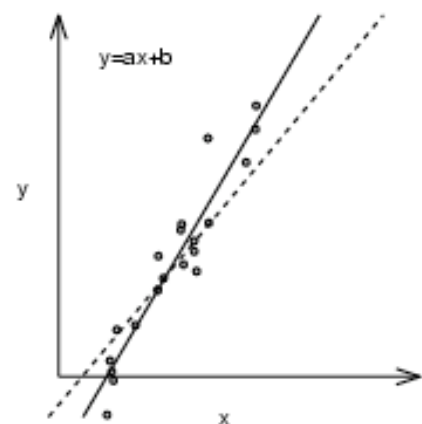


A dense path coverage minimizes the amount of a priori information needed

## Recipe, Step 3: Regularize!

---

### Over-determined: *More data than unknowns*



Define a *penalty function*  $\Phi$  on the *error*  $\mathbf{e}$ ,  
and minimize, by least-squares:

$$\Phi = [\mathbf{G} \cdot \mathbf{m} - \mathbf{d}]^2 = \mathbf{e}^T \cdot \mathbf{e} \quad \text{by} \quad \frac{\partial \Phi}{\partial m_i} = 0. \quad (7)$$

This is a minimization in the *data space*.

### Under-determined: *More unknowns than data*

Add equations that minimize some norm in the *model space*:

$$\Phi = \mathbf{e}^T \cdot \mathbf{e} + \mathbf{m}^T \cdot (\mathbf{A}^T \cdot \mathbf{A}) \cdot \mathbf{m}. \quad (8)$$

If  $\mathbf{A} = \mathbf{I}$  the identity matrix  $\rightarrow$  minimum model norm: **norm damping**.

If  $\mathbf{A} = \mathbf{D}$  a difference matrix  $\rightarrow$  minimum-roughness: **smoothing**.

# Regularization: the Mathematics

## Numerical Methods for the Solution of Ill-Posed Problems

by

A. N. Tikhonov  
A. V. Goncharenko  
V. V. Stepanov  
A. G. Yagola

Moscow State University,  
Moscow, Russia



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## ILL-POSED AND INVERSE PROBLEMS

DEDICATED TO ACADEMICIAN MIKHAIL MIRONOVICH LUR'YEV  
ON THE OCCASION OF HIS 70<sup>th</sup> BIRTHDAY



Editors:

V.G. Romanov, S.I. Kabanikhin,  
Yu.E. Anikonov and A.L. Bukhgeim

/// VSP ///

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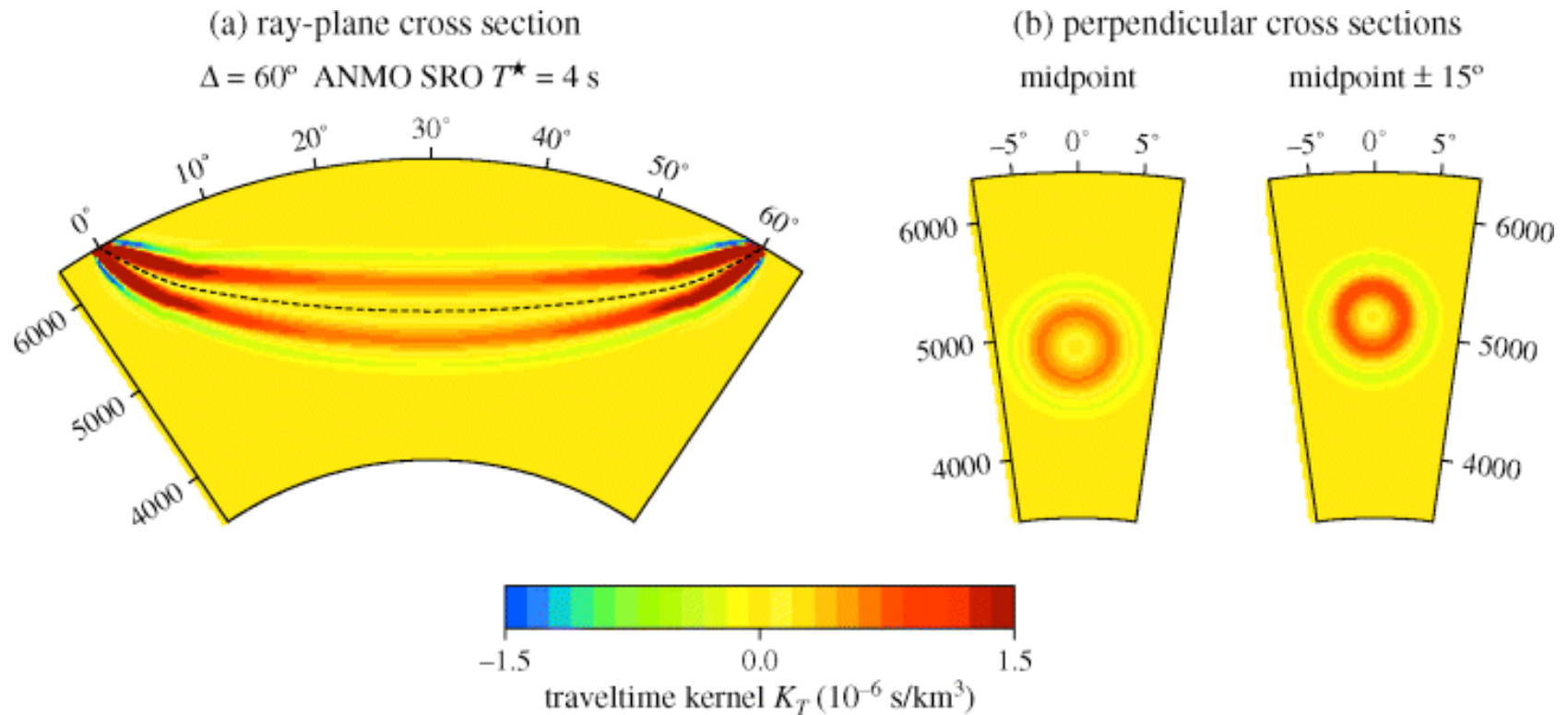
Methods in Geochemistry and Geophysics, 36

## GEOPHYSICAL INVERSE THEORY AND REGULARIZATION PROBLEMS

M.S. ZHDANOV

ELSEVIER

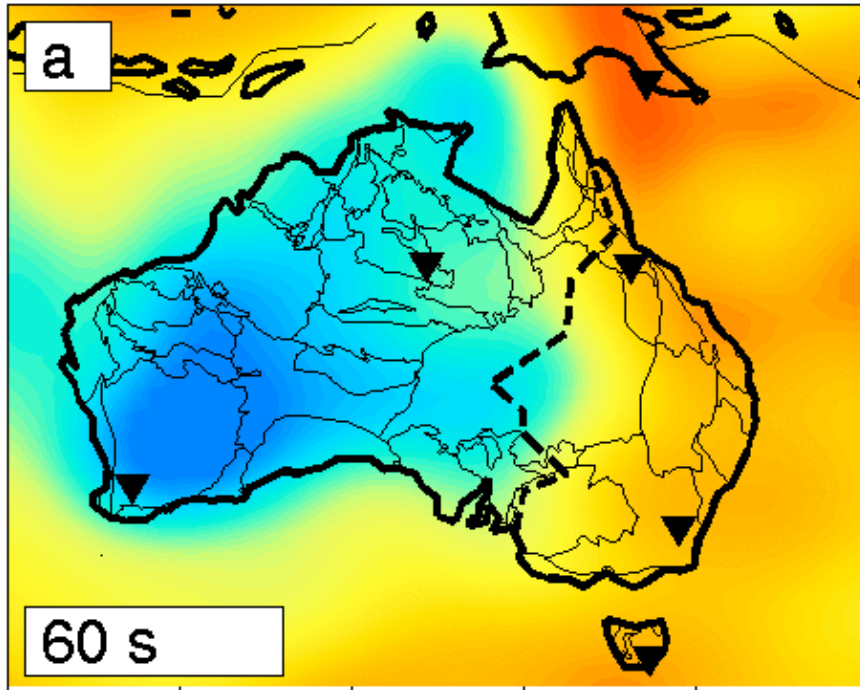
# Regularization: the Physics



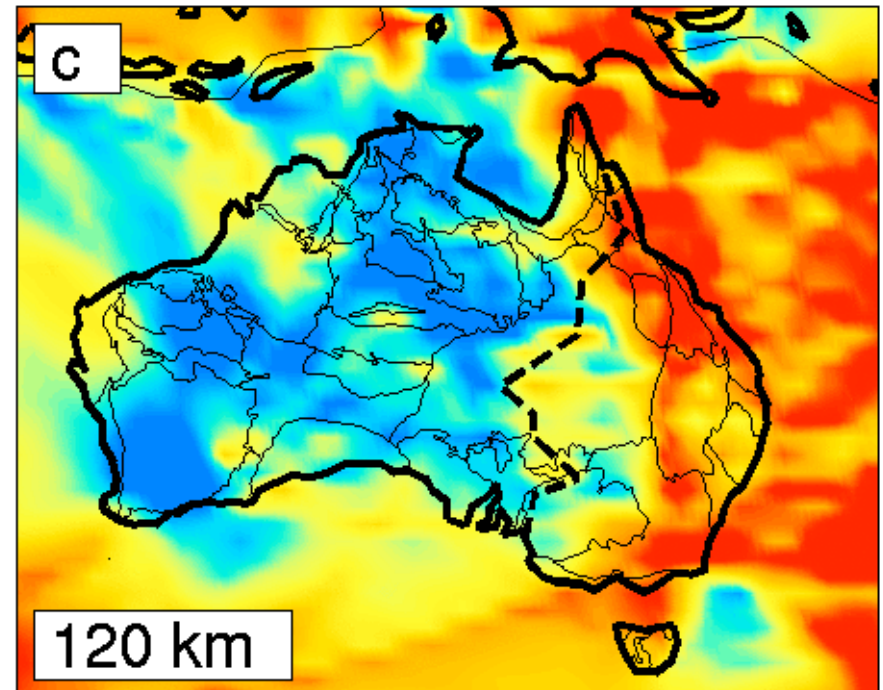
Such “fat” rays sample more of the Earth and thus we need fewer of them to have a well-constrained tomographic problem.



# Regularization: the Art

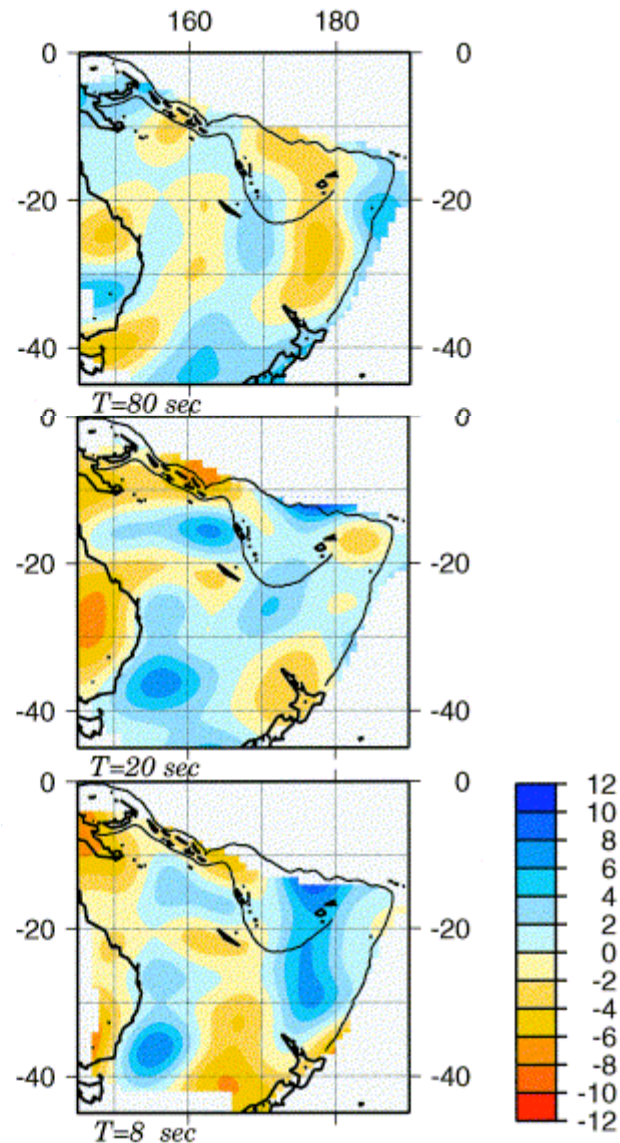


Too much?  
Too smooth?

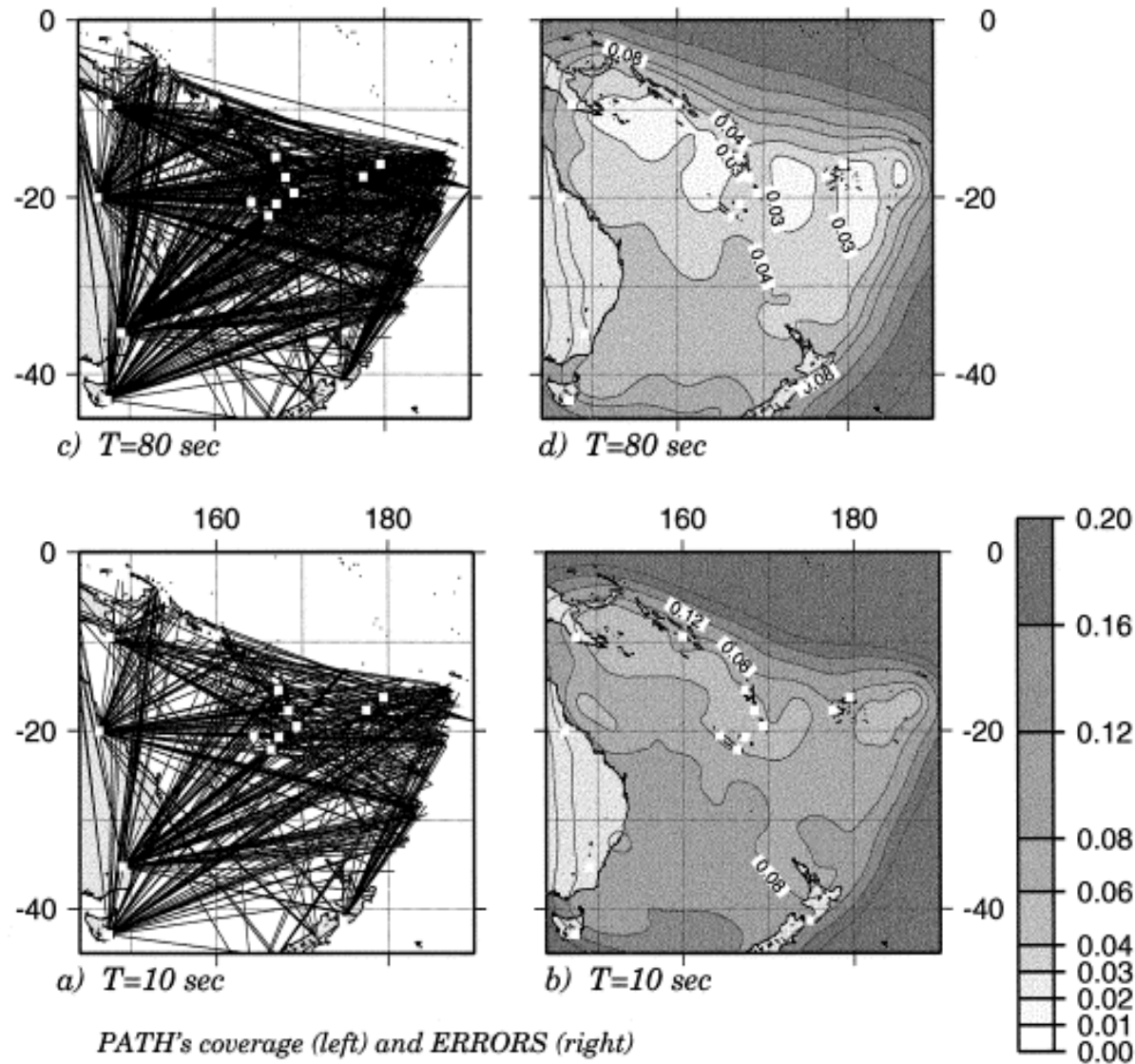


Too little?  
Too rough?

# How to interpret seismic models

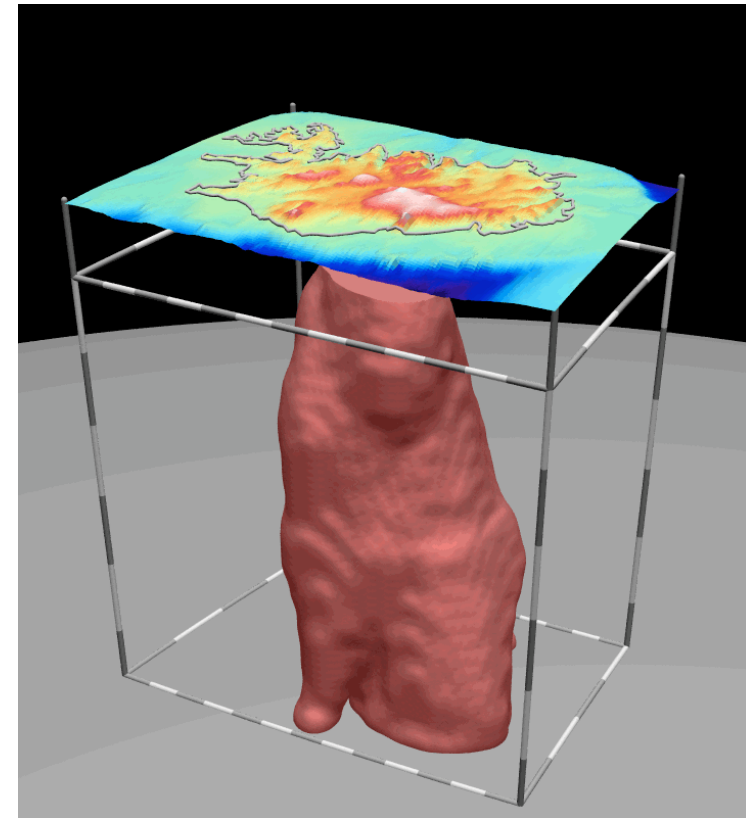
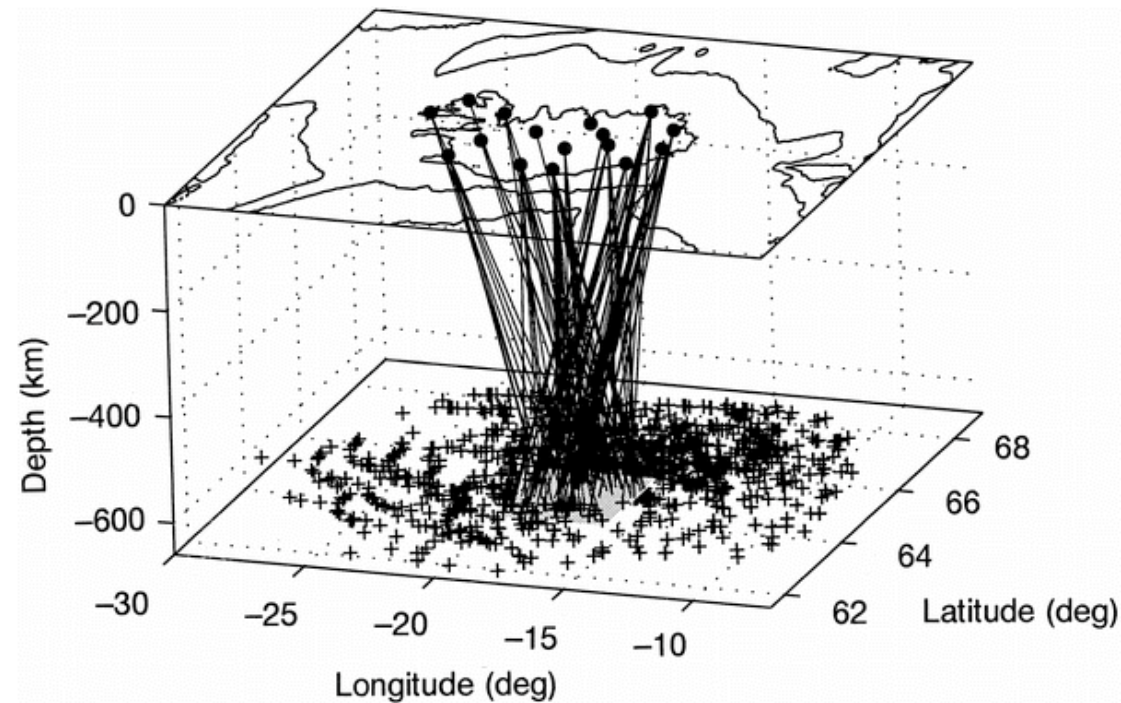


# Demand to see the ray paths



*PATH's coverage (left) and ERRORS (right)  
RAYLEIGH WAVES*

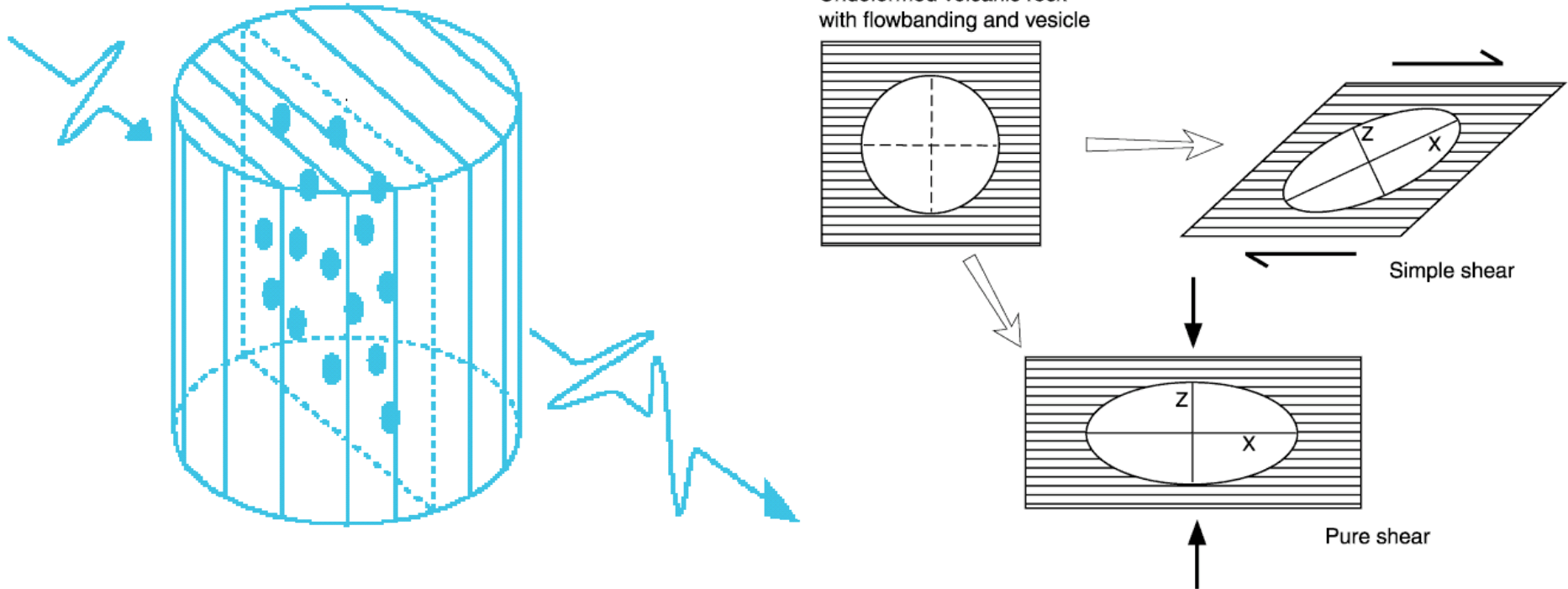
# Nature isn't always kind



# Seismic anisotropy

Wave speeds depend on

propagation direction and polarization:



No surprise: elasticity maps stress and strain,  
and both depend on three directions



# **Polarization anisotropy**

---

- The particles of Love and Raleigh surface waves move in orthogonal directions
- SH and SV body waves sometimes exhibit clear splitting

# **Azimuthal anisotropy**

---

- It's usually very hard to separate whether the time difference arises from an anisotropic direction or an isotropic wave speed difference (aka heterogeneity)

# Why is this so hard?

For 3-D heterogeneity and slight anisotropy:

$$\delta\hat{\beta}_V = \delta\beta_V^{TI} + \frac{G_c}{2\rho\beta_V} \cos 2\theta + \frac{G_s}{2\rho\beta_V} \sin 2\theta \quad (3)$$

Maximum direction is related to fast axis of anisotropic minerals:

$$G = \sqrt{G_c^2 + G_s^2} \quad \text{and} \quad \Psi_{\max} = \frac{1}{2} \arctan \frac{G_s}{G_c} \quad (4)$$

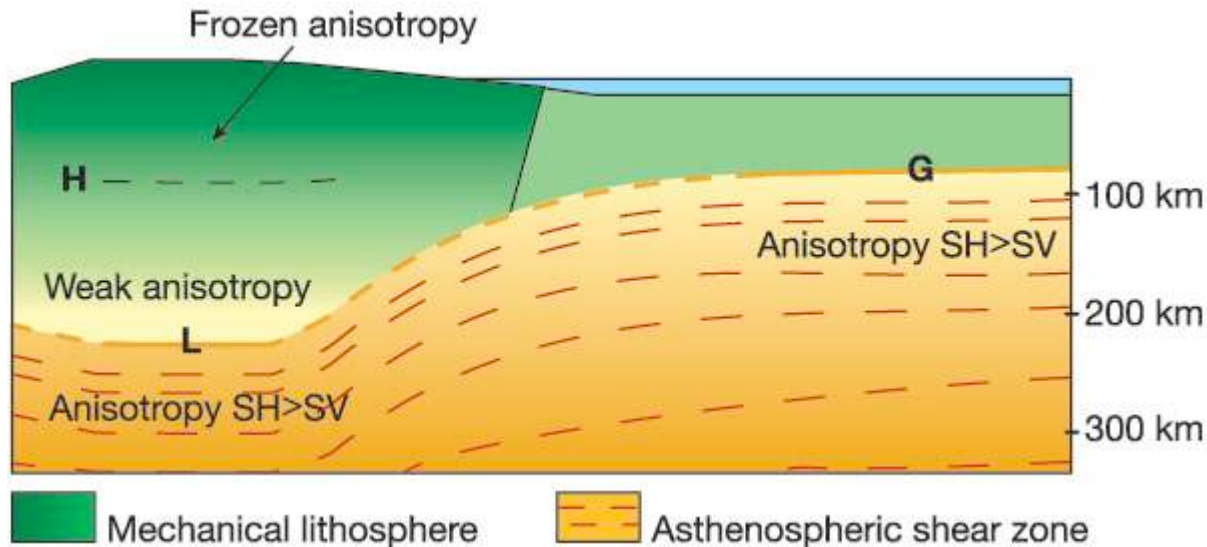
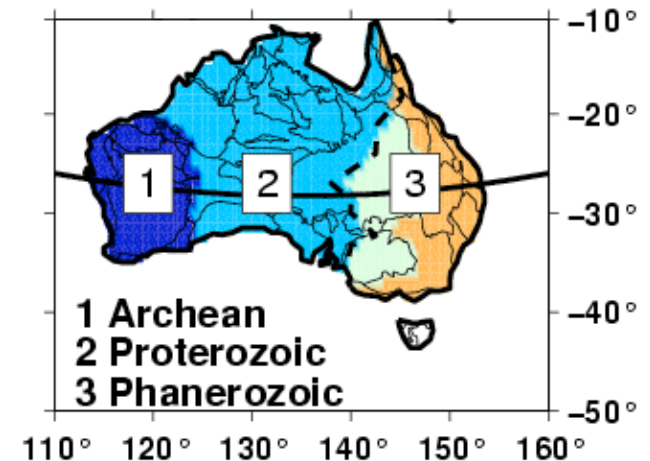
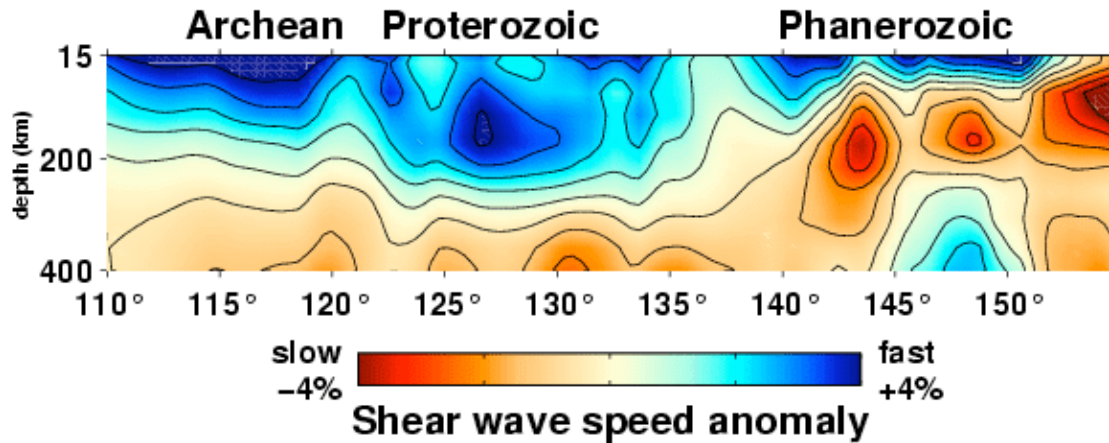
It's very hard to tell whether a phase comes in early because it went through a fast patch or because it came in a fast direction – heterogeneity and anisotropy “trade off.”

# Questions to ask of the tomographer

- How is the forward model computed?
- What is the ray coverage?
- What (sort of) damping did you use?
- What does velocity estimation trade off with?
- What is the grid size / the correlation length?
- How are different data sets weighted?
- How far is the final from the starting model?
- Does the starting model have discontinuities?
- How is the surface/depth parameterization
- Is your sensitivity 1-D, 2-D, or 3-D?

## Journey to Middle Earth, Part I:

# The continental lithosphere

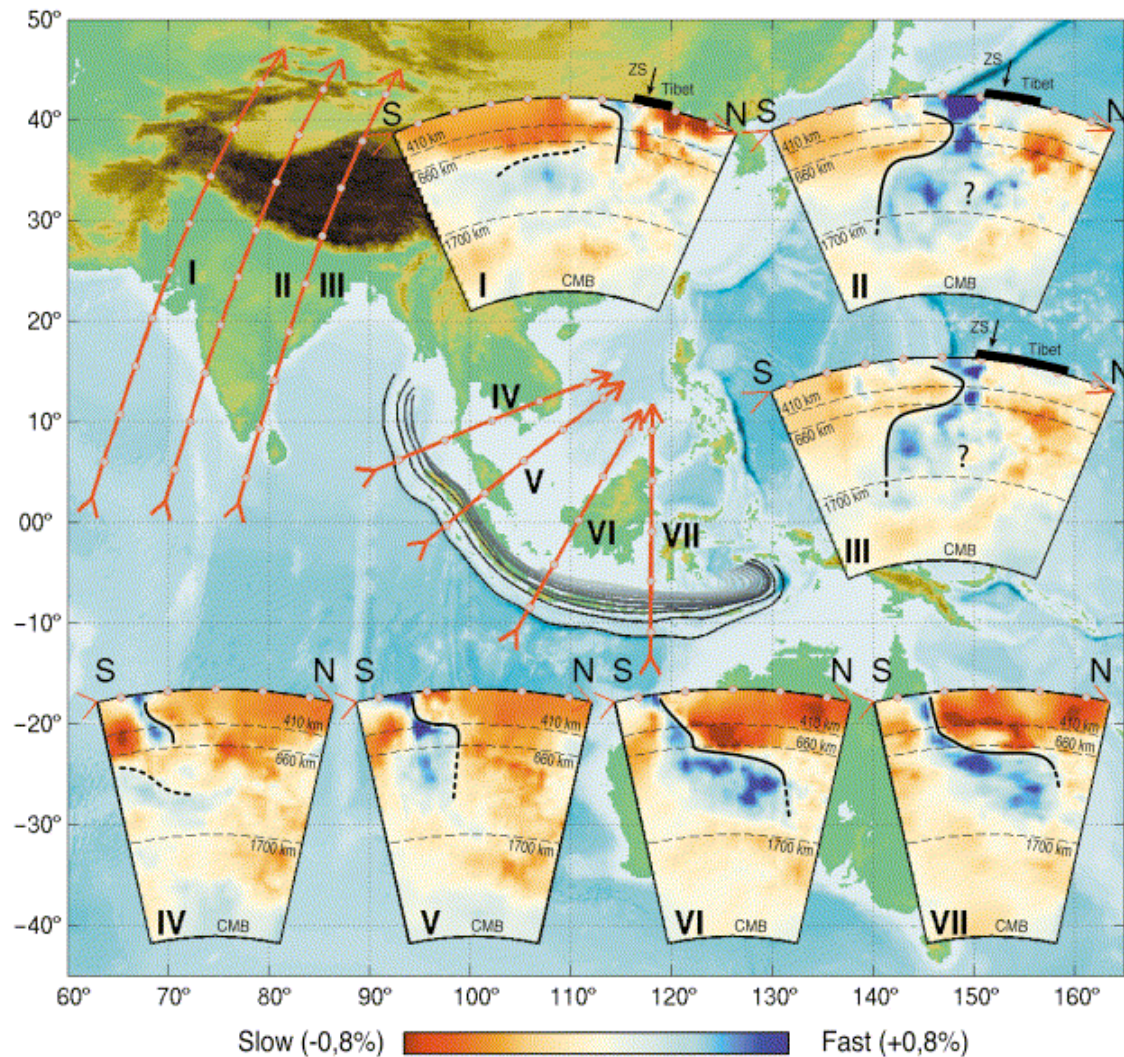


Simons, *GRL*, 2002

Gung, *Nature*, 2003

# Journey to Middle Earth, Part II:

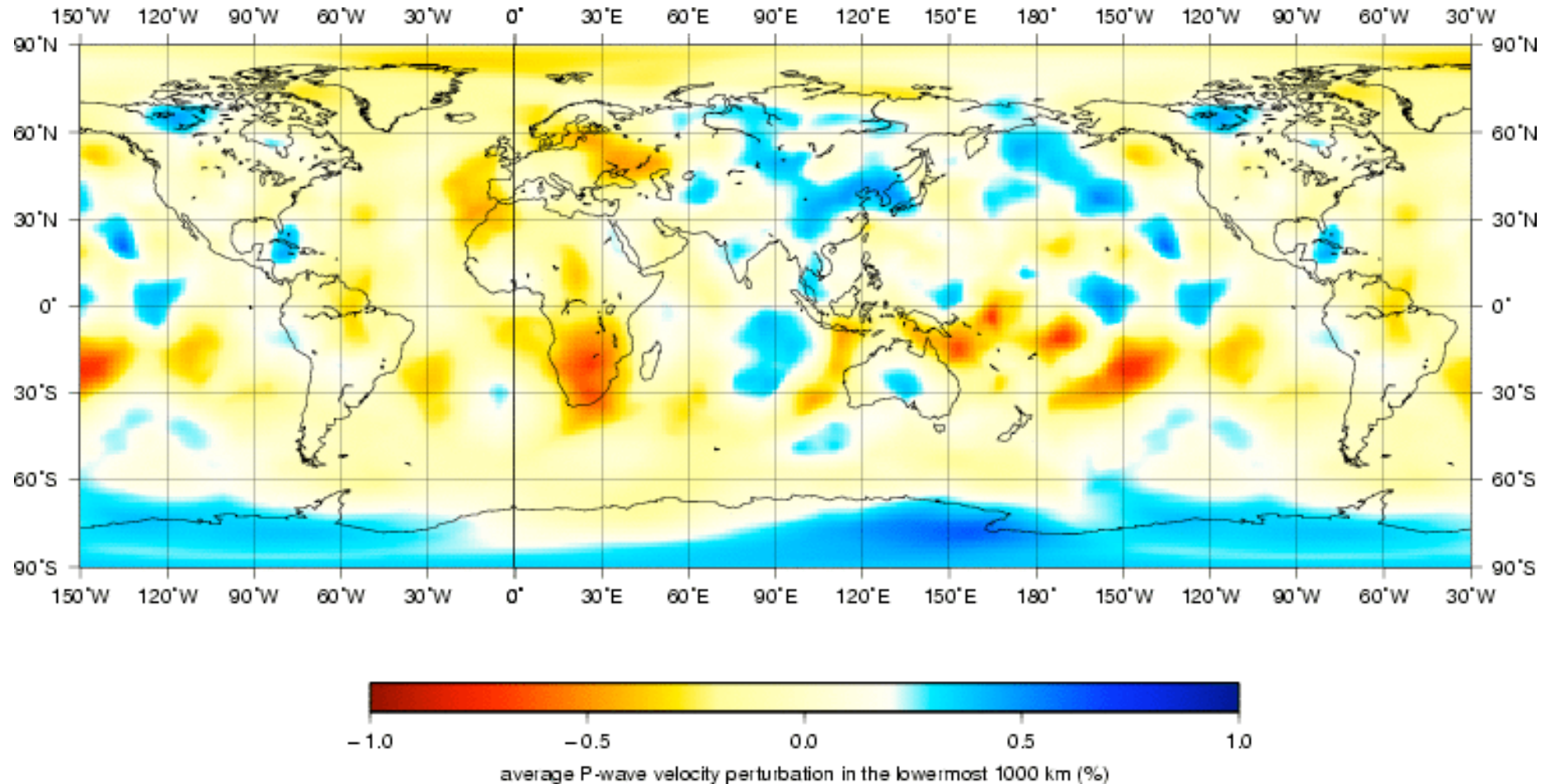
## Subduction zones





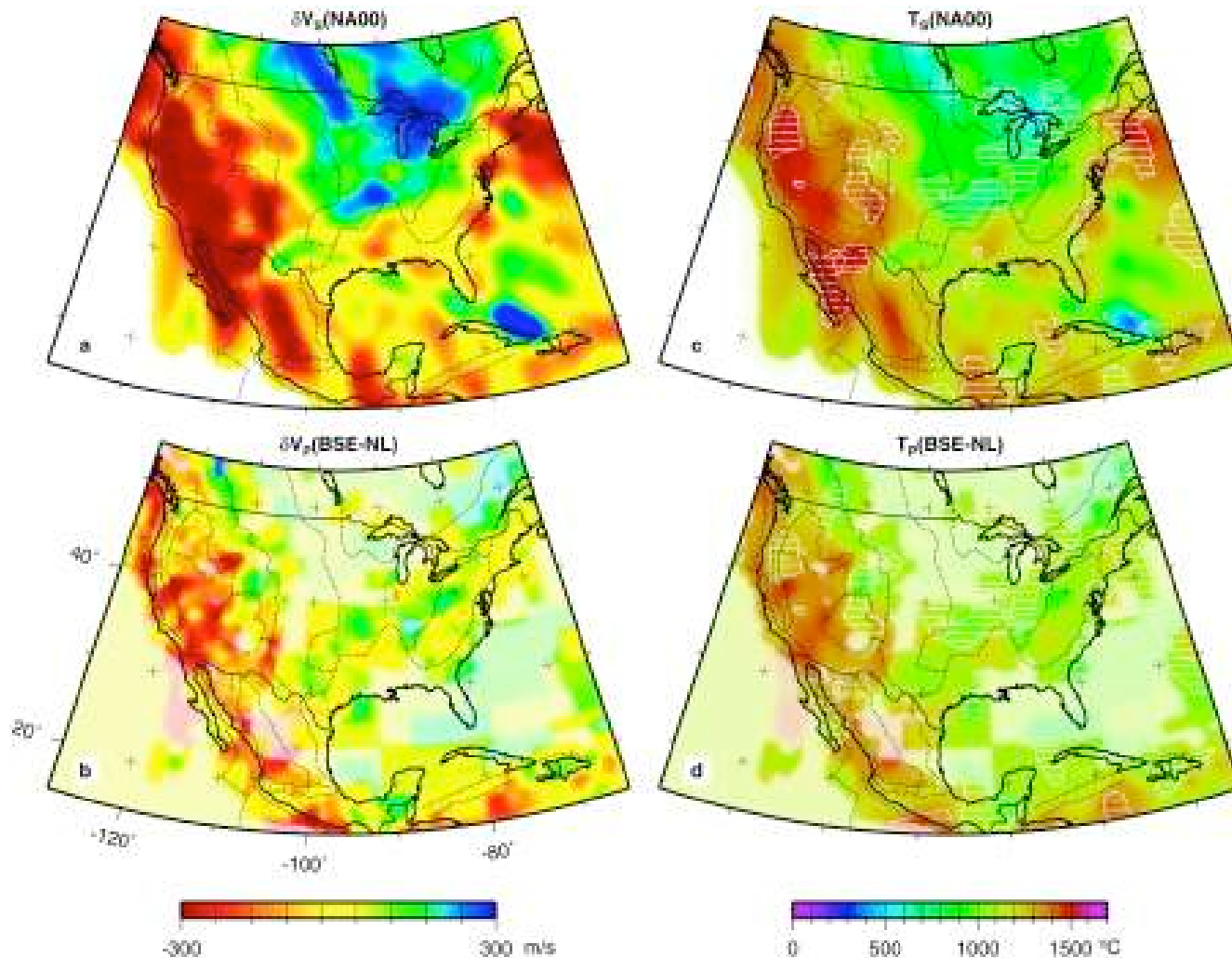
# Journey to Middle Earth, Part III:

# Deep mantle plumes



# What does it all mean? Part I:

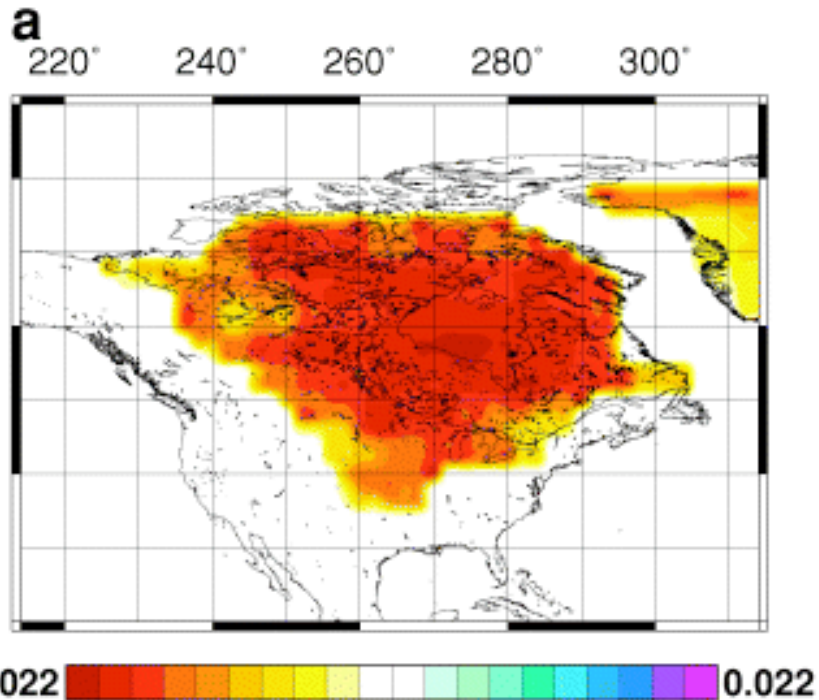
## Temperature anomalies



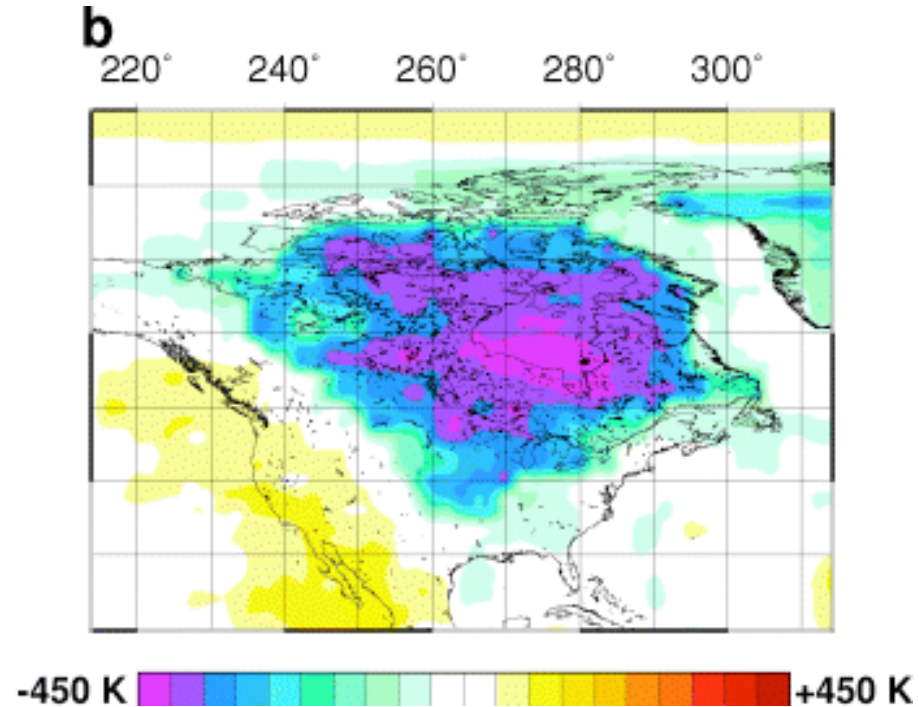
110 km

# What does it all mean? Part II:

## Compositional anomalies



$\Delta\text{Fe}/(\text{Fe} + \text{Mg})$

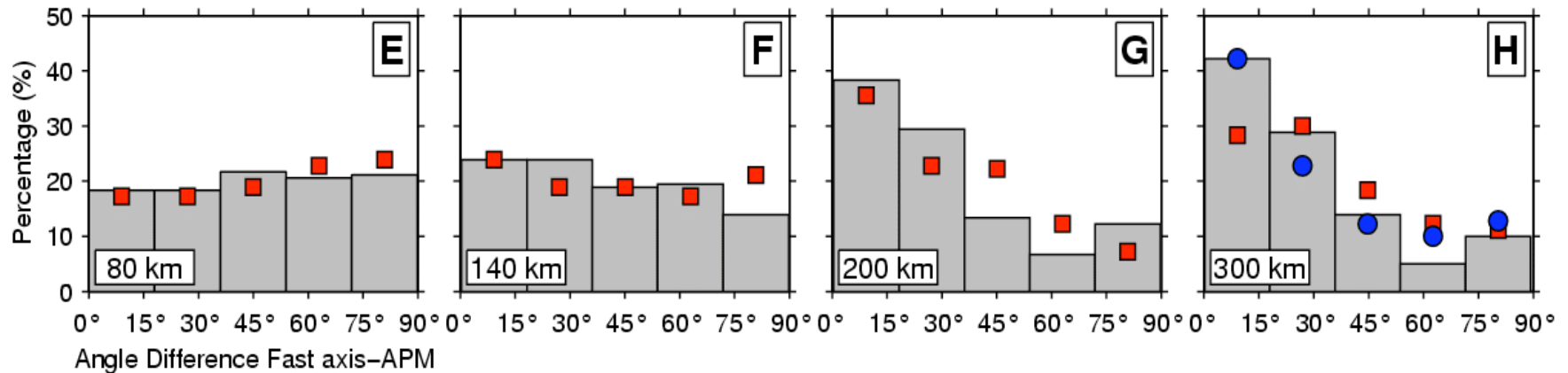
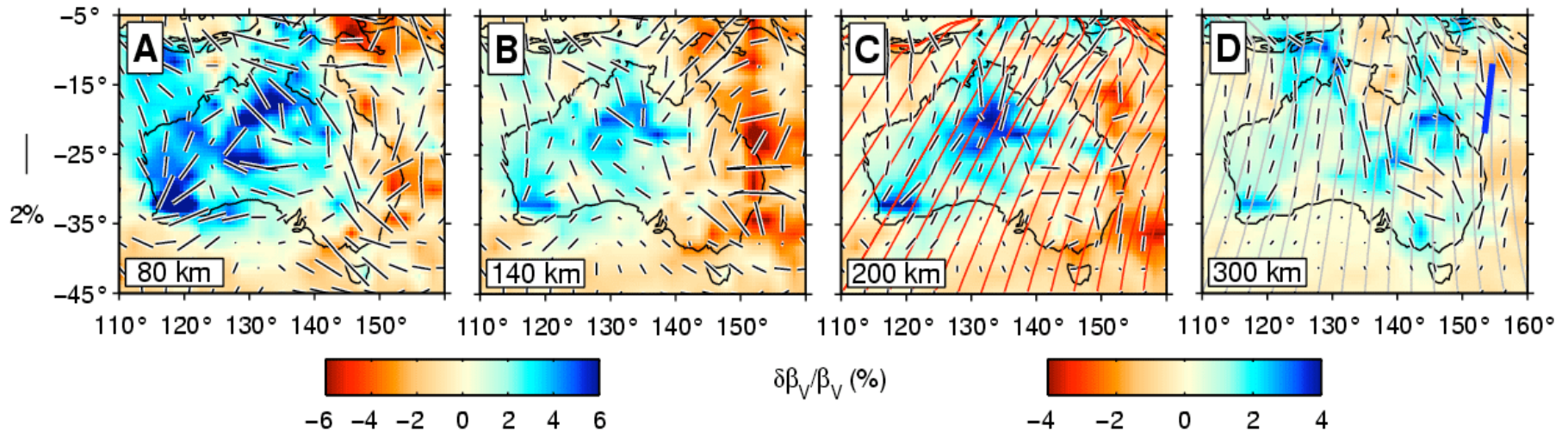


$\Delta T$

150 km

# What does it all mean? Part III:

## Deformation in the mantle



Fossil

Contemporaneous

Simons, *EPSL*, 2003

# Conclusions

- Ultimately, seismology can only tell us where, or in which direction, wave propagation is faster or slower than a reference model
- The non-seismologist has to know the basics of inverse problem modeling, understand the sometimes poor constraints, and be critical
- Improvements are being made: better data, better forward models, better inversions
- As much as with the *a posteriori* interpretation, the community needs to help defining *a priori* acceptable starting models



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**More equations, for  
completeness**

# A linear system of equations

We're attempting to solve

$$\mathbf{d} = \mathbf{G} \cdot \mathbf{m} \quad (1)$$

Minimize penalty function of weighted error and model norms

$$\begin{aligned} \Phi = & (\mathbf{d} - \mathbf{G} \cdot \mathbf{m}) \cdot \mathbf{A}^{-1} \cdot (\mathbf{d} - \mathbf{G} \cdot \mathbf{m}) \\ & + \mathbf{m} \cdot \mathbf{B}^{-1} \cdot \mathbf{m} \end{aligned} \quad (2)$$

In matrix form, solve

$$\begin{bmatrix} \mathbf{A}^{-1/2} \cdot \mathbf{G} \\ \mathbf{B}^{-1/2} \end{bmatrix} \cdot \mathbf{m} = \begin{bmatrix} \mathbf{A}^{-1/2} \cdot \mathbf{d} \\ 0 \end{bmatrix} \quad (3)$$

Solution

$$\boxed{\mathbf{m} = (\mathbf{B}^{-1} + \mathbf{G}^T \cdot \mathbf{A}^{-1} \cdot \mathbf{G})^{-1} \cdot \mathbf{G}^T \cdot \mathbf{A}^{-1} \cdot \mathbf{d}} \quad (4)$$



# Norm and first gradient regularization

For  $\mathbf{A}^{-1}$ , use the inverse of the data covariance matrix  $\mathbf{C}_d$  (BLUE)

For  $\mathbf{B}^{-1}$ , use the identity matrix  $\mathbf{I}$  plus the squared first derivative

$$\mathbf{D}_1 = \begin{pmatrix} \dots & 0 & 0 \\ -1 & 1 & 0 \\ 0 & -1 & 1 \\ 0 & 0 & \dots \end{pmatrix} \quad (5)$$

Minimize weighted penalty function

$$\begin{aligned} \Phi = & (\mathbf{d} - \mathbf{G} \cdot \mathbf{m}) \cdot \mathbf{C}_d^{-1} \cdot (\mathbf{d} - \mathbf{G} \cdot \mathbf{m}) \\ & + \alpha \mathbf{m} \cdot \mathbf{I} \cdot \mathbf{m} + \beta \mathbf{m} \cdot \mathbf{D}_1^2 \cdot \mathbf{m} \end{aligned} \quad (6)$$

Solution

$$\boxed{\mathbf{m} = (\alpha \mathbf{I} + \beta \mathbf{D}_1^2 + \mathbf{G}^T \cdot \mathbf{C}_d^{-1} \cdot \mathbf{G})^{-1} \cdot \mathbf{G}^T \cdot \mathbf{C}_d^{-1} \cdot \mathbf{d}} \quad (7)$$

# Bayesian inversion

Gaussian *a priori* probability function on the model parameters

$$\rho(\mathbf{m}) \propto \exp \left( -\frac{1}{2} \mathbf{m} \cdot \mathbf{C}_m^{-1} \cdot \mathbf{m} \right) \quad (8)$$

Maximize joint distribution of data, model, subject to  $\mathbf{d} = \mathbf{G} \cdot \mathbf{m}$

$$\boxed{\mathbf{m} = \mathbf{C}_m \cdot \mathbf{G}^T \cdot (\mathbf{C}_d + \mathbf{G} \cdot \mathbf{C}_m \cdot \mathbf{G}^T)^{-1} \cdot \mathbf{d}} \quad (9)$$

Equivalent to (using a trivial matrix identity)

$$\mathbf{m} = (\mathbf{C}_m^{-1} + \mathbf{G}^T \cdot \mathbf{C}_d^{-1} \cdot \mathbf{G})^{-1} \cdot \mathbf{G}^T \cdot \mathbf{C}_d^{-1} \cdot \mathbf{d} \quad (10)$$

So in choosing norm and gradient regularization we've identified

$$\mathbf{C}_m^{-1} = \alpha \mathbf{I} + \beta \mathbf{D}_1^2 \quad (11)$$

This imposes a particular form of the *a priori* covariance  $\mathbf{C}_m$

# To Bayes or not to Bayes, what's the question?

*A priori* model covariance function with correlation length  $L$

$$C_{\mathbf{m}}(\mathbf{r}_1, \mathbf{r}_2) = \sigma^2 \exp \left( -\frac{|\mathbf{r}_1, \mathbf{r}_2|^2}{2L^2} \right) \quad (12)$$

The following equivalence holds [*Yanovskaya and Ditmar, 1990*]

$$\mathbf{m} \cdot \mathbf{C}_{\mathbf{m}}^{-1} \cdot \mathbf{m} = \frac{1}{2\pi} \frac{1}{(\sigma L)^2} \sum_{n=0}^{\infty} \frac{1}{n!} \left( \frac{L^2}{2} \right)^n \nabla^n \mathbf{m} \cdot \nabla^n \mathbf{m} \quad (13)$$

So indeed

$$\mathbf{C}_{\mathbf{m}}^{-1} = \alpha \mathbf{I} + \beta \mathbf{D}_1^2 + \text{higher-order terms} \quad (14)$$

# Exact resolution computation

For the linear problem, in a generalized sense,

$$\mathbf{m}^{\text{est}} = \mathbf{G}^{-\text{g}} \cdot \mathbf{d}^{\text{obs}} = \mathbf{G}^{-\text{g}} \cdot \mathbf{G} \cdot \mathbf{m}^{\text{true}} \quad (15)$$

The resolution matrix is given by

$$\mathbf{R} = \mathbf{G}^{-\text{g}} \cdot \mathbf{G} \quad (16)$$

In the Bayesian framework [*Montagner, 1986*]

$$R(\mathbf{r}, \mathbf{r}') = \delta(\mathbf{r} - \mathbf{r}') - \frac{C_{\text{p}}(\mathbf{r}, \mathbf{r}')}{C_{\text{m}}(\mathbf{r}, \mathbf{r}')} \quad (17)$$

This represents the degree to which we are able to reduce the *a priori* covariance  $C_{\text{m}}$  of the model parameters (the null-state of information) by obtaining the *a posteriori* covariance structure  $C_{\text{p}}$  after the inversion.