

So: for  $\textcircled{1} \rightarrow \textcircled{2}$  along path A  $\Delta S + \left[ \int_1^2 \left( \frac{\delta Q}{T} \right)_{\text{surround}} \right] > 0$  Irr.

$$\Delta S + \left[ \int_1^2 \left( \frac{\delta Q}{T} \right)_{\text{sum}} \right] = 0 \quad \text{Reversible.}$$

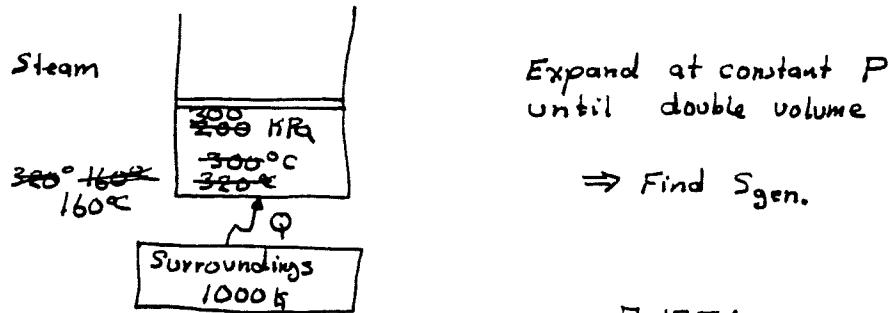
This principle is frequently written as  $\Rightarrow$

$$S_{\text{generated}} = \Delta S_{\text{system}} + \left[ \int_1^2 \left( \frac{\delta Q}{T} \right)_{\text{surroundings}} \right]$$

where  $S_{\text{gen}} \geq 0$

The implication here is that  $S_{\text{surround}}$  always increases for any real operation.  $\Delta S_{\text{system}}$  may actually decrease, but increase in entropy of surroundings more than makes up for it, always.

Example:



① System:  $300-200 \text{ KPa}$  } Superheat  $A-4$   $\rightarrow$   $S_1 = 7.7722 \frac{\text{KJ}}{\text{kg-K}}$   
 $160^\circ \rightarrow 320^\circ$   $U_1 = 0.651 \text{ m}^3/\text{kg}$   
 $h_1 = 2587.1 \frac{\text{KJ}}{\text{kg}}$   
 $h_2 = 2782.3 \frac{\text{KJ}}{\text{kg}}$

② System:  $300 \text{ KPa}$  } Superheat  $A-4$   $\rightarrow$   $S_2 = 8.522 \frac{\text{KJ}}{\text{kg-K}}$   
 $U_2 = 2U_1 = 1.302 \text{ m}^3/\text{kg}$   $h_2 = 3648.2 \frac{\text{KJ}}{\text{kg}}$

$$\Delta S_{\text{sys}} = S_2 - S_1 = 1.394 \frac{\text{KJ}}{\text{kg-K}}$$

Surroundings:  $\int_1^2 \left( \frac{\delta Q}{T} \right)_{\text{surr}} = ? \quad \frac{Q_{\text{surr}}}{T_{\text{surr}}} \Rightarrow Q = h_2 - h_1 \quad (\text{const } P)$   
 $= 865.9 \frac{\text{KJ}}{\text{kg}}$

But in this case, heat flow is referenced to surroundings, not system.

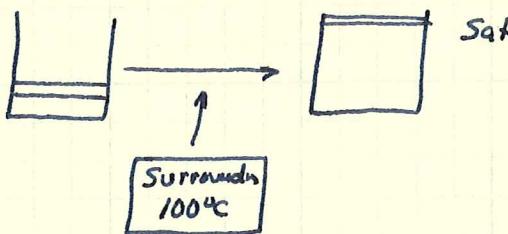
$$\frac{Q_{\text{surr}}}{T_{\text{surr}}} = - \frac{865.9 \frac{\text{KJ}}{\text{kg}}}{1000 \text{ K}} = -0.8659 \frac{\text{KJ}}{\text{kg-1K}}$$

$$S_{\text{gen}} = 1.394 - 0.8659 = 0.5281 \frac{\text{KJ}}{\text{kg-K}}$$

⇒ Some comments with respect to this —

So we made entropy... process cannot be reversed... Why?

Suppose Sat Liquid  
100°C  
1 kg.



Sat. Vapor - 100°C.

Reversible? Yes!  $S_{gen} = 0$   
probably

$$\text{Check quickly } \Delta S_{\text{system}} = S_g - S_f = 6.0480 \text{ kJ/kg-K}$$

$$Q = h_2 - h_1 = h_g - h_f = 2257 \text{ kJ/kg}$$

$$S_{gen} = \Delta S_{\text{sys}} + \left[ \int_1^2 \frac{\delta Q}{T} \right]_{\text{surroundings}} = 6.0480 + \frac{(-2257 \text{ kJ/kg})}{373.15 \text{ K}} = 0$$

Here, " $\rightarrow$ "  
because  $Q$  is ref'd  
to surroundings.

May 21 Got through increase in entropy principle: Three forms (for surroundings at constant temperature).

Closed System:  $\sigma_{\text{in booth}}$

$$S_{\text{gen}} = m(s_2 - s_1) + \frac{Q_{\text{surr}}}{T_{\text{surr}}} \geq 0$$

$$\frac{\text{System}}{- \int_1^2 \left( \frac{\delta Q}{T} \right)_{\text{sys body}}}$$

Steady Flow:

$$\dot{S}_{\text{gen}} = \sum m_i s_e - \sum m_i s_i + \frac{Q_{\text{surr}}}{T_{\text{surr}}} \geq 0$$

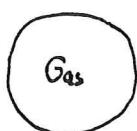
$$- \frac{\dot{Q}_{\text{sys}}}{T_{\text{sys}}}$$

$$\text{Unsteady (Fixed End states)} \quad S_{\text{gen}} = \sum m_e s_e - \sum m_i s_i + (m_2 s_2 - m_1 s_1)_{\text{cv}} \geq 0$$

$\Rightarrow$  Key point:  $Q_{\text{surr}}$  is referenced to surroundings so  $\frac{Q_{\text{surr}}}{T_{\text{surr}}}$  is positive.

One way to think of entropy is as a measure of the disorganization of energy.

Case 1:



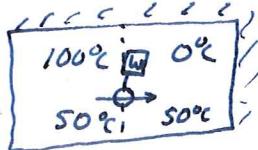
$\rightarrow$  Moving: Has kinetic energy.

$\Rightarrow$  On top of random motion, we have an organized motion. Could bring the gas to a rest and recover work  $\equiv$  organized energy that we can use elsewhere.

Case 2:



KE from case 1 is converted to internal energy



$\Rightarrow$  Now energy is random and disorganized. Cannot directly convert it to organized energy  $\Rightarrow$  Work.

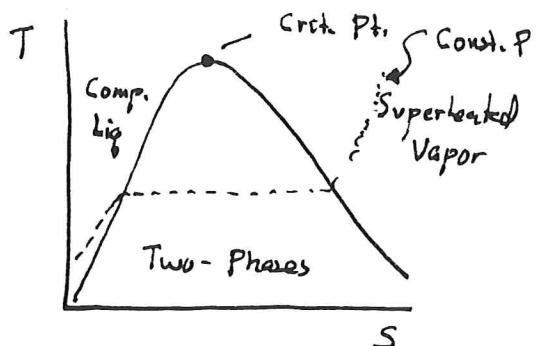
$\Rightarrow$   $w = \text{flow of organized energy} \xrightarrow{100\%} \text{Heat} \equiv \text{Disorganized Energy}$

$Q = \text{flow of disorganized energy} \xrightarrow{<100\%} \text{Work.}$

$\Rightarrow$  Its a one way street, and the increase in S principle expresses this.

Entropy Data: Steam, R-12, R-134a, etc. Read it off the chart.

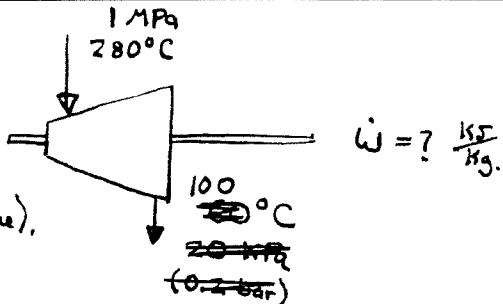
The T-S diagram:



Advantage: Constant entropy processes (isentropic processes) are represented as straight lines.

Simple example:

### Steam Turbine



You've gone to the catalog, and ordered a reversible, adiabatic turbine. (Top of the line).

How do I solve this?  $-\frac{\dot{\omega}}{m} = h_e - h_i$

→ know state 1. But only know  $T_2$  in could be any pressure.

→ Go back to increase in entropy principle for flow system (one inlet ≠ outlet).

$$S_{gen} = m \left( \frac{S_e - S_i}{T_e} \right) + \frac{\dot{Q}_{surr}}{T_{surr}} \geq 0$$

- ① What does reversible mean?  $S_{gen} = 0$  ! }  $\Rightarrow S_e - S_i$
- ② How about adiabatic?  $\dot{Q}_{surr} = 0$  }

So for a reversible, adiabatic process,  $S_i = S_f$ .

(i)  $\left. \begin{array}{l} 1 \text{ MPa} \\ 280^\circ\text{C} \end{array} \right\} \xrightarrow{\text{A-4}} \left\{ \begin{array}{l} h_i = 3008.2 \text{ kJ/kg} \\ s_i = 7.0465 \text{ kJ/kg-K} \end{array} \right.$

(e)  $T_e = 100^\circ\text{C}$        $s_e = s_i = 7.0465 \text{ kJ/kg-K}$       Two-Phase or superheated?

$$\text{A-2 } 100^\circ\text{C} \quad \frac{s_f}{1.3069} \quad \frac{s_g}{7.3549}$$

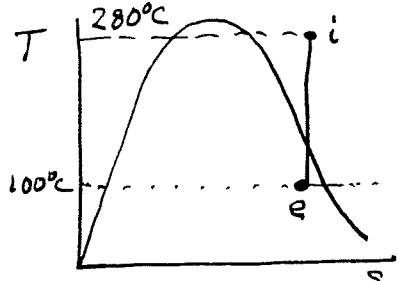
⇒ Since we are between, we are in two-phase.  
⇒ If  $s_e > s_g \rightarrow$  superheat.

$$\text{Get } x_e = \frac{s_e - s_f}{s_g - s_f} = \frac{7.0465 - 1.3069}{7.3549 - 1.3069} = 0.949$$

$$\text{Get } h_e = h_f + x_e h_{fg} = 419.04 + (0.949)(2257) = 2561 \text{ kJ/kg}$$

$$-\frac{\dot{\omega}}{m} = h_e - h_i = 2561 - 3008.2 \text{ kJ/kg} \Rightarrow \boxed{\frac{\dot{\omega}}{m} = 447.25 \text{ kJ/kg}}$$

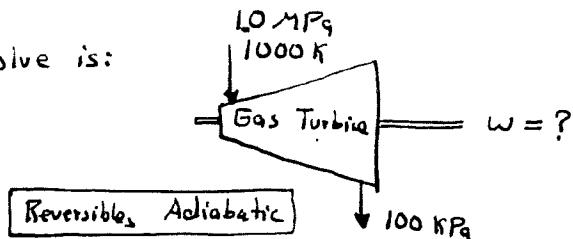
On the T-s diagram, what would this look like? Where would we start? end?



May 13

Gases are a little bit trickier. Need to go back to fundamentals.

For gases, the basic problem we wish to solve is:



If we knew  $T_e$ , then we could solve  $\Rightarrow -\frac{\dot{w}}{\dot{m}} = h_e - h_i$

$\Rightarrow$  Don't have tabular entropy data for air like we do for steam (do have  $S^\circ$ , but this isn't directly entropy).

Start with definition of entropy:  $\Delta S = \int_1^2 \left( \frac{\delta Q}{T} \right)_{rev}$  or  $dS = \left( \frac{\delta Q}{T} \right)_{rev}$

$$\text{So: } (\delta Q)_{rev} = T dS$$

Now first law, written specifically for case of reversible process:

$$(\delta Q)_{rev} - (\delta W)_{rev} = dU$$

But for a reversible path:  $(\delta W)_{rev} = PdV \dots$  so substituting for  $\delta Q$ ,  $\delta W$ .

$$TdS - PdV = dU \quad \boxed{dU = TdS - PdV} \quad \text{One of the Gibbs Relations.}$$

$$\text{or } \boxed{TdS = dU + PdV}$$

Can get the other from the definition of enthalpy.  $H = U + PV$

$\Rightarrow$  Can you figure out what the correct substitution into this equation would lead to?  $TdS = dH - PdV$ ?

Chain rule:  $dH = dU + PdV + VdP$

$$\Rightarrow dH - VdP = dU + PdV \quad \boxed{\text{Subst.}}$$

$$\boxed{TdS = dH - VdP} \Rightarrow \text{Second Gibbs Relation.}$$

~~Second~~ A major utility of these equations is that they allow ~~s~~ data to be developed from properties that are easier to measure.

Work with the ~~2nd~~ equation:  $ds = \frac{dH}{T} - \frac{V}{T}dP$   $ds = \frac{dh}{T} - \frac{v}{T}dp$

So far, we have made no assumptions in model.

Assume ideal gas:  $Pv = RT$ ,  $\frac{U}{T} = \frac{R}{P}$

$$ds = \frac{dh}{T} - \frac{R}{P}dP \quad \text{Integrate:}$$

$$\int_i^e ds = \int_i^e \frac{dh}{T} - \int_i^e \frac{R}{P}dP$$

Ask about each term

$$S_e - S_i = \int_i^e \frac{dh}{T} - R \ln \left[ \frac{P_e}{P_i} \right]$$

Cannot directly integrate because  $T$  varies with  $h$ .

$$\text{but } dh = C_p dT$$

$$S_e - S_i = \int_i^e \frac{C_p}{T} dT - R \ln \left[ \frac{P_e}{P_i} \right]$$

Key point. While  $U, h$  depend only on  $T$ , not on  $P$  for ideal gases,  $S = S(T, P)$ .

For constant  $C_p$  (inert gases, or real gases over small  $\Delta T$ ) can integrate this:

$$S_e - S_i = C_p \ln \left[ \frac{T_e}{T_i} \right] - R \ln \left[ \frac{P_e}{P_i} \right] \quad \text{- Ideal gas, const } C_p.$$

Conceptually, can you see how you would solve our original turbine problem?

$\Rightarrow$  Know ① Pressure Ratio }  $\Rightarrow T_e$   
②  $\Delta S = 0$

First Law

But if  $C_p$  varies with  $T$ , then we need to do the integral from  $C_p$  vs  $T$  data.

$\Rightarrow$  The integral depends only on  $T$ , so can define:

$$S^o(T) = \int_0^T \frac{C_p}{T} dT \quad \text{performed at 1 atm.}$$

$\Rightarrow$  This is tabulated in Table A-16.

$\Rightarrow$  So for variable  $C_p \Rightarrow$

$$S_e - S_i = S_e^o - S_i^o - R \ln \left[ \frac{P_e}{P_i} \right]$$

$$S_e - S_i = \int_i^e \frac{C_p}{T} dT - R \ln \left[ \frac{P_e}{P_i} \right]$$

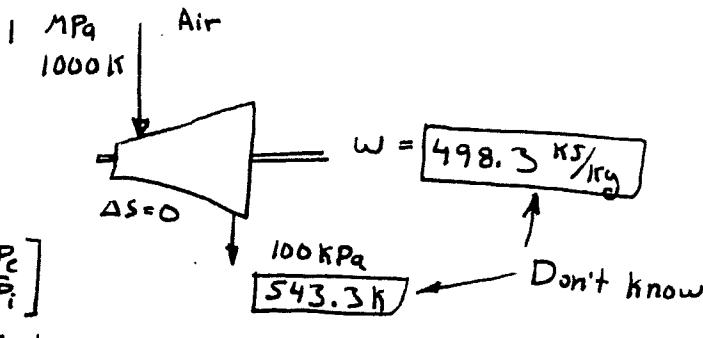
↳ Expresses  $T$       ↳ Expresses  $P$

$$\int_0^e \frac{C_p}{T} dT - \int_0^i \frac{C_p}{T} dT - \left[ R \ln \frac{P_e}{1\text{atm}} - R \ln \frac{P_i}{1\text{atm}} \right]$$

$S_e$        $S_i$

(66) (65)

Now one more trick:



Solved by  $\Delta S = 0 = S_e^0 - S_i^0 - R \ln \left[ \frac{P_e}{P_i} \right]$   
 $\rightarrow A-16 \rightarrow T_e, h_e.$

For isentropic:  $\exp \left[ \frac{S_e^0 - S_i^0}{R} \right] = \frac{P_e}{P_i} = \frac{\exp \left[ \frac{S_e^0}{R} \right]}{\exp \left[ \frac{S_i^0}{R} \right]}$

Define  $P_r = \exp \left[ \frac{S_e^0}{R} \right]$  → In table A-16.  
 (relative pressure) not reduced pressure  $P_r = P/P_c$   
 $\Rightarrow$  For isentropic processes, ( $\Delta S = 0$ )

Depends only on  $T_e$

$$\frac{P_e}{P_i} = \frac{P_{r,e}}{P_{r,i}}$$

Depends only  
on  $T_e$  in  
Table A-16.

For our problem: What do we know? All but  $P_{r,e} \Rightarrow$

$$P_{r,e} = P_{r,i} \left[ \frac{P_e}{P_i} \right] = 114 \left[ \frac{0.1}{1.0} \right] = 11.4 \xrightarrow{\text{A-16}} T_e = 543.3 \text{ K}$$

$\xrightarrow{\text{A-16, 1000 K}}$

$$h_e = 547.75 \text{ kJ/kg}$$

$$-\frac{\dot{m}}{\dot{m}} \dot{w} = h_e - h_i \quad \xrightarrow{\dot{w} = 498.3 \text{ m/s/kg}}$$

⇒ Same answer as before!

There is a reduced volume function that serves the same purpose.

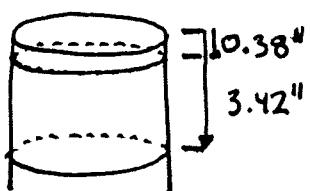
$$\frac{U_z}{V_i} = \frac{U_{r,z}}{U_{r,i}}$$

- For  $\Delta S = 0$
- $U_r$ 's are a function only of  $T_e$  can be found in Table A-16.

Example: Automobile Engine: 5.0 liter displacement V-8. }  
 Compression Ratio = 9 }

$$V_{max} - V_{min} = \frac{0.005 \text{ m}^3}{8} \quad \frac{V_{max}}{V_{min}} = 9 \quad 8V_{min} = \frac{0.005}{8} \Rightarrow V_{min} = 7.81 \times 10^{-5} \text{ m}^3$$

$$V_{max} = 7.03 \times 10^{-4} \text{ m}^3$$



$$V_1 = 7.81 \times 10^{-5} \text{ m}^3 \quad P_1 = 7.0 \text{ MPa}$$

$$V_2 = 7.03 \times 10^{-4} \text{ m}^3 \quad T_1 = 2000 \text{ K.}$$

⇒ Find how much work you get out of one expansion.

May 4

- Assume its reversible,
- Time fast enough so  $Q=0$
- Pretend gas has same props as air.

Preliminaries. Mass in cylinder:

$$m = \frac{PV}{RT} = \frac{7000 \text{ kPa}}{2000 \text{ K}} \left| \frac{7.81 \times 10^{-5} \text{ m}^3}{8.314 \text{ kPa} \cdot \text{m}^3} \right| \text{ kmole-1} \left| \frac{28.9 \text{ kg}}{\text{kmole}} \right|$$

$$m = 9.50 \times 10^{-4} \text{ kg}$$

What is first law?

$$Q-W = m(u_2 - u_1) \Rightarrow W = m(u_1 - u_2) \quad \text{① } 2000\text{K} \rightarrow \text{u}_1 = 1678.7 \text{ kJ/kg.}$$

Since  $\Delta S = 0$ , Have available:

$$\frac{P_2}{P_1} = \frac{P_{r,2}}{P_{r,1}} \quad \frac{V_2}{V_1} = \frac{V_{r,2}}{V_{r,1}} \quad \begin{cases} \Rightarrow \text{Both are valid.} \\ \Rightarrow \text{Which has the greater utility now?} \end{cases}$$

$$V_{r,2} = V_{r,1} \left[ \frac{V_2}{V_1} \right] = 2.776 \left[ \frac{7.03 \times 10^{-4}}{7.81 \times 10^{-5}} \right] = 24.99 \xrightarrow{\text{A-16}} T_2 = 1002.5 \text{ K} \quad \text{u}_2 = 761.1 \text{ kJ/kg}$$

$$W = m(u_1 - u_2) = [9.50 \times 10^{-4} \text{ kg}] [1678.7 - 761.1 \text{ kJ/kg}] = \underline{\underline{0.87 \text{ kJ}}}$$

How would you find the final pressure?

One Way

$$\frac{P_1 V_1}{T_1} = \frac{P_2 V_2}{T_2} \quad P_2 = P_1 \left[ \frac{V_1}{V_2} \right] \left[ \frac{T_2}{T_1} \right] = (7000 \text{ kPa}) \left[ \frac{7.81 \times 10^{-5}}{7.03 \times 10^{-4}} \right] \left[ \frac{1002.5}{2000} \right] = 390 \text{ kPa.}$$

Alternate way:

$$\frac{P_2}{P_1} = \frac{P_{r,2}}{P_{r,1}} \Rightarrow P_2 = P_1 \left[ \frac{P_{r,2}}{P_{r,1}} \right] = (7000 \text{ kPa}) \left[ \frac{115.2}{2068} \right] = 390 \text{ kPa.}$$

Next topic:

This describes how we handle variable  $C_p$ ,  $\Delta S = 0$  problems. Now we want to get into constant  $C_p$ .

$$S_2 - S_1 = C_p \ln \left[ \frac{T_2}{T_1} \right] - R \ln \left[ \frac{P_2}{P_1} \right] \quad \text{so for } \Delta S = 0$$

$$\cancel{e^{P_1}} \quad \ln \left[ \frac{T_2}{T_1} \right] = \frac{R}{C_p} \ln \left[ \frac{P_2}{P_1} \right] = \ln \left[ \left( \frac{P_2}{P_1} \right)^{\frac{R}{C_p}} \right]$$

$$\text{Take exp of both sides} \Rightarrow \frac{T_2}{T_1} = \left( \frac{P_2}{P_1} \right)^{\frac{R}{C_p}}$$

$$R = C_p - C_v \text{ for ideal gas.}$$

$$\frac{R}{C_p} = \frac{C_p - C_v}{C_p} = 1 - \frac{C_v}{C_p} = 1 - \frac{1}{k} = \frac{k-1}{k} \Rightarrow$$

$$\text{but } k = \frac{C_p}{C_v}$$

↑ Similar Derivation

$$\frac{T_2}{T_1} = \left( \frac{P_2}{P_1} \right)^{\frac{k-1}{k}}$$

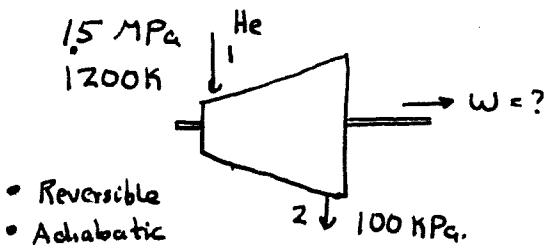
(6)  
(7)

Note that the  $P-V$  relationship says that  $P_1 V_1^k = P_2 V_2^k = \frac{PV}{\text{For all points between.}}$

So this is where the Polytropic Processes we discussed way back come in:

For isentropic processes, ideal gas, const.  $C_p \Rightarrow PV^k = \text{const}$   
Air,  $k=1.4$  at room T.

Example: Nuclear power plant operating on helium cycle.



$$\text{First Law} \quad -\frac{\dot{w}}{\dot{m}} = h_2 - h_1$$

No helium tables... why?

$$\Rightarrow \frac{\dot{w}}{\dot{m}} = C_p(T_1 - T_2)$$

Isentropic relations say:

$$\frac{T_2}{T_1} = \left[ \frac{P_2}{P_1} \right]^{\frac{k-1}{k}}$$

$$\bar{C}_p = \frac{5}{2} R$$

$$\bar{C}_v = \bar{C}_p - R = \frac{3}{2} R$$

$$T_2 = T_1 \left[ \frac{P_2}{P_1} \right]^{\frac{k-1}{k}} = [1200 \text{ K}] \left[ \frac{0.1 \text{ MPa}}{1.5 \text{ MPa}} \right]^{\frac{0.667}{1.667}} = 406.07 \text{ K} \Rightarrow \frac{\bar{C}_p}{\bar{C}_v} = \frac{5/2 R}{3/2 R} = 1.667$$

All inert gases.

$$\frac{\dot{w}}{\dot{m}} = \frac{5}{2} \left| \frac{8.314 \text{ kJ}}{\text{Kmole-K}} \right| \frac{1200 - 406.07 \text{ K}}{4 \text{ kg}} = 4,125.46 \frac{\text{kJ}}{\text{kg}}$$

Now, same example but its an internally reversible turbines with a heat loss of 1000 kJ/kg.

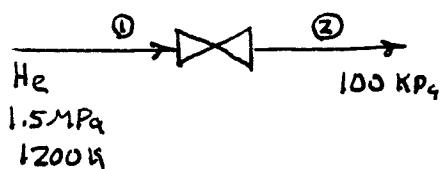
First Law:

$$\frac{\dot{w}}{\dot{m}} = \frac{\dot{Q}}{\dot{m}} + C_p(T_1 - T_2)$$

Same problem,  
but how to  
find  $T_2$ ?

Now another example:

Throttle



What is  $T_2$ ?

$$T_2 = 1200 \text{ K},$$

How about non-ideal?  $\Rightarrow h_1 = h_2$ ,  
but  $T_1 \neq T_2$ .

Find entropy generation for this process two ways

- ① Normal Way
- ② from definition.

Ideal gas: Constant  $C_p, C_v \Rightarrow \frac{s_2 - s_1}{T_2 - T_1}$

$$(\dot{S}_{\text{gen}})_{\text{sys}} = \dot{m}(s_2 - s_1) - \int_1^2 \left( \frac{\partial S}{\partial T} \right)_{\text{sys}} dT$$

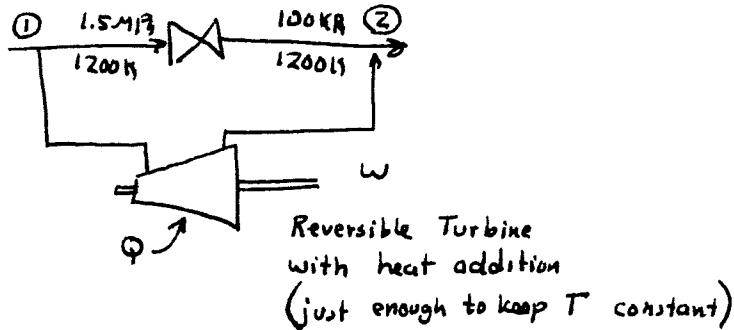
$$s_2 - s_1 = C_p \ln \left[ \frac{T_2}{T_1} \right] - R \ln \left[ \frac{P_2}{P_1} \right] = - \frac{8.314 \text{ kJ}}{\text{Kmole-K}} \left| \frac{100}{1.5} \right| \ln \left[ \frac{0.1}{1.5} \right] = 5.63 \frac{\text{kJ}}{\text{Kmole-K}}$$

May 6  
DTP on  
 $N_{Co}$ ,  $N_T$

From definition:

$$\Delta S = \int_1^2 \left( \frac{\delta Q}{T} \right)_{rev.}$$

Need to construct a reversible path between ① and ②, which the throttle isn't!  
This will be a little bit different than last time:



First Law:  $\dot{Q} = \dot{W}$  ( $\Delta T = 0$ ) [Note that this isn't a cycle... in a cycle  $\dot{Q}_{in} > \dot{W}_{net}$ ]

How to get the work... only thing we have is  $P_1, P_2$ .

Within the turbine, for each parcel of gas  $\frac{P_i V}{T} = \text{const.}$  out since  $T = \text{const}$

$$\Rightarrow PV = \text{const.} \quad \dot{W} = \int P dV = \int \frac{P_i V_i}{V} dV = P_i V_i \ln \left[ \frac{V_2}{V_1} \right]$$

$$\text{but since } PV = RT \Rightarrow \dot{W} = RT_i \ln \left[ \frac{P_1}{P_2} \right] = \frac{8.314 \text{ kJ}}{\text{kmole-K}} \left| \frac{1200 \text{ K}}{100 \text{ K}} \right| \ln \left[ \frac{1.5}{0.1} \right] \\ \frac{V_2}{V_1} = \frac{P_1}{P_2} \quad \dot{W} = 6754 \text{ kJ/kg} = \dot{Q}$$

$$\Delta S = \int_1^2 \left( \frac{\delta Q}{T} \right)_{rev.} = \frac{\dot{Q}}{T} = \frac{6754 \text{ kJ/kg}}{1200 \text{ K}} = \boxed{5.63 \text{ kJ/kg-K}}$$

Same as before

### Review of Major Relations:

Definition of entropy:

$$\Delta S = \int_1^2 \left( \frac{\delta Q}{T} \right)_{rev.}$$

→ Can be used directly if path is internally reversible

→ If process between 1 and 2 is not reversible - must construct alternate reversible path to evaluate.

From definition

$$(S_{gen})_{\text{system}} = (S_2 - S_1)_{\text{system}} - \int_1^2 \left( \frac{\delta Q}{T} \right)_{\text{system.}}$$

- ① Note that you have to use correct sign on  $Q$
- ② Other forms include steady flow & unsteady process.

$$(S_{gen})_{\text{sys + surroundings}} = (S_2 - S_1)_{\text{system}} + \frac{Q_{\text{surr}}}{T_{\text{surr}}}$$

→ Assumes surroundings are at constant  $T_0$

For reversible ( $S_{gen}=0$ ), adiabatic ( $Q=0$ )  
Processes -  $\Rightarrow$  This converts to  $S_2 = S_1$

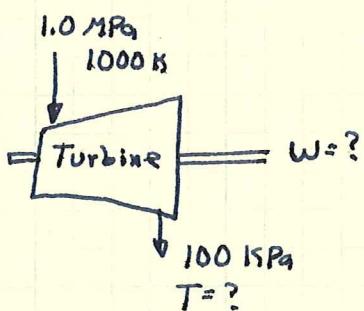
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PCM gave May 13.

Exam on May 15.

May 17

Brief Summary:



$$\text{First Law: } \delta Q - \delta W = dU$$

$$\text{Assume reversible paths } (\delta Q)_{\text{rev}} - (\delta W)_{\text{rev}} = dU$$

$$(\delta W)_{\text{rev}} = PdV \quad dS = \frac{(\delta Q)_{\text{rev}}}{T}$$

$$\Rightarrow TdS = PdV = dU \quad \text{or} \quad TdS = \underbrace{dU + PdV}_{H} \stackrel{\text{Gibbs, 1}}{=}$$

$$H = U + PV \quad \text{so} \quad dH = dU + PdV + VdP \quad \text{or} \quad dH - VdP = dU + PdV$$

$$\Rightarrow TdS = dH - VdP \quad \text{Gibbs 2.}$$

$$dS = \frac{dh}{T} - \frac{V}{T}dp \quad \text{Get } S \text{ from simple prop.}$$

$$\text{Integrate: } \Delta S = dh = C_p dT, \quad \frac{T_2}{T_1} = \frac{R}{P}$$

$$\Delta S = \int_{T_1}^{T_2} \frac{C_p}{T} dT - \int_{P_1}^{P_2} \frac{R}{P} dP = \int_{T_1}^{T_2} \frac{C_p}{T} dT - R \ln \frac{P_2}{P_1}$$

Now can go a couple of ways:

①  $C_p = \text{constant}$  (He, Air when  $\Delta T$  isn't too big in I will make it clear)

$$\Delta S = C_p \ln \frac{T_2}{T_1} - R \ln \frac{P_2}{P_1}$$

$$\text{If turbine is reversible, adiabatic} \quad [S_{\text{gen}} = (\Delta S)_{\text{sys}} + \frac{Q_{\text{gen}}}{T_{\text{gen}}}] \Rightarrow \Delta S = 0$$

$$\ln \left[ \frac{T_2}{T_1} \right] = \frac{R}{C_p} \ln \left[ \frac{P_2}{P_1} \right] \quad \text{or} \quad \frac{T_2}{T_1} = \left[ \frac{P_2}{P_1} \right]^{\frac{R}{C_p}} \quad R = C_p - C_v$$

$$\frac{T_2}{T_1} = \left[ \frac{P_2}{P_1} \right]^{\frac{K-1}{K}}$$

$$\text{where } K = \frac{C_p}{C_v} \approx 1.4 \text{ for air} \\ = 1.67 \text{ for inert,} \\ \approx 1.3 \text{ for hot air.}$$

Similar ~~Relation~~  
Derivations.

$$\frac{P_2}{P_1} = \left( \frac{V_1}{V_2} \right)^K$$

$$\frac{T_2}{T_1} = \left[ \frac{P_2}{P_1} \right]^{\frac{K-1}{K}}$$

②  $C_p = \text{variable}$ :

Get from Table A-17

$$\Delta S = \underbrace{S_2^o - S_1^o}_{T\text{-only}} - R \ln \frac{P_2}{P_1} \quad \underbrace{\text{P-only}}$$

$$\int_0^{T_2} \frac{C_p}{T} dT = \int_0^{T_1} \frac{C_p}{T} dT$$

Note that relation says that for

- Ideal gas
  - Const  $C_p, C_v$
  - $\Delta S > 0$  (max work)
- $\Rightarrow P_1 V_1^K = P_2 V_2^K = \text{const}$
- $\Rightarrow$  Polytropic Process, where  $K = C_p/C_v$

$$\text{For turbine where } \Delta S = 0; \exp\left[\frac{S_2^o - S_1^o}{R}\right] = \frac{P_2}{P_1}$$

$$\Rightarrow \frac{\exp\left[\frac{S_2^o}{R}\right]}{\exp\left[\frac{S_1^o}{R}\right]} = \frac{P_2}{P_1}$$

Depends only  
on  $T$

$$\text{Define } \exp\left[\frac{S_2^o}{R}\right] = P_{r,2} \Rightarrow \text{Table A-17}$$

$\Rightarrow$  For  $\Delta S = 0$

$$\frac{P_{r,2}}{P_{r,1}} = \frac{P_2}{P_1}$$

$\Rightarrow$  Gives you way to  
solve turbine problem.

• Know  $P_2/P_1, T_1 \Rightarrow P_{r,2}$

$\Rightarrow$  Equation  $P_{r,2} \Rightarrow T_2$

$$\left. \begin{array}{l} \Rightarrow \\ \Rightarrow h_2 \end{array} \right\} \quad \boxed{h = h_1 - h_2}$$

Similar idea for volume:

$$\frac{V_{r,2}}{V_{r,1}} = \frac{V_2}{V_1}$$

For  $\Delta S = 0$

$V_r = f(T)$  only - Table A-17

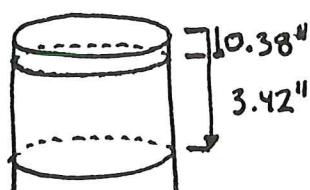
Example: Automobile Engine:

5.0 liter displacement V-8.  
Compression Ratio = 9

$$V_{\max} - V_{\min} = \frac{0.005 \text{ m}^3}{8} \quad \begin{matrix} 1 & 4'' & 1 \end{matrix}$$

$$\frac{V_{\max}}{V_{\min}} = 9 \quad 8V_{\min} = \frac{0.005}{8} \Rightarrow V_{\min} = 7.81 \times 10^{-5} \text{ m}^3$$

$$V_{\max} = 7.03 \times 10^{-4} \text{ m}^3$$



$$\left. \begin{array}{l} V_1 = 7.81 \times 10^{-5} \text{ m}^3 \\ V_2 = 7.03 \times 10^{-4} \text{ m}^3 \end{array} \right\} \quad \begin{matrix} P_1 = 7.0 \text{ MPa} \\ T_1 = 2000 \text{ K.} \end{matrix}$$

$\Rightarrow$  Find how much work you get  
out of one expansion.

- Assume it's reversible,
  - Time fast enough so  $Q=0$
  - Pretend gas has same props as air.
- Preliminaries. Mass in cylinder:

$$m = \frac{PV}{RT} = \frac{7000 \text{ kPa}}{2000 \text{ K}} \left| \frac{7.81 \times 10^{-5} \text{ m}^3}{8.314 \text{ kPa} \cdot \text{m}^3} \right| \text{ kmole-lf} \left| \frac{28.9 \text{ kg}}{\text{kmole}} \right|$$

$$m = 9.50 \times 10^{-4} \text{ kg}$$

What is first law?

$$Q-W = m(u_2 - u_1) \Rightarrow W = m(u_1 - u_2) \quad \text{① } 2000\text{K} \rightarrow \cancel{u_1} \quad u_2 = 1678.7 \text{ kJ/kg.}$$

Since  $\Delta S=0$ , Have available:

$$\frac{P_2}{P_1} = \frac{P_{r,2}}{P_{r,1}} \quad \frac{V_2}{V_1} = \frac{V_{r,2}}{V_{r,1}} \quad \begin{cases} \Rightarrow \text{Both are valid.} \\ \Rightarrow \text{Which has the greater utility now?} \end{cases}$$

$$V_{r,2} = V_{r,1} \left[ \frac{V_2}{V_1} \right] = 2.776 \left[ \frac{7.03 \times 10^{-4}}{7.81 \times 10^{-5}} \right] = 24.99 \xrightarrow{\text{A-16}} T_2 = 1002.5 \text{ K}$$

Can look up @ 2000K, Table A-16.

$$\cancel{u_1} \quad u_2 = 761.1 \text{ kJ/kg}$$

$$W = m(u_1 - u_2) = [9.50 \times 10^{-4} \text{ kg}] [1678.7 - 761.1 \text{ kJ/kg}] = \underline{\underline{0.87 \text{ kJ}}}$$

May 20 How would you find the final pressure?

One Way

$$\frac{P_1 V_1}{T_1} = \frac{P_2 V_2}{T_2} \quad P_2 = P_1 \left[ \frac{V_1}{V_2} \right] \left[ \frac{T_2}{T_1} \right] = (7000 \text{ kPa}) \left[ \frac{7.81 \times 10^{-5}}{7.03 \times 10^{-4}} \right] \left[ \frac{1002.5}{2000} \right] = 390 \text{ kPa.}$$

Alternate way:

$$\frac{P_2}{P_1} = \frac{P_{r,2}}{P_{r,1}} \Rightarrow P_2 = P_1 \left[ \frac{P_{r,2}}{P_{r,1}} \right] = (7000 \text{ kPa}) \left[ \frac{115.2}{2068} \right] = 390 \text{ kPa.}$$

~~A heat topic:~~

This describes how we handle variable  $C_p$ ,  $\Delta S=0$  problems. Now we want to get into constant  $C_p$ .

$$S_2 - S_1 = C_p \ln \left[ \frac{T_2}{T_1} \right] - R \ln \left[ \frac{P_2}{P_1} \right] \quad \text{so for } \Delta S=0$$

$$\cancel{C_p} \quad \ln \left[ \frac{T_2}{T_1} \right] = \frac{R}{C_p} \ln \left[ \frac{P_2}{P_1} \right] = \ln \left[ \left( \frac{P_2}{P_1} \right)^{\frac{R}{C_p}} \right]$$

$$\text{Take exp. of both sides} \Rightarrow \frac{T_2}{T_1} = \left( \frac{P_2}{P_1} \right)^{\frac{R}{C_p}}$$

$R = C_p - C_v$  for ideal gas.

$$\frac{R}{C_p} = \frac{C_p - C_v}{C_p} = 1 - \frac{C_v}{C_p}$$

$$\text{but } k = \frac{C_p}{C_v}$$

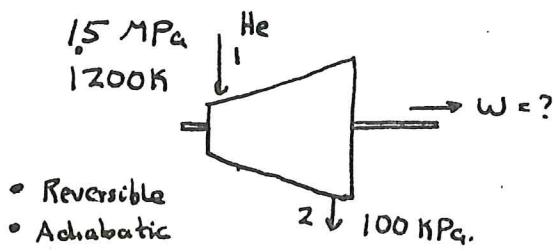
$$= 1 - \frac{1}{k} = \frac{k-1}{k} \Rightarrow$$

$$\frac{P_2}{P_1} \left( \frac{V_1}{V_2} \right)^k \quad \frac{T_2}{T_1} = \left( \frac{P_2}{P_1} \right)^{\frac{k-1}{k}}$$

↑ Similar Derivation

$$\frac{T_2}{T_1} = \left( \frac{P_2}{P_1} \right)^{\frac{k-1}{k}}$$

Example: Nuclear power plant operating on helium cycle.



$$\text{First Law} \quad -\frac{\dot{W}}{m} = h_2 - h_1$$

No helium tables... why?

$$\Rightarrow \frac{\dot{W}}{m} = C_p(T_1 - T_2)$$

ISENTROPIC relations say:

$$T_2 = T_1 \left[ \frac{P_2}{P_1} \right]^{\frac{K-1}{K}} = [1200 \text{ K}] \left[ \frac{0.1 \text{ MPa}}{1.5 \text{ MPa}} \right]^{\frac{0.667}{1.667}} = 406.07 \text{ K}$$

$$\frac{T_2}{T_1} = \left[ \frac{P_2}{P_1} \right]^{\frac{K-1}{K}} = \frac{C_p}{C_v} = \frac{5/2 R}{3/2 R} = 1.667$$

All inert gases.

$$\frac{\dot{W}}{m} = \frac{5}{2} \frac{2.08 \text{ kJ/K}}{t_{\text{inert}} - K} \frac{1200 - 406.07 \text{ K}}{4 \text{ kg}} = 4,125.46 \frac{\text{kJ}}{\text{kg}}$$

Now we can solve the turbine problem:

$$S_e - S_i = 0 = S_e^o - S_i^o - R \ln \frac{P_e}{P_i}$$

$$S_e^o = S_i^o + R \ln \frac{P_e}{P_i} = 2.96770 \frac{\text{KJ}}{\text{kg}\cdot\text{K}} + \frac{8.314 \text{ KJ}}{\text{kmole}\cdot\text{K}} \left| \frac{1 \text{ kmole}}{28.9 \text{ kg}} \right| \ln \left[ \frac{0.1 \text{ MPa}}{1.0 \text{ MPa}} \right]$$

$$S_e^o = 2.305 \frac{\text{KJ}}{\text{kg}\cdot\text{K}}$$

A-16

1000K →

$$T_e = 543.3 \text{ K}$$

$$h_e = 547.75 \text{ KJ/kg}$$

$$h_i = 1046.04 \text{ KJ/kg}$$

$$-\frac{\dot{W}}{\dot{m}} = h_e - h_i = 547.75 - 1046.04 \Rightarrow$$

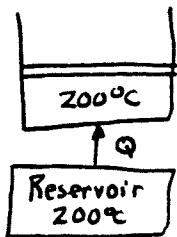
$$\frac{\dot{W}}{\dot{m}} = 498.3 \frac{\text{KJ}}{\text{kg}}$$

A slightly different derivation leads you to another useful expression:

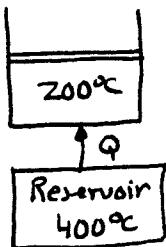
$$\Delta S =$$

Another concept: Internal Reversibility: Everything inside boundaries is reversible.

Isothermal Expansion

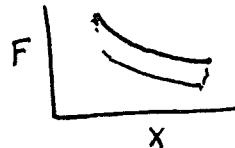
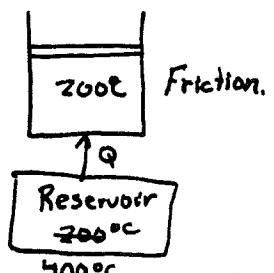


Totally Reversible



Internally Reversible

(Q across  $\Delta T$  is outside body)

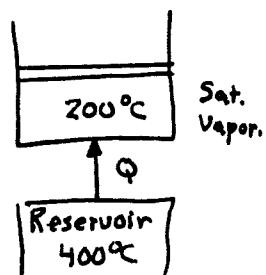


Less W out  
Beginning and  
end states  
same.  
 $\Rightarrow$  less Q in.

Also works

$$\Delta S = \int \left( \frac{\delta Q}{T} \right)_{\text{Internally Reversible}}$$

Let's you find Q.



Find

- Q Added until volume doubles.
- T stays constant.
- Internally Reversible.

$$\begin{aligned} P_i &= 1.554 \text{ MPa} & \text{KJ} \\ u_2 &= 2595.3 \text{ mJ} & \text{kg} \\ T_i &= 200^\circ\text{C} \\ S_i &= 6.4323 \text{ KJ} & \text{kg}\cdot\text{K} \end{aligned}$$

Find Q, W, entropy gen. for sys.  
entropy gen for universe.

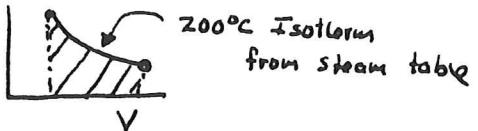
$$P_2 = 0.8441 \text{ MPa}$$

$$u_2 = 2628.6 \text{ KJ/kg}$$

$$s_2 = 6.794 \text{ KJ/kg}\cdot\text{K}$$

For  $Q, w$ ,  $\Rightarrow$  First Law looks like  $Q - w = m(u_2 - u_1)$   $w = \int P dV$ , but  $P$  varies.

One approach is to use brute force  $\Rightarrow P$



$\Rightarrow$  Get  $w_{rev}$  then  $Q$  from first law.

Alternate. Since process is internally reversible:  $\Delta S = \int \left( \frac{\delta Q}{T} \right)$

$$\text{Since } T = \text{constant: } T(s_2 - s_1) = Q = [200 + 273 K] [6.794 - 6.4323 \frac{KJ}{kg \cdot K}] = +171.1 \frac{KJ}{kg}$$

$$\text{Now use first law to get } w = Q - m(u_2 - u_1) = 171.1 \frac{KJ}{kg} - [2628.6 - 2595.3 \frac{KJ}{kg}] = 137.78 \frac{KJ}{kg}$$

$\Rightarrow$  This gives us an additional tool in our arsenal  $\Rightarrow$  If  $T$  is  
 $\Rightarrow$  If  $T$  is constant & process is internally reversible, can quickly  
get  $Q$ .  
 $\Rightarrow$  If  $P$  is constant, can quickly get  $w$ .  
 $\Rightarrow$  Whichever you get, use first law to find the other.

$$(S_{gen})_{\text{system}} = m(s_2 - s_1) - \int \left( \frac{\delta Q}{T} \right)_{\text{sys}}$$

$$= m(s_2 - s_1) - \frac{Q_{\text{sys}}}{T_{\text{sys}}} = 6.794 - 6.4323 \frac{KJ}{kg \cdot K} - \frac{171.1 \frac{KJ}{kg}}{473 K}$$

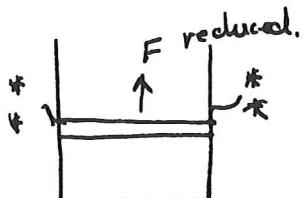
$$(S_{gen})_{\text{sys}} = 0$$

$$(S_{gen})_{\text{sys}} = m(s_2 - s_1) + \frac{Q_{\text{surr}}}{T_{\text{surr}}} = 6.794 - 6.4323 + \frac{(-171.1 \frac{KJ}{kg})}{473 K}$$

$$(S_{gen})_{\text{sys}} = 0.105 \frac{KJ}{kg \cdot K}$$

Since  $S$  was generated, what was the source of irreversibility?  
 $Q$  across  $\Delta T: \Rightarrow 400 \rightarrow 200^\circ C$ .

Now think about case where there is friction so that  $w < \int P dV$ .  
Beginning: End states the same:



$$F = 0.75 F_{\text{no friction}}$$

$$w = 0.75 w_{\text{Reversible}} = 0.75 [137.78 \frac{KJ}{kg}] = 103.34 \frac{KJ}{kg}.$$

$$\text{First Law: } Q = m(u_2 - u_1) + w = [2628.6 - 2595.3 \frac{KJ}{kg}] + 103.34 \frac{KJ}{kg}$$

$$Q = 136.64 \frac{KJ}{kg}.$$

$$(S_{gen})_{\text{system}} = m(s_2 - s_1) - \frac{Q_{\text{sys}}}{T_{\text{sys}}} = 6.794 - 6.4323 \frac{KJ}{kg \cdot K} - \frac{136.64}{473 K} = 0.073 \frac{KJ}{kg \cdot K}$$

From definition of  $S$ , and first law for reversible systems  $\Rightarrow$  Gibbs relations.

$$\begin{aligned} Tds &= du + PdV \\ Tds &= dh - VdP \end{aligned}$$

Gives us a way to find  $\Delta S$  from measurable props.  
 ① Develop  $S$  values in steam tables via numerical integration.

② For ideal gas - can numerically analytically integrate.

2<sup>nd</sup> Equation:

$$I. \quad (S_2 - S_1)_{sys} = \int_{T_1}^{T_2} \frac{dh}{T} - \int_{P_1}^{P_2} \frac{VdP}{T}$$

Variable  $C_p$ .

$$(S_2 - S_1)_{sys} = S_2^o - S_1^o - R \ln \left[ \frac{P_2}{P_1} \right]$$

[- From gas tables,  
- Function of T only.]

For isentropic processes:  $(\Delta S)_{sys} = 0$

$$\frac{P_2}{P_1} = \frac{P_{r,2}}{P_{r,1}} \quad \frac{V_2}{V_1} = \frac{V_{r,2}}{V_{r,1}}$$

- Function of T only  
- Found in Table A-16.

$\Rightarrow$  For other gases that do not list  $P_r, V_r$  values (e.g.  $O_2, N_2$ ), need to use  $S^o$

$$(S_2 - S_1)_{sys} = \int_{T_1}^{T_2} \frac{C_p}{T} dT - R \ln \left[ \frac{P_2}{P_1} \right]$$

$\Rightarrow$  For ideal gas -  $S$  depends on both  $P, T$ .

Constant  $C_p$  (inerts & small  $\Delta T$ )  
 $\Rightarrow$  Integrate directly

$$(S_2 - S_1)_{sys} = C_p \ln \left[ \frac{T_2}{T_1} \right] - R \ln \left[ \frac{P_2}{P_1} \right]$$

for convenience in some problems, a slight transformation yields.

$$(S_2 - S_1)_{sys} = C_v \ln \left[ \frac{T_2}{T_1} \right] + R \ln \left[ \frac{V_2}{V_1} \right]$$

For isentropic processes  $(\Delta S)_{sys} = 0$

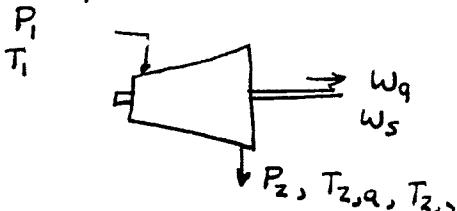
$$\frac{T_2}{T_1} = \left[ \frac{P_2}{P_1} \right]^{\frac{K-1}{K}}$$

$$\frac{T_2}{T_1} = \left[ \frac{V_1}{V_2} \right]^{K-1}$$

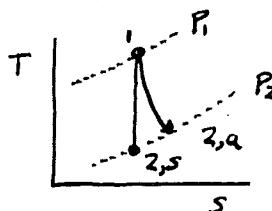
$$\frac{P_2}{P_1} = \left[ \frac{V_1}{V_2} \right]^K$$

$K = C_p/C_v$

Compressor and Turbine efficiency:



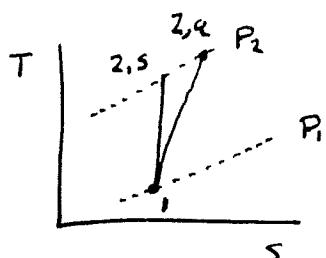
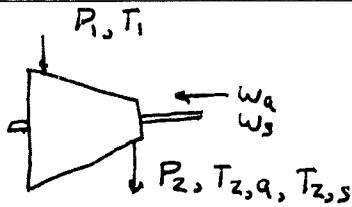
$$\eta_T = \frac{\text{actual } w}{\text{isentropic } w} = \frac{h_1 - h_{2,a}}{h_1 - h_{2,s}} < 1$$



Real  $\eta \Rightarrow$  Think of some useful work being lost to friction

$\Rightarrow$  Less work comes out because exiting fluid is hotter.

Compressor:



Defined as inverse

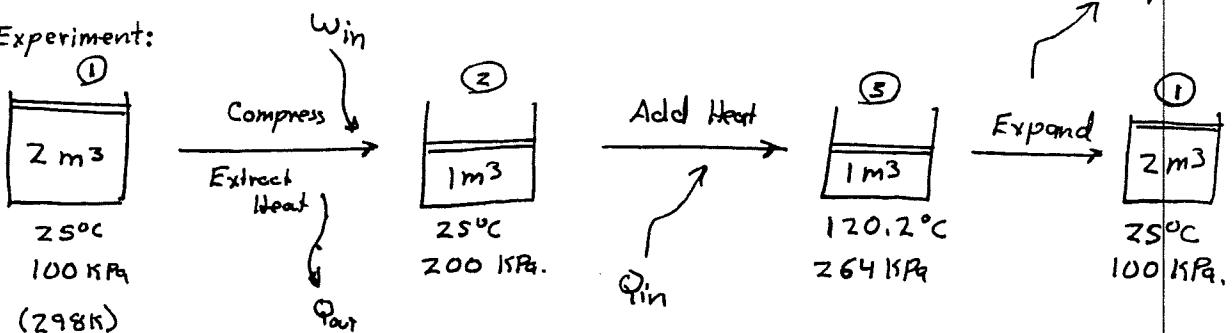
$$\eta_c = \frac{\text{isentropic work needed to do comp}}{\text{actual work needed.}}$$

$$= \frac{h_{2,s} - h_1}{h_{2,q} - h_1} < 1$$

Think of some win going to increase  $P_2$   
and some going to friction  
→ Exhaust gas is hotter ⇒ need more works

### Cycle Analysis - An Introduction:

Thought Experiment:



First Law: Neglect KE, PE.

$$Q - W = \Delta U \Rightarrow \text{For entire process, start and end at same place.} \\ \Rightarrow \Delta U = 0.$$

$$\Rightarrow \sum Q_{\text{cycle}} = \sum W_{\text{cycle}}.$$

$$4(Q_{1-2} + Q_{2-3}) = 4(W_{1-2} + W_{2-3}) \\ - + - +$$

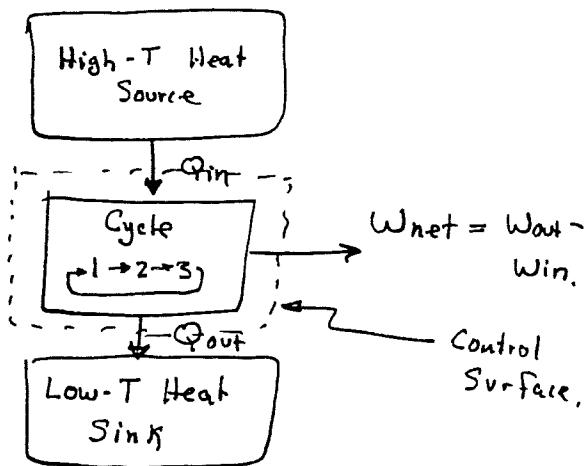
numerically

Last question: ~~Is~~ Which is larger:  $W_{\text{in}}$  ( $W_{1-2}$ ) or  $W_{\text{out}}$  ( $W_{2-3}$ )?

$$W = \int P dV \quad \begin{array}{l} \text{for } W_{\text{in}}, \quad P = 100 \rightarrow 200 \\ \text{for } W_{\text{out}}, \quad P = 264 \rightarrow 100 \end{array}$$

$\Rightarrow$  Average pressure is higher for  $W_{\text{out}}$ :  $\Rightarrow$  Net work for cycle is positive.

$\Rightarrow$  ~~For~~ Simply process to this:



Think of your automobile as a power cycle:

- Heat Engine: Turns heat into work.

- First Law: (Convention for cycles).

$$Q_m = Q_{\text{out}} + W_{\text{net}}$$

$\Rightarrow$  Drawing the cycle boxes is the one place where we ignore usual first law sign convention. Pos. in direction of arrow.

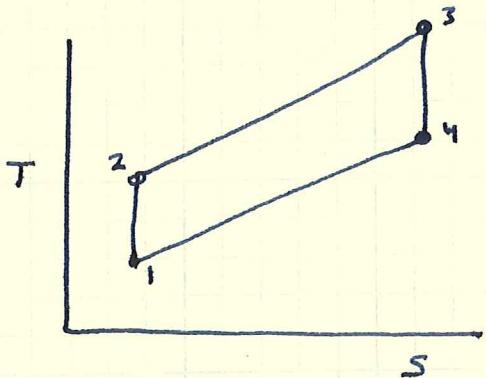
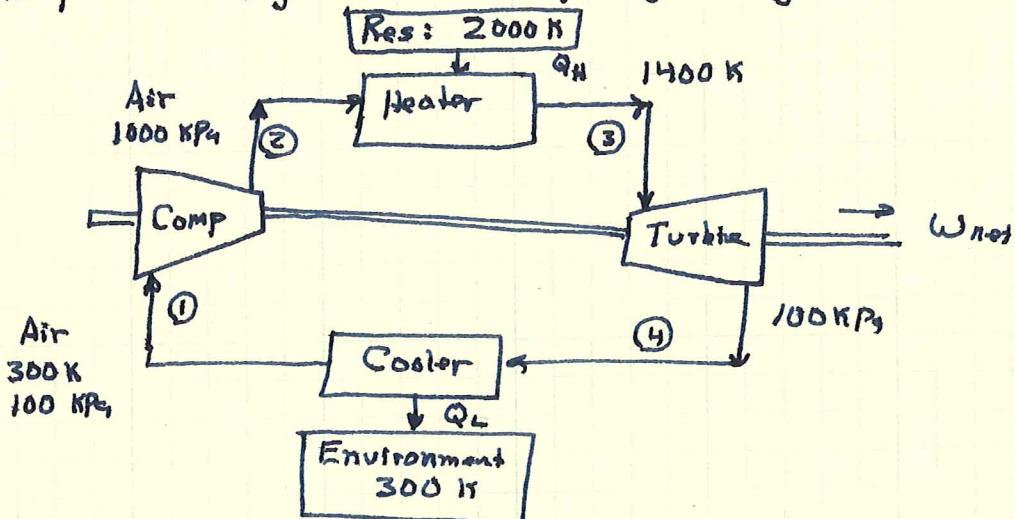
$Q_{\text{in}}$  = Heating value associated with the fuel.

$W_{\text{net}}$  = shaft work.

$Q_{\text{out}}$  = hot exhaust + radiator.

May 22

Let's put these together into a simple cycle analysis: Gas turbine cycle.



Add 1 → 2 Compress:  $\Delta S = 0$

2 → 3 Add heat,  $\Delta P = 0$

$$\Delta S = \int \frac{\delta Q}{T} \quad Q > 0$$

3 → 4 Expand,  $\Delta S = 0$

4 → 1 Lose heat,  $\Delta P = 0$

**1 → 2** Isentropic Compression:  $T_1 = 300 \text{ K} \Rightarrow h_1 = 300.19 \text{ kJ/kg}$

$$\frac{P_2}{P_1} = \frac{P_{r,2}}{P_{r,1}} \Rightarrow P_{r,2} = P_{r,1} \left[ \frac{P_2}{P_1} \right] = 1.386 \left[ \frac{1000}{100} \right] = 13.86$$

**A-17, 300 K**

$T_2 = 574.1 \text{ K}$

$h_2 = 579.9 \text{ kJ/kg}$

$$W_{\text{comp}} = h_2 - h_1 = 579.9 - 300.19 = 279.7 \text{ kJ/kg}$$

$$\text{Q} = h_3 - h_2 = 1515.42 - 579.9 = 935.5 \text{ kJ/kg}$$

$$\frac{P_4}{P_3} = \frac{P_{r,4}}{P_{r,3}} \Rightarrow P_{r,4} = P_{r,3} \left[ \frac{P_4}{P_3} \right] = 450.5 \left[ \frac{100}{1000} \right] = 45.05$$

$T_4 = 787.7 \text{ K}$

$$W_{\text{Turb}} = h_3 - h_4 = 1515.42 - 808.5 = 706.9 \text{ kJ/kg}$$

$$\text{Q} = h_4 - h_1 = 808.5 - 300.19 = 508.3 \text{ kJ/kg}$$

$$\text{Net Work} = W_{\text{Turb}} - W_{\text{comp}} = 706.9 - 279.7 = 427.2 \text{ kJ/kg}$$

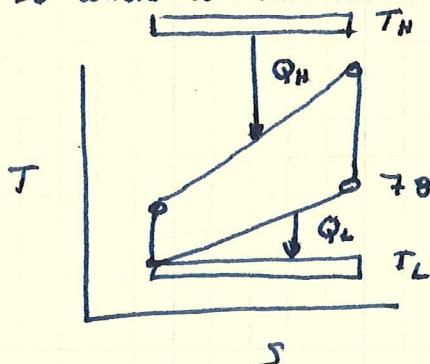
$$\eta = 1 - \frac{Q_{\text{out}}}{Q_{\text{in}}} = 1 - \frac{508.3}{935.5} = 45.7\%$$

How about a Carnot between the same two reservoirs?

$$\eta = 1 - \frac{T_L}{T_H} = 1 - \frac{300}{2000} = 85\% \Rightarrow \text{Best any engine could do.}$$

So where is the ~~s~~

So where is the source of inefficiency?



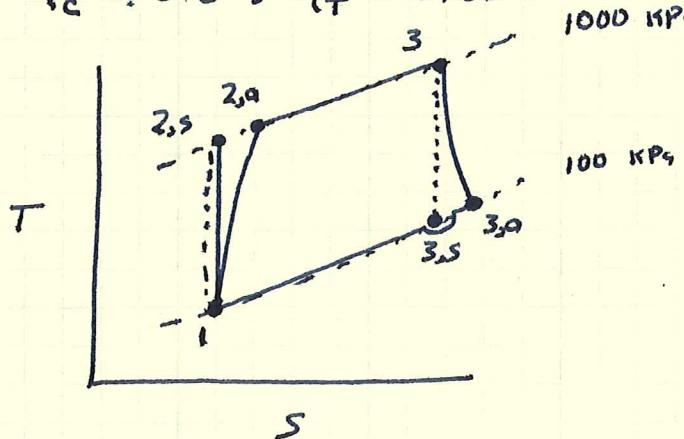
- It is the transfer of heat across finite  $\Delta T \Rightarrow$  Lost work, like water over dam.

- Which is bigger problem...  $Q_H$  or  $Q_L$ ?  
 $\Rightarrow$  Don't answer.  
 $\Rightarrow$  It's  $Q_L$ . (Chapter 7).

- What to do? Insert cycle... use this energy.

Now look at finite ~~cycle~~ compressor & turbine efficiencies.

$$\eta_c = 0.8 ; \eta_t = 0.85$$



$$\textcircled{1} \quad 100 \text{ kPa}, 300 \text{ K} \Rightarrow h_1 = 300.19 \text{ kJ/kg}$$

$$\boxed{1 \rightarrow 2s} \quad \text{same as before} \quad \frac{P_2}{P_1} = \frac{P_{T,2}}{P_{T,1}} \Rightarrow h_{2,s} = 579.9 \text{ kJ/kg}$$

$$\eta_c = \frac{W_{\text{ideal}}}{W_{\text{actual}}} = \frac{h_{2,s} - h_1}{h_{2,a} - h_1} \Rightarrow h_{2,a} = h_1 + \frac{h_{2,s} - h_1}{\eta_c}$$

$$h_{2,a} = 300.19 + \frac{579.9 - 300.19}{0.8} = 649.83 \text{ kJ/kg.}$$

$$W_{1 \rightarrow 2} = h_{2,a} - h_1 = 649.83 - 300.19 = 349.64 \text{ kJ/kg} \quad [279.7 \text{ kJ/kg}]$$

Think about this for a second:  $\begin{cases} \text{Both ideal and real compressors go from } 100 \rightarrow 1000 \text{ kPa, ideal needs less work,} \\ \text{Real needs more, this shows up as extra } h \text{ (higher T) -- Think of as internal friction.} \end{cases}$

2a → 3

$$g_{in} = h_3 - h_{2a} = 1515.42 - 649.93 = 865.59 \text{ kJ/kg.}$$

3 → 4a

3 → 4s; same as before;

$$\frac{P_4}{P_3} = \frac{P_{n4}}{P_{n3}} \Rightarrow h_{4s} = 808.5 \text{ kJ/kg}$$

$\Delta s = 0$

~~to~~

$$\eta_T = \frac{\omega_{actual}}{\omega_{ideal}} = \frac{h_3 - h_{4,a}}{h_3 - h_{4,s}} \Rightarrow h_{4,a} = h_3 - \eta_T(h_3 - h_{4,s})$$

$$h_{4,a} = 1515.42 - 0.85[1515.42 - 808.5] = 914.54 \text{ kJ/kg.}$$

$$\Rightarrow \omega_{T,a} = h_3 - h_{4,a} = 1515.42 - 914.54 = 600.88 \text{ kJ/kg}$$

[706.9 kJ/kg]

4a → 1

$$g = h_{4,a} - h_1 = 914.54 - 300.19 = 614.35 \text{ kJ/kg.}$$

$$\Rightarrow \text{Net Work} = \omega_T - \omega_c = \frac{600.88}{614.35} - 349.64 = 264.74 \text{ kJ/kg}$$

[427.2 kJ/kg]

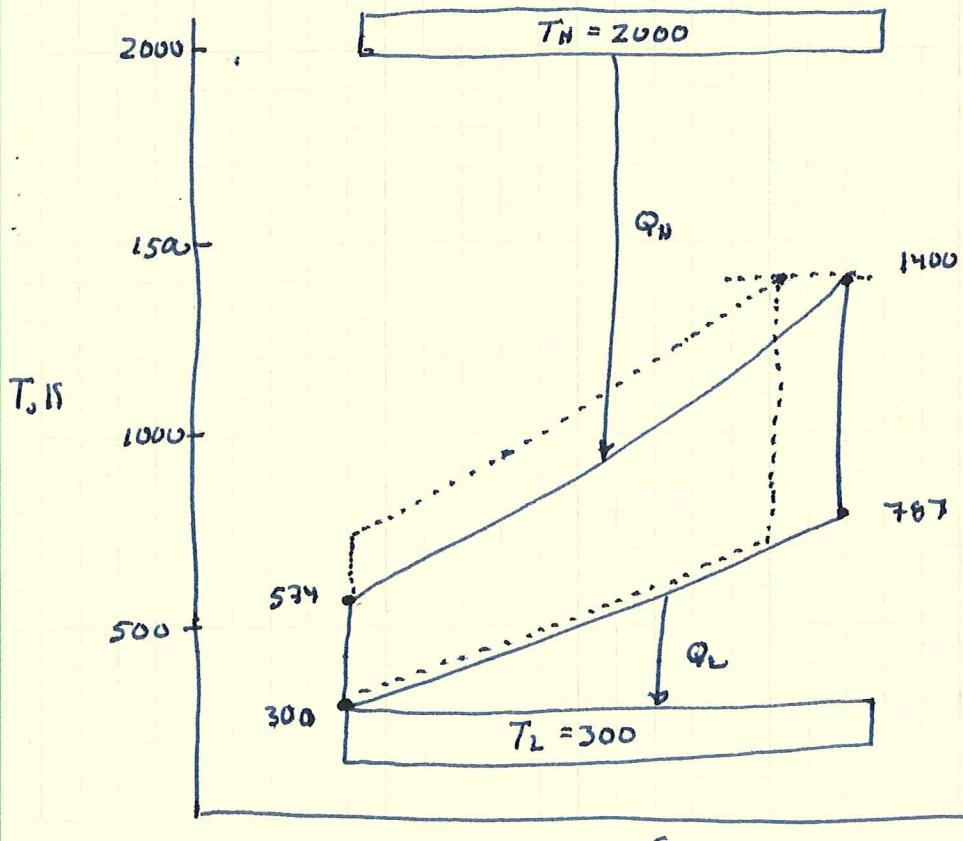
$$\eta = 1 - \frac{Q_{out}}{Q_{in}} = 1 - \frac{614.35}{865.59} = 29\%$$

[45.7%]

Lets think qualitatively about ways to improve cycle efficiency, say  $\eta_T = \eta_c = 1$

May 24

Draw to scale



### Reservoirs Fixed

① Reduce  $\Delta T$  for C's

⇒ higher compression ratio.

$\Delta T$  for both  $Q_H$  and  $Q_L$  are lower.

⇒ Increase turbine inlet temperature  $> 1400\text{K}$ .  
(Materials limits)

②  $T_4 > T_2$ , use exhaust to preheat gas from compressor outlet.